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# Energy Bands of Vanadium Oxide.

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## ABSTRACT

Energy bands of vanadium oxide have been calculated by the tight binding method. For this purpose the method of Lafon and Lin has been extended to include 3p, 3d, and 4s functions, so that three-center integrals could be included in the calculations. All matrix elements were summed to convergence, and the core functions of both vanadium and oxygen are explicitly included.

The crystal potential is constructed from a superposition of spherically averaged atomic potentials. The exchange potential used consists of the  $\rho^{1/3}$  approximation of Slater, together with an adjustable factor  $\lambda$ .

Band structures are calculated using three different crystal potentials. The density of states is computed from a sample of 505 points in  $1/48$  of the Brillouin zone. The effect of including core states is examined in detail, as are the effects of varying the factor  $\lambda$  in the exchange potential. Using the three different band structures which have been computed, an estimate of the self-consistent potential is made.

# CHAPTER I

## INTRODUCTION

In any investigation of the electronic properties of the solid state, it is necessary to make some approximations in order to overcome the mathematical difficulties involved. One approach to the problem is that of energy band theory, in which one considers the energy states of a single electron in a rigid, infinite, periodic lattice. The other electrons are treated as a group of noninteracting particles in the periodic potential of the lattice. Valence and conduction electrons are assumed to be in Bloch states; that is, states extended throughout the crystal and obeying the Bloch condition that

$$\psi_K(r + R_1) = e^{iK \cdot R_1} \psi_K(r)$$

where  $\psi_K(r)$  is the wave function,  $R_1$  is a direct lattice vector, and  $K$  is a vector in the reciprocal lattice (Bloch, 1928). This approach has been successful in describing the properties of a wide variety of materials.

Within the approximations imposed by energy band theory, there are several methods of calculation. One is interested in solving the one electron Schrodinger equation

$$H\psi_K = E_K\psi_K ,$$

where  $H$  is the sum of the kinetic energy operator plus whatever crystal potential is being used. The various methods of calculation attack this

problem by expanding the wave function in terms of some set of known functions, such as plane waves, spherical waves, Hartree-Fock atomic wave functions, or combinations of these.

In this paper we are concerned with the energy band structure of vanadium monoxide (VO), as determined by the tight binding method. In this scheme we expand the wave function  $\psi_K$  in terms of self-consistent atomic wave functions. The details of the tight binding method are presented in the second chapter, and we will not discuss them further here. Instead we pass on to a discussion of transition-metal oxides, with particular attention to be paid to vanadium monoxide.

The transition-metal oxides have been responsible for some of the most notable failures of the band theory technique. Perhaps the best known case is that of nickel oxide (NiO). This compound has a set of partially filled overlapping d-bands and should therefore be a metal according to band theory. In reality, it is a good insulator. The same situation occurs in the monoxides of manganese (MnO) and cobalt (CoO).

The properties of many compounds in this class can, however, be explained on the basis of energy band theory. In particular, this should be true of the oxides of vanadium (Adler, 1968). The vanadium oxides are of interest because many of them undergo sudden metal-semiconductor transitions as temperature is lowered. The change in conductivity can be as large as  $10^7$  within a few hundredths of a degree change in temperature.

Vanadium monoxide was chosen for consideration because it was reported to undergo a very sharp transition and because it has the simplest crystal structure. At room temperature it has the sodium chloride structure, and it is stable over a wide range of composition, from  $\text{VO}_{0.87}$

to  $\text{VO}_{1.20}$  (Westman and Nordmark, 1960). Stoichiometric samples show a high percentage of defects on the cation and anion lattices. It has been concluded by Ariya and Golomalzina (1962), on the basis of infrared spectra, that the VO lattice consists of alternating regions of composition  $\text{V}_2\text{O}_3$  and pure vanadium.

The lattice parameter is a function of composition, of course, and several different ranges for it has been found by different investigators. Typical values are in the range  $4.024\text{\AA} - 4.123\text{\AA}$  for the composition range quoted above (Stringer, 1964). The value  $4.062$  was used in this calculation.

The low temperature structure is not known. It has been suggested (Schonberg, 1954) that VO breaks down into a body-centered tetragonal phase and some other phase. It was later reported that this other phase is a body-centered cubic structure (Gurevich and Ormont, 1957) or a body-centered tetragonal phase with a  $c/a$  ratio very close to unity (Westman and Nordmark, 1960). An even later study (Geld et al., 1961) concluded that this phase did not exist.

Electrical properties of VO have been studied by Morin (1959), who found that the material was semiconducting at low temperatures. At a temperature  $T_0 = 126^\circ\text{K}$ , it abruptly became metallic, with a conductivity increase of  $10^6$ . Austin (1962) also found a semiconductor-metal transition at  $T_0 = 126^\circ\text{K}$  with the same increase in conductivity. His sample, however, was not exactly VO but was closer to  $\text{VO}_{0.9}$ .

Kawano et al. (1966) studied samples of various compositions in the temperature range  $100^\circ - 250^\circ\text{K}$ , and did not find a transition. He found VO to be metallic throughout the range. Kawano also measured magnetic susceptibility in the temperature ranges  $1.2 - 35^\circ\text{K}$  and  $77^\circ\text{K} -$  to

300°K. He found that  $\text{VO}_{1.147}$  and  $\text{VO}_{1.257}$  behaved like antiferromagnets having Neel temperatures 7.0° and 4.6°K respectively. No such behavior was found for the other samples.

A later study by Warren et al. (1967) found a discontinuity of  $10^4$  in electrical conductivity at a temperature near 125°K. They also obtained evidence for the presence of antiferromagnetism at low temperatures.

It is readily apparent that a good deal of work remains to be done experimentally before the low temperature phase of VO can be described theoretically. Although present evidence tends to favor the existence of the metal-insulator transition, Adler (1968) describes two other situations which also fit the observed data. The first is that VO is actually a metal at all temperatures, but at some temperature  $T_0$  it decomposes into the two phases originally suggested by Klemm & Grimm (1940); that is, a body-centered tetragonal structure of composition  $\text{V}_3\text{O}$  and a higher oxide of composition  $\text{V}_3\text{O}_4$ . The phase boundaries in this material would then lead to the high resistivity. The second possibility is that samples of VO contain regions of  $\text{V}_2\text{O}_3$  in filamentary form, which would imply that the transition observed is actually due to the  $\text{V}_2\text{O}_3$ . This is supported by the striking resemblance between VO and  $\text{V}_2\text{O}_3$  in both electrical and magnetic behavior. The fact that VO and  $\text{V}_2\text{O}_3$  exhibit significantly different transition temperatures seems to discourage this view. However, if one views the VO system as being  $\text{V}_2\text{O}_3$  with a large vanadium excess, the crystalline distortion model which Adler applied to  $\text{V}_2\text{O}_3$  (1967) would predict a lowering of the transition temperature. This possibility can be checked by determining the crystal structure for some crystal which actually undergoes the transition.

Assuming that the metal-insulator transition in VO does exist, we still cannot propose a model with any degree of confidence until the low temperature structure is determined. There exist two possibilities which are attractive, however. A typical face-centered cubic type antiferromagnetism at low temperatures would split the sixfold-degenerate  $\Gamma'_{25}$  band into two threefold degenerate bands. Since VO has 3 3d electrons per cation, this could explain its low temperature behavior. This model, of course, depends upon the transition taking place at the Neel temperature. Another possibility is the crystal distortion model proposed by Adler (1967) for  $V_2O_3$ . In the case of VO Adler suggests that the distortion is a rhombohedral pairing of alternate (111) planes. Of course a third model could be a combination of the two. Since it is not certain that VO actually undergoes a metal-insulator transition, no model will be proposed in this paper.

The second chapter of this paper will be devoted to a discussion of the tight binding method and the crystal potential which was used in this calculation. The third chapter will be concerned with the details of the calculation, the fourth chapter will give the results obtained, and the fifth chapter will give the conclusions to be drawn.

## CHAPTER II

In this chapter we will discuss the tight binding, or LCAO, method of calculating energy bands. The matrix elements for this method will be derived and the integrals necessary to compute them will be given. The crystal potential and exchange potential will also be discussed. Much of the following discussion is taken from Callaway (1964).

The tight binding method was proposed by Bloch in 1928. Until recently, however, use of the method has been largely restricted either to calculations of a highly empirical nature or to ones in which it serves as an interpolation scheme after energy bands have been calculated at symmetry points by other methods. The technique has been improved, however, by Lafon and Lin (1966) and this, together with the increased efficiency of modern computers, has made the tight binding approximation an effective tool for calculating energy bands.

The method is based upon the expansion of the crystalline electronic wave function in terms of bound atomic states. The atomic wave functions must be combined to form functions which have the periodicity required by Bloch's theorem. It can be easily shown that the Bloch function  $\phi_K$  given by

$$\phi_K(\underline{r}) = \frac{1}{\sqrt{N}} \sum_{\underline{v}} e^{i\mathbf{K} \cdot \mathbf{R}_{\underline{v}}} u(\underline{r} - \mathbf{R}_{\underline{v}})$$

where  $u(\underline{r} - \mathbf{R}_{\underline{v}})$  is an atomic function centered at lattice site  $\mathbf{R}_{\underline{v}}$ , does satisfy the Bloch condition

$$\phi_K(\underline{r} + \mathbf{R}_{\underline{1}}) = e^{i\mathbf{K} \cdot \mathbf{R}_{\underline{1}}} \phi_K(\underline{r}) ,$$

where  $\underline{R}_j$  is a direct lattice vector.

The set consisting of only the bound atomic states does not form a complete set, so an expansion of the wave function using this set as a basis does not lead to an exact solution of the Schrodinger equation. However, the inclusion of all the bound states can be expected to yield a good approximation to the actual wave function.

We will now consider the matrix elements which occur in the tight binding formulation. Since atomic states on different atoms are not orthogonal, the secular determinant has the form

$$|H_{i,j}(K) - E S_{i,j}(K)| = 0,$$

where  $H_{i,j}$  and  $S_{i,j}$  are the matrix elements of the Hamiltonian and of unity with respect to the basis of Bloch functions formed from atomic functions having symmetry properties  $i$  and  $j$ , respectively. The expressions for these matrix elements will now be derived.

Writing the Hamiltonian to yield energy in Rydbergs, we have

$$\begin{aligned} H &= -\nabla^2 + \sum_{\underline{v}} V(\underline{r} - \underline{R}_{\underline{v}}) + VX \\ &= -\nabla^2 + V + VX \end{aligned}$$

where  $V(\underline{r} - \underline{R}_{\underline{v}})$  is the potential centered on the atom at lattice site  $\underline{R}_{\underline{v}}$  and  $VX$  is the exchange. We will temporarily ignore the exchange term.

$$\begin{aligned} H_{i,j}(K) &= \int \phi_{K,i}^* H \phi_{K,j} d^3r \\ &= \int \frac{1}{\sqrt{N}} \sum_{\underline{v}} e^{-i\mathbf{K} \cdot \underline{R}_{\underline{v}}} u_i^*(\underline{r} - \underline{R}_{\underline{v}}) [-\nabla^2 + V] \\ &\quad \frac{1}{\sqrt{N}} \sum_{\underline{v}''} e^{i\mathbf{K} \cdot \underline{R}_{\underline{v}''}} u_j(\underline{r} - \underline{R}_{\underline{v}''}) d^3r \end{aligned}$$



$$= \frac{1}{N} \sum_{\mathbf{v}, \mathbf{v}''} e^{-i\mathbf{K} \cdot (\mathbf{R}_{\mathbf{v}} - \mathbf{R}_{\mathbf{v}''})} \int u_1^*(\mathbf{r} - \mathbf{R}_{\mathbf{v}}) [-\nabla^2 + V] u_j(\mathbf{r} - \mathbf{R}_{\mathbf{v}''}) d^3r$$

Now the summations over  $\mathbf{v}$  and  $\mathbf{v}''$  are both over the entire crystal, so one of them is redundant. We can simplify the expression by setting one of them (say  $\mathbf{v}''$ ) equal to 0 and deleting the multiplier  $N$ . This yields

$$\begin{aligned} H_{1,j}(\mathbf{K}) &= \sum_{\mathbf{v}} e^{-i\mathbf{K} \cdot \mathbf{R}_{\mathbf{v}}} \int u_1^*(\mathbf{r} - \mathbf{R}_{\mathbf{v}}) [-\nabla^2 + V] u_j(\mathbf{r}) d^3r \\ &= \sum_{\mathbf{v}} e^{-i\mathbf{K} \cdot \mathbf{R}_{\mathbf{v}}} \int u_1^*(\mathbf{r} - \mathbf{R}_{\mathbf{v}}) (-\nabla^2) u_j(\mathbf{r}) d^3r \\ &\quad + \sum_{\mathbf{v}} e^{-i\mathbf{K} \cdot \mathbf{R}_{\mathbf{v}}} \int u_1^*(\mathbf{r} - \mathbf{R}_{\mathbf{v}}) V u_j(\mathbf{r}) d^3r \\ &= T_{1,j}(\mathbf{K}) + V_{1,j}(\mathbf{K}) \end{aligned}$$

where

$$T_{1,j}(\mathbf{K}) = \sum_{\mathbf{v}} e^{-i\mathbf{K} \cdot \mathbf{R}_{\mathbf{v}}} \int u_1^*(\mathbf{r} - \mathbf{R}_{\mathbf{v}}) (-\nabla^2) u_j(\mathbf{r}) d^3r \quad (1)$$

$$V_{1,j}(\mathbf{K}) = \sum_{\mathbf{v}} e^{-i\mathbf{K} \cdot \mathbf{R}_{\mathbf{v}}} \int u_1^*(\mathbf{r} - \mathbf{R}_{\mathbf{v}}) V u_j(\mathbf{r}) d^3r \quad (2)$$

The matrix element of unity is given by

$$\begin{aligned} S_{1,j}(\mathbf{K}) &= \int \phi_{K,1}^* \phi_{K,j}^* d^3r \\ &= \frac{1}{N} \sum_{\mathbf{v}, \mathbf{v}'} e^{-i\mathbf{K} \cdot (\mathbf{R}_{\mathbf{v}} - \mathbf{R}_{\mathbf{v}'})} \int u_1^*(\mathbf{r} - \mathbf{R}_{\mathbf{v}}) u_j^*(\mathbf{r} - \mathbf{R}_{\mathbf{v}'}) d^3r \\ &= \sum_{\mathbf{v}} e^{-i\mathbf{K} \cdot \mathbf{R}_{\mathbf{v}}} \int u_1^*(\mathbf{r} - \mathbf{R}_{\mathbf{v}}) u_j(\mathbf{r}) d^3r \quad (3) \end{aligned}$$

The integrals in both  $T_{i,j}(K)$  and  $S_{i,j}(K)$  are two-center integrals; that is, they consist of one wave function centered on a site  $R_v$  and another on the origin, with an operator which is spatially independent between them. Standard techniques of performing these integrals have been devised and are in general use, as we will discuss shortly.

The potential energy integrals in  $V_{i,j}(K)$  are three-center integrals; that is, the operator between the wave functions depends upon position also. Evaluation of these integrals is therefore much more difficult. It is the problems associated with the computation of these integrals that have restricted the use of the tight binding method.

The contribution of Lin and Lafon was to devise a scheme to compute these integrals. It is based on the technique of Gaussian transformation, and will be covered in detail here.

We proceed by expanding the potential  $V$  in a Fourier series. As we will discuss shortly,  $V$  is given by a sum of atomic potentials centered on each lattice site and so has the periodicity of the lattice. Hence the Fourier expansion of  $V$  contains only reciprocal lattice vectors. We choose our origin on a vanadium atom, which gives inversion symmetry. Therefore the Fourier series contains only cosine terms. We can then write

$$V(\underline{r}_C) = \sum_{K_v} V(K_v) \cos(K_v \cdot \underline{r}_C)$$

where  $V(K_v)$  is a Fourier coefficient of the potential, the summation is over all reciprocal lattice vectors, and  $\underline{r}_C$  is the radius vector from any given vanadium site in the crystal.

We will now consider the integral of two 1s orbitals centered on sites A and B with the potential on site C. Using the Fourier series expansion, this becomes

$$\langle 1s(A) | V | 1s(B) \rangle = \sum_{K_v} V(K_v) \langle 1s(A) | \cos \underline{K_v} \cdot \underline{r_C} | 1s(B) \rangle.$$

The wave functions used in this calculation consist of linear combinations of Slater type orbitals, so that a 1s orbital has the form  $\exp(-\alpha r)$ . Expressing the 1s orbitals in a Laplace transform gives

$$e^{-\alpha_1 r_A} = \frac{\alpha_1}{2\sqrt{\pi}} \int_0^\infty s_1^{-\frac{3}{2}} \exp\left(-\frac{\alpha_1^2}{4s_1} - r_A^2 s_1\right) ds_1$$

and similarly for the orbital on site B. This result, as well as results for several other integrals to be used in this derivation, is to be found in Tables of Integral Transforms, edited by Erdeli (1954). Putting in these expressions gives

$$\begin{aligned} \langle 1s(A) | \cos \underline{K_v} \cdot \underline{r_C} | 1s(B) \rangle &= \frac{\alpha_1 \alpha_2}{4\pi} \int_0^\infty (s_1 s_2)^{-\frac{3}{2}} \exp\left(-\frac{\alpha_1^2}{4s_1} - \frac{\alpha_2^2}{4s_2}\right) \\ &\quad \int \exp(-s_1 r_A^2 - s_2 r_B^2) \cos \underline{K_v} \cdot \underline{r_C} d\tau ds_1 ds_2 \end{aligned} \quad (4)$$

The product of two Gaussians centered on sites A and B can be written in terms of a third Gaussian centered at a point D along the line AB; that is,

$$\exp(-s_1 r_A^2 - s_2 r_B^2) = \exp\left(-\frac{s_1 s_2 r_{AB}^2}{s_1 + s_2}\right) \exp(-(s_1 + s_2) r_D^2)$$

where

$$D_i = \frac{s_1 A_i + s_2 B_i}{s_1 + s_2}, \quad i = x, y, z.$$

By writing the vector  $r_C$  as

$$r_C = r_D + r_{CD}$$

and expanding the cosine as

$$\cos \underline{K}_v \cdot \underline{r}_C = \cos \underline{K}_v \cdot \underline{r}_D \cos \underline{K}_v \cdot \underline{r}_{CD} - \sin \underline{K}_v \cdot \underline{r}_D \sin \underline{K}_v \cdot \underline{r}_{CD}$$

the spatial integral in equation (4) can be evaluated.

$$\begin{aligned} \int \exp(-s_1 r_A^2 - s_2 r_B^2) \cos \underline{K}_v \cdot \underline{r}_C d\tau = \\ \exp\left(-\frac{s_1 s_2 r_{AB}^2}{s_1 + s_2}\right) [\cos \underline{K}_v \cdot \underline{r}_{CD} \int \exp[-(s_1 + s_2) r_D^2] \cos \underline{K}_v \cdot \underline{r}_D d\tau \\ - \sin \underline{K}_v \cdot \underline{r}_{CD} \int \exp[-(s_1 + s_2) r_D^2] \sin \underline{K}_v \cdot \underline{r}_D d\tau] \end{aligned} \quad (5)$$

The last integral in equation (5) is the integral of an odd function throughout all space, and so it vanishes. Then equation (5) becomes

$$\begin{aligned} \int \exp(-s_1 r_A^2 - s_2 r_B^2) \cos \underline{K}_v \cdot \underline{r}_C d\tau = \\ \exp\left(-\frac{s_1 s_2 r_{AB}^2}{s_1 + s_2}\right) \cos \underline{K}_v \cdot \underline{r}_{CD} 2\pi \int_0^\infty \int_0^\infty \exp[-(s_1 + s_2) r_D^2] \\ \cos(\underline{K}_v \cdot \underline{r}_D \cos \theta) r_D^2 \sin \theta d\theta dr_D \end{aligned}$$

$$\begin{aligned}
&= \exp \left( - \frac{s_1 s_2 r_{AB}^2}{s_1 + s_2} \right) \cos \tilde{K}_v \cdot \tilde{r}_{CD} \frac{4\pi}{K_v} \int_0^\infty r_D \exp [-(s_1 + s_2) r_D^2] \\
&\quad \sin K_v r_D \, dr_D \\
&= \exp \left( - \frac{s_1 s_2 r_{AB}^2}{s_1 + s_2} \right) \cos \tilde{K}_v \cdot \tilde{r}_{CD} \left( \frac{\pi}{s_1 + s_2} \right)^{\frac{3}{2}} \exp \left( - \frac{K_v^2}{4(s_1 + s_2)} \right) \quad (6)
\end{aligned}$$

Now we make the following change of variable;

$$Z = s_1 + s_2$$

$$u = s_1 / (s_1 + s_2)$$

$$r = u(1-u)r_{AB}^2$$

$$g = K_v^2 + \alpha_2^2 / (1-u) + \alpha_1^2 / u$$

The Jacobian of this transformation is given by

$$\begin{aligned}
\frac{d(s_1, s_2)}{d(u, Z)} &= \begin{vmatrix} \frac{ds_2}{du} & \frac{ds_1}{dZ} \\ \frac{ds_1}{du} & \frac{ds_2}{dZ} \end{vmatrix} \\
&= Z
\end{aligned}$$

Using this fact together with the evaluation of the spatial integral which we completed in equation (6), and noting that

$$1-u = \frac{s_2}{s_1 + s_2},$$

it is a simple matter to reduce the original three-center integral to

$$\langle 1s(A) | \cos \mathbf{K}_v \cdot \mathbf{r}_C | 1s(B) \rangle =$$

$$\frac{1}{4}\pi^{1/2} \alpha_1 \alpha_2 \int_0^1 [u(1-u)]^{3/2} \cos \mathbf{K}_v \cdot \mathbf{r}_{CD} du \int_0^\infty z^{-7/2} \exp(-fz - \frac{g}{4z}) dz =$$

$$2\pi \alpha_1 \alpha_2 r_{AB}^3 \int_0^1 F \left[ \frac{3}{(fg)^{5/2}} + \frac{3}{(fg)^2} + \frac{1}{(fg)^{3/2}} \right] \cos \mathbf{K}_v \cdot \mathbf{r}_{CD} e^{-(fg)^{1/2}} du$$

Now we have reduced a three dimensional integral over all space to a one dimensional integral over the interval zero to one, simply by expanding the potential in a Fourier series and by using the technique of Gaussian transformation. The integral may now be easily evaluated by numerical methods. The only price we must pay for this reduction is that we have to sum the Fourier series to convergence, calculating a new integral for each value of  $\mathbf{K}_v$ . This is quite feasible, as will be discussed later.

Potential integrals involving higher Slater orbitals than  $1s$  can be obtained from this integral by performing the proper partial differentiations with respect to  $\alpha_1$ ,  $\alpha_2$ ,  $A_x$ ,  $A_y$ ,  $A_z$ ,  $B_x$ ,  $B_y$ ,  $B_z$ . For example,

$$\begin{aligned} \langle 2P_x(A) | \cos \mathbf{K}_v \cdot \mathbf{r}_C | 1s(B) \rangle = \\ \int x_A e^{-\alpha_1 r_A} \cos \mathbf{K}_v \cdot \mathbf{r}_C e^{-\alpha_2 r_B} d\tau = \\ - \frac{d}{d\alpha_1} \left( \frac{1}{\alpha_1} \frac{d}{dA_x} \langle 1s(A) | \cos \mathbf{K}_v \cdot \mathbf{r}_C | 1s(B) \rangle \right). \end{aligned}$$

For a more detailed treatment of this procedure and a tabulation of the integrals involved, see Appendix A.

Integrals computed in this fashion are compiled as previously indicated in equation (2) to form the matrix element  $V_{i,j}(K)$ .

The overlap integrals appear as the first term of the sum over the reciprocal lattice vectors and are compiled according to equation (3) to form the matrix element  $S_{i,j}(K)$ . The kinetic energy integrals could be obtained by applying the operator

$$T = -\nabla^2$$

$$= - \left( \frac{d^2}{dB_x^2} \frac{d^2}{dB_y^2} \frac{d^2}{dB_z^2} \right)$$

to the corresponding overlap integrals. However, methods of computing these integrals have been developed along the lines begun by Roothaan (1951). For this calculation a computer program written by Roney et al. (1968) was used for this purpose. The kinetic energy integrals were then combined according to equation (1) to form the matrix element  $T_{i,j}(K)$ .

This concludes our discussion of the tight binding method. We now pass on to a discussion of the crystal potential.

It is appropriate at this point to discuss the structure of vanadium oxide. VO has the sodium chloride structure; that is, it is face-centered cubic with a basis of two atoms per unit cell. The primitive translation vectors of the direct lattice are

$$\underline{a}_1 = \frac{a}{2} (\underline{i} + \underline{j})$$

$$\underline{a}_2 = \frac{a}{2} (\underline{j} + \underline{k})$$

$$\underline{a}_3 = \frac{a}{2} (\underline{k} + \underline{i})$$

where  $a$  is the lattice constant.

The reciprocal lattice vectors are given by

$$\underline{b}_1 = \frac{2\pi}{a} (\underline{i} + \underline{j} - \underline{k})$$

$$\underline{b}_2 = \frac{2\pi}{a} (-\underline{i} + \underline{j} + \underline{k})$$

$$\underline{b}_3 = \frac{2\pi}{a} (\underline{i} - \underline{j} + \underline{k})$$

For this calculation the origin was chosen at a vanadium site. The oxygen atom in the unit cell is located at

$$\underline{R}_s = \frac{1}{2} (\underline{a}_1 + \underline{a}_2 + \underline{a}_3)$$

For a general reciprocal lattice vector given by

$$\underline{K}_n = n_1 \underline{b}_1 + n_2 \underline{b}_2 + n_3 \underline{b}_3$$

we have

$$\underline{R}_s \cdot \underline{K}_n = \pi(n_1 + n_2 + n_3)$$

With this background, we will consider the crystal potential.

The crystal potential was formed from a superposition of spherically averaged atomic potentials. That is,

$$V(\underline{r}) = \sum_{\underline{l}} V_v(\underline{r} - \underline{R}_l) + \sum_{\underline{m}} V_o(\underline{r} - \underline{R}_m)$$

where  $V_v(\underline{r} - \underline{R}_l)$  is an atomic potential centered on a vanadium atom at



site  $R_1$  and  $V_o(r - R_m)$  is an atomic potential centered on the oxygen atom at site  $R_m$ . For reasons already discussed, we are interested in expanding this potential in a Fourier series. The Fourier coefficients are given by

$$\begin{aligned} V(K) &= \frac{1}{N\Omega_o} \int V(r) e^{-iK \cdot r} d^3r \\ &= \frac{1}{N\Omega_o} \left[ \sum_l \int e^{-iK \cdot r} V_v(r - R_l) d^3r + \sum_m \int e^{-iK \cdot r} V_o(r - R_m) d^3r \right] \end{aligned}$$

where  $\Omega_o$  is the unit cell volume and  $N$  is the number of unit cells.

Now make a change of variable from  $r$  to  $r + R_l$  in the first integral and from  $r$  to  $r + R_m$  in the second.

$$\begin{aligned} V(K) &= \frac{1}{N\Omega_o} \left[ \sum_l \int e^{-iK \cdot (r + R_l)} V_v(r) d^3r + \sum_m \int e^{-iK \cdot (r + R_m)} V_o(r) d^3r \right] \\ &= \frac{1}{N\Omega_o} \left[ \sum_l e^{-iK \cdot R_l} \int V_v(r) e^{-iK \cdot r} d^3r + \sum_m e^{-iK \cdot R_m} \int V_o(r) e^{-iK \cdot r} d^3r \right] \end{aligned}$$

By our choice of origin in the unit cell, we have

$$R_m = R_l + R_s$$

Then

$$\begin{aligned} V(K) &= \frac{1}{N\Omega_o} \left[ \sum_l e^{-iK \cdot R_l} \int V_v(r) e^{-iK \cdot r} d^3r + \sum_l e^{-iK \cdot (R_l + R_s)} \int V_o(r) e^{-iK \cdot r} d^3r \right] \\ &= \frac{1}{N\Omega_o} \sum_l e^{-iK \cdot R_l} \left[ \int V_v(r) e^{-iK \cdot r} d^3r + e^{iK \cdot R_s} \int V_o(r) e^{-iK \cdot r} d^3r \right] \end{aligned}$$

Now

$$\sum_{\underline{1}} e^{-i\mathbf{K} \cdot \mathbf{R}_{\underline{1}}} = N \delta_{\mathbf{K}, \mathbf{K}_n}$$

where  $\mathbf{K}_n$  is a reciprocal lattice vector. As already pointed out,

$$\mathbf{K}_n \cdot \mathbf{R}_s = \pi(n_1 + n_2 + n_3)$$

Then

$$\begin{aligned} V(\mathbf{K}) &= \frac{1}{\Omega_0} \left[ \int V_v(\mathbf{r}) e^{-i\mathbf{K}_n \cdot \mathbf{r}} d^3r + e^{-i\mathbf{K}_n \cdot \mathbf{R}_s} \int V_o(\mathbf{r}) e^{-i\mathbf{K}_n \cdot \mathbf{r}} d^3r \right] \\ &= \frac{1}{\Omega_0} \left[ \int V_v(\mathbf{r}) e^{-i\mathbf{K}_n \cdot \mathbf{r}} d^3r + (-1)^{n_1+n_2+n_3} \int V_o(\mathbf{r}) e^{-i\mathbf{K}_n \cdot \mathbf{r}} d^3r \right] \end{aligned}$$

Each of the atomic potentials has the form

$$V(\mathbf{r}_2) = -\frac{Z}{r_2} + 2 \int \frac{\rho(\mathbf{r}_1)}{r_{12}} d^3r_1$$

where  $Z$  is the atomic number and  $\rho$  is the charge density. The factor of 2 converts the energy into Rydbergs. The charge density is given by

$$\begin{aligned} \rho(\mathbf{r}) &= \sum_1 a_1 |\psi_1|^2 \\ &= 2|\psi_{1s}|^2 + 2|\psi_{2s}|^2 + 6|\psi_{2p}|^2 + 2|\psi_{3s}|^2 + 6|\psi_{3p}|^2 + 3|\psi_{3d}|^2 + 2|\psi_{4s}|^2 \end{aligned}$$

for vanadium in the  $3d^3 4s^2$  configuration. The  $a_1$ 's are the occupation numbers of the 1-th states. For vanadium,

$$v(r_2) = -\frac{46}{r_2} + 2 \int \frac{\rho(r_1)}{|r_{12}|} d^3r_1$$

The Fourier coefficient of the vanadium potential is

$$\begin{aligned} v_v(k) &= \int v_v(r_2) e^{-i\mathbf{k} \cdot \mathbf{r}_2} d^3r_2 \\ &= -46 \int \frac{e^{-i\mathbf{k} \cdot \mathbf{r}_2}}{r_2} d^3r_2 + 2 \iint \frac{\rho(r_1)}{|r_{12}|} e^{-i\mathbf{k} \cdot \mathbf{r}_2} d^3r_1 d^3r_2 \end{aligned}$$

Use the identity

$$\int \frac{e^{-i\mathbf{k} \cdot \mathbf{r}_2}}{|r_{12}|} d^2r_2 = \frac{4\pi}{k^2} e^{-i\mathbf{k} \cdot \mathbf{r}_1}.$$

Then

$$\begin{aligned} v_v(k) &= -\frac{184\pi}{k^2} + \frac{8\pi}{k^2} \int \rho(r_1) e^{-i\mathbf{k} \cdot \mathbf{r}_1} d^3r_1 \\ &= -\frac{184\pi}{k^2} + \frac{32\pi^2}{k^3} \int_0^\infty \rho(r_1) r_1 \sin(kr_1) dr_1 \end{aligned}$$

Now take a spherical average of the charge density.

$$\rho(r_1) = \frac{1}{4\pi} \sum_l a_l |u_l|^2$$

where  $u_l$  is the radial part of the wave function for the  $l$  - th state.

Define

$$\begin{aligned} f_v(r_1) &= 2|u_{1s}|^2 + 2|u_{2s}|^2 + 6|u_{2p}|^2 + 2|u_{3s}|^2 + 6|u_{3p}|^2 \\ &\quad + 3|u_{3d}|^2 + 2|u_{4s}|^2 \end{aligned}$$

Then

$$V_v(K) = -\frac{184\pi}{K^2} + \frac{8\pi}{K^3} \int_0^{\infty} r_1 f_v(r_1) \sin(Kr_1) dr_1$$

Similarly

$$V_o(K) = -\frac{64\pi}{K^2} + \frac{8\pi}{K^3} \int_0^{\infty} r_1 f_o(r_1) \sin(Kr_1) dr_1$$

where

$$f_o(r_1) = 2|u_{1s}|^2 + 2|u_{2s}|^2 + 4|u_{2p}|^2$$

From these expressions it is not at all obvious that  $\lim_{K \rightarrow 0} V(K)$  exists. To show this, we will consider this term separately. Expand the sine term in a power series.

$$\begin{aligned} V_v(0) &= \lim_{K \rightarrow 0} V_v(K) \\ &= \lim_{K \rightarrow 0} \left[ -\frac{184\pi}{K^2} + \frac{8\pi}{K^3} \int_0^{\infty} f_v(r_1) \left[ Kr_1 - \frac{(Kr_1)^3}{3!} + \cdots \right] r_1 dr_1 \right] \\ &= -\frac{184\pi}{K^2} + \frac{8\pi}{K^3} \int_0^{\infty} K f_v(r_1) r_1^2 dr_1 - \frac{8\pi}{6K^3} K^3 \int_0^{\infty} f_v(r_1) r_1^4 dr_1 \\ &= -\frac{184\pi}{K^2} + \frac{184\pi}{K^2} - \frac{4\pi}{3} \int_0^{\infty} f_v(r_1) r_1^4 dr_1 \\ &= -\frac{4\pi}{3} \int_0^{\infty} f_v(r_1) r_1^4 dr_1 \end{aligned}$$

A similar expression holds, of course, for oxygen.

The radial parts of the wave functions consist of Slater-type orbitals, as tabulated by Clementi (1965). The various integrals involved in the potential can therefore be evaluated in closed form with the aid of Laplace transforms from Erdelyi (1954).

The exchange potential used consisted of the  $\rho^{1/3}$  approximation of Slater, together with a multiplicative factor of  $\lambda$  to be adjusted. The charge density for the crystal was found at selected points within spheres centered about vanadium and oxygen atoms and having radii equal to half the nearest neighbor distance. These points were chosen so that a 96 point Gaussian formula for integration could be employed. The charge density was found by summing contributions from atomic charge densities out through nine sets of neighbors.

The cube root of the total charge density was computed along three independent axes and the results were averaged to find the final exchange result. Fourier coefficients of the exchange were calculated within the spheres centered on the vanadium and oxygen atoms and these were combined in the same manner as with the crystal potential to form the full exchange coefficient. Therefore the full Hamiltonian matrix element is given by

$$H_{i,j}(K) = V_{i,j}(K) + T_{i,j}(K) + \lambda VX_{i,j}(K)$$

where  $VX_{i,j}(K)$  is the exchange matrix element.

### Chapter III

In this chapter we will be concerned with the details of the calculation. The use of group theory will be discussed first. Then we will consider the problems involved in the computation of the three-center integrals, such as the convergence of the various sums.

The group theory for the cubic lattice has been worked out in detail (Slater, 1965) and will not be discussed at length in this chapter. A few points of particular interest for this calculation will be found in Appendix B. Here we will be more concerned with the application of group theory than with the mathematical details.

It is useful to expand the crystal wave function in terms of functions which transform according to the rows of the irreducible representations of the cubic group. This is due to the fact that the matrix elements between functions which belong to different irreducible representations, or to different rows of the same irreducible representation, are zero. Mathematically speaking,

$$\begin{aligned}
 H_{s,t}^{p,q} &= \int f_s^{p*} H f_t^q d^3r \\
 &= \frac{1}{dq} \sum_{s,t} H_{s,t}^{p,q} d_{p,q} d_{s,t} \\
 H_{s,s}^{p,p} &= \frac{1}{dq} \sum_t H_{t,t}^{q,q}
 \end{aligned}$$

where  $f_s^p$  is a function transforming as the  $s$  - th row of the  $p$  - th irreducible representation and similarly for  $f_t^q$ . Also  $dq$  is the dimension of the  $q$  - th irreducible representation, and the sum is over

all rows of the  $q$  - th irreducible representation.

Therefore in energy band calculations it is frequently necessary to form linear combinations of the basis functions which will transform in the proper manner. In our case the basis consisted of atomic orbitals, which already have the correct symmetry properties for a cubic lattice, so this was unnecessary. It was still necessary to combine them into Bloch sums to insure that they have the proper periodicity, as was discussed in Chapter II in connection with the tight binding method.

The matrix elements of the Hamiltonian have been given as

$$\begin{aligned} H_{i,j}(K) &= \sum_{\mathbf{R}_1} e^{-i\mathbf{K} \cdot \mathbf{R}_1} \int u_i^*(\mathbf{r}-\mathbf{R}_1) [-\nabla^2 + V + \lambda VX] u_j(\mathbf{r}) d^3r \\ &= \sum_{\mathbf{R}_1} e^{-i\mathbf{K} \cdot \mathbf{R}_1} \int u_i^*(\mathbf{r}-\mathbf{R}_1) H u_j(\mathbf{r}) d^3(\mathbf{r}) \\ &= \sum_{\mathbf{R}_1} e^{-i\mathbf{K} \cdot \mathbf{R}_1} E_{i,j}(\mathbf{R}_1) \end{aligned}$$

where  $u_i(\mathbf{r}-\mathbf{R}_1)$  and  $u_j(\mathbf{r})$  are atomic orbitals centered at  $\mathbf{R}_1$  and the origin, respectively,  $V$  is the crystal potential,  $\lambda$  is an adjustable parameter, and  $VX$  is the exchange.

As it stands, the sum is over all direct lattice vectors, so that we must consider separately the contribution from each atom in the crystal. The number of terms in this series can be greatly reduced by the use of group theory, however. Let  $R$  be any operation of the cubic group. Then the following is true;

$$R^{-1} = R^{\dagger} = R^*$$

$$R^{-1} H R = H$$

where  $H$  is the Hamiltonian. The quantity  $E_{i,j}(R_1)$  is given by

$$\begin{aligned}
 E_{i,j}(R_1) &= \int u_i^*(r-R_1) H u_j(r) d^3r \\
 &= \int u_i^*(r-R_1) R^{-1} H R u_j(r) d^3r \\
 &= \int [R u_i(r-R_1)]^* H [R u_j(r)] d^3r
 \end{aligned}$$

As an example let

$$E_{i,j}(R_1) = E_{x^2-y^2, x^2-y^2}(abc) \quad ,$$

where  $(abc)$  is a general point of the direct lattice, and let  $R$  be the operator  $R_{19}$  as defined in Appendix B. Then from the transformation properties of the function  $x^2 - y^2$ , as given also in Appendix B, we have

$$\begin{aligned}
 E_{x^2-y^2, x^2-y^2}(abc) &= \int [R_{19} u_{x^2-y^2}(r-abc)]^* H [R_{19} u_{x^2-y^2}(r)] d^3r \\
 &= \frac{3}{4} E_{3z^2-r^2, 3z^2-r^2}(acb) - \frac{3^{\frac{1}{2}}}{4} E_{3z^2-r^2, x^2-y^2}(acb) \\
 &\quad + \frac{1}{4} E_{x^2-y^2, x^2-y^2}(acb)
 \end{aligned}$$

where we have related an integral between  $d$  functions transforming like  $x^2 - y^2$  at a point  $(abc)$  to a group of integrals at a point  $(acb)$ . By relations of this type we can express all integrals arising from a given set of neighbors in terms of the integrals arising from any one member of that set. This means that the sum over direct lattice vectors can be restricted to include only one vector from each set of neighbors, which



reduces the sum by a factor of 48 for a general set of neighbors.

For special sets of neighbors, by which is meant a set of neighbors containing less than 48 members, there exist some operations in the group which transform a vector of that set into itself. That means that not all the integrals at that point are independent; i.e., there exist relations of the type above for integrals on the same point. For maximum efficiency we could discard all but one member of each set of dependent integrals. This would substantially reduce the number of parameters involved. For example, in summing contributions from the first four sets of neighbors for functions of d symmetry, only 27 integrals need be included, out of a total of 62.

For overlap integrals the number is even fewer, since they have the symmetry properties of two-center integrals as expressed by Slater and Koster (1954). It is convenient to choose the overlap integrals in the same manner as that used for the three-center integrals, however. This allows the use of the same set of matrix elements to construct both energy and overlap matrices.

Use of group theory to obtain relations between integrals was of the utmost importance when programs to compute the three-center integrals were being tested. At that time all integrals were calculated, and the relations were used to detect errors in programming.

The band structure diagrams such as Figure 1 were drawn using group theory. These curves were plotted by connecting energy values whose eigenvectors transformed according to the same irreducible representation, subject to the condition that curves representing the same representations do not cross. Therefore it was necessary to know the basis functions for the various irreducible representations. Some of

these were obtained from Callaway (1964), and the remainder were computed using the character tables given by Slater (1965). At points where the symmetry changes, such as X or W, the bands were drawn with the assistance of compatibility relations found in both of the above references.

The calculation included as basis functions the wave functions for the 1s, 2s, 3s, 4s, 2p, 3p, 3d states of vanadium and the 1s, 2s, 2p states of oxygen. The eigenvalue problem therefore required the diagonalization of a 20 by 20 matrix. At symmetry points the order of the matrix could have been reduced using the orthogonality property of functions transforming according to different irreducible representations. At the point  $\Gamma$ , for example, we had the two dimensional representation  $\Gamma_{12}$ , the three dimensional representation  $\Gamma'_{25}$ , the one dimensional representation  $\Gamma_1$  which occurred six times, and the three dimensional representation  $\Gamma_{15}$  which occurred three times. Therefore the matrix could have been block diagonalized into the direct sum of a two by two, three by three, six by six, and nine by nine matrix. Even these would have been further reducible, since functions transforming according to different rows of the same irreducible representation are also orthogonal. Similar results hold for other symmetry points.

However, a 20 X 20 matrix is not difficult to diagonalize, so this procedure was unnecessary. Herein lies one of the great advantages of the tight binding method over other methods. Once the necessary integrals have been computed, it is as easy to find the energies at a general point in the Brillouin zone as at a symmetry point. In other methods, such as the orthogonalized plane wave (O P W) or augmented plane wave (A P W), the secular determinant may become so large that

solutions are possible only at symmetry points where block diagonalization is possible.

We now turn our attention to the problem of computing the three-center integrals, which formed the major part of the task. Since the crystal potential was constructed from a superposition of atomic potentials, the site at which it was centered was arbitrary provided it had the same symmetry as the origin, a vanadium site. Therefore C, the lattice site of the potential, was taken to be the origin. The site B was also taken as the origin, and the direct lattice sum was performing using the point A, where A, B, and C are defined as in equation (4) of Chapter II.

Expanding the potential in a Fourier series made it necessary to sum the series to convergence in the reciprocal lattice, which was done before performing the direct lattice sum. This meant summing over every reciprocal lattice vector at each magnitude of  $\underline{K}_n$ . The angular dependence of the sum was carried in the factors  $\cos(\underline{K}_n \cdot \underline{A})$ ,  $\sin(\underline{K}_n \cdot \underline{A})$ , and multiples of  $K_x, K_y$ , and  $K_z$  times these, where  $K_x, K_y$ , and  $K_z$  are the components of  $\underline{K}_n$ . It was possible to eliminate all terms in the sum arising from vectors having one or more negative components, by the method to be described now. Expand the cosine term as

$$\begin{aligned} \cos(\underline{K}_n \cdot \underline{A}) &= \cos K_x A_x \cos K_y A_y \cos K_z A_z - \sin K_x A_x \sin K_y A_y \\ &\quad \cos K_z A_z - \sin K_x A_x \cos K_y A_y \sin K_z A_z - \cos K_x A_x \\ &\quad \sin K_y A_y \sin K_z A_z. \end{aligned}$$

Hold the magnitudes of  $K_x, K_y$ , and  $K_z$  fixed and sum over their positive

and negative values. There are eight terms in this series, assuming that all three components are non-zero. The result is

$$\sum \cos(K_x A_x + K_y A_y + K_z A_z) = 8 \cos K_x A_x \cos K_y A_y \cos K_z A_z$$

where the sum is done as described. Similar results hold for  $K_x \sin(K_n \cdot A)$  and for terms such as  $K_x K_y \cos(K_n \cdot A)$ ,  $K_x^2 K_y \sin(K_n \cdot A)$ ,  $K_x K_y K_z \sin(K_n \cdot A)$ , etc. If one of the components  $K_x$ ,  $K_y$ ,  $K_z$  is 0, the factor 8 becomes 4. If two of them are 0 it becomes 2, and for the origin it is 1. These results allowed us to sum only over reciprocal lattice vectors having all components non-negative, which greatly reduced the sum.

For the special case of integrals where both functions had s symmetry, the sum was simplified even more. These functions have no angular dependence, so that it was possible to use only one vector from each set of neighbors in the reciprocal lattice. This could also have been done for the other integrals, but the gain in computer time would have been very little, since more integrals would have had to be computed for each term in the series.

The Fourier series was summed until each integral had converged to three significant figures. The Fourier series for the potential converged very slowly; a sum over 30,000 positive vectors in the reciprocal lattice yielded two place convergence at best and for many values of  $r$  it did not converge to even one place. The integrals contained a factor of  $\exp(-Kr)$  divided by  $K^n$ , where  $n$  varied from one integral to another, which made the Fourier series for the integrals converge much more rapidly. For all integrals involving p and d functions, the sum

was carried through the first 92 positive vectors, which covered 28 sets of neighbors. For integrals involving two s functions, the sum was carried out through 53 sets of neighbors. Even this did not produce convergence in the low lying s states, but functions representing these states had almost no overlap, so that the final result was unaffected by this lack of convergence. Table I contains a list of the first 53 reciprocal lattice vectors and the corresponding Fourier coefficients of potential and exchange for the assumed configuration  $3d^4 4s^1$ . The first 28 of these are the values over which p and d integrals were summed, although all positive permutations of  $K_x, K_y, K_z$  were used.

For the case in which both functions were located on the same site, the factor  $\exp(-k \cdot r)$  disappeared, and convergence was much slower. For d - d integrals, for example, it was necessary to sum 4000 positive vectors to obtain three place convergence in the potential, although exchange converged faster than this. For s - s integrals, 15,000 terms were included in the series. This yielded convergence only for the  $4s - 4s$  integral. The sum was extended to 30,000 terms in one case, but the integrals of core functions still did not converge.

Convergence for these central cell sums was checked not only by comparing results of successive partial sums but by comparing the result to the Hartree - Fock atomic energies. For the core states the two numbers should be very close. Since this was not the case for most integrals, the Hartree - Fock energy was used instead of the Fourier series sum for all central cell integrals except  $4s, 3d$ , and  $3p$  vanadium and  $2p$  and  $2s$  oxygen. In these latter sums the exchange potential played a major role, indicating that the Hartree - Fock energy would have been a poor approximation, since it does not permit the energy to

vary as  $\lambda$  changes. Table II contains the values used for various central cell integrals. For those integrals where the Hartree - Fock energy was not used, contributions from potential plus kinetic energy and exchange are listed separately. The exchange parameter for these results is equal to 1.

The sum over direct lattice vectors was carried out until each matrix element had converged to three significant figures. In some cases the convergence was better than this. The  $4s - 4s$  matrix element required 32 terms in the direct lattice sum before convergence was obtained. All matrix elements were summed through at least eighth neighbor contributions, although this was unnecessary for the core states. The programs for calculating the three-center integrals were set up in such a way that all integrals involving functions of the same symmetry from the same atom were computed at the same time. For example, the same program which computed  $3p - 3p$  integrals for two vanadium  $p$  functions also computed  $3p - 2p$  and  $2p - 2p$  integrals. Therefore the core states were carried through the same number of neighbors as were the higher functions of the same symmetry.

In order to give some idea of the extension of some of the more important wave functions, values for potential plus kinetic energy, exchange, and overlap integrals for some of the functions are given. Table III lists these values for the independent parameters for the  $3d - 3d$  integrals through 4 neighbors. This was enough to give convergence in the eigenvalues, although not in the matrix elements. This table does not include the two central cell parameters, which were given in Table II. Table IV gives the independent parameters for the  $2p - 2p$  integrals of oxygen through fourth neighbors, once again omitting the

central cell result. Table V gives potential plus kinetic energy, exchange and overlap for  $4s - 4s$  integrals through 32 neighbors.

In doing integrals where both functions were centered on the oxygen sublattice, a change of variable was made to place the oxygen atom at the origin. This moved the potential to the site  $(a/2, 0, 0)$ . Therefore oxygen-oxygen integrals appear to have the same coordinates as vanadium-vanadium integrals. This greatly simplified the computation of matrix elements. A similar change of variable was made so that oxygen-vanadium and vanadium-oxygen integrals use the same set of coordinates.

In the following chapter the results of the calculations described here will be presented. We now proceed to a discussion of those results.

## CHAPTER IV

In this chapter the results of the energy band calculation will be given. The effects of including core states will be examined, and an attempt will be made to estimate the configuration of a self-consistent potential.

Wave functions used in this calculation were generated by Clementi (1965) for the atomic configuration  $3d^3 4s^2$  for vanadium and  $2p^4$  for oxygen. We were interested in the configuration  $3d^4 4s^1 2p^4$  as a starting point for finding the energy band structure in the solid, so the occupation numbers were changed accordingly. Using these wave functions, we proceeded by means already described to calculate the energy band structure shown graphically in Figure 1. The notation for the irreducible representations is that of Bouckaert, Smoluchowski, and Wigner (1936).

Only the bands within 0.5 Rydbergs of the Fermi energy are depicted. These arise primarily from the 3d levels of vanadium. For example, at the maximum of bands  $\Delta_1$  and  $\Delta_5$ , the band  $\Delta_1$  is about 30 per cent s and 10 per cent p, the rest being 3d, and the  $\Delta_5$  band is about 15 per cent p and 85 per cent 3d. A band arising from the 4s level of vanadium appears in the vicinity of  $\Gamma$ , but it quickly rises above the 3d bands. A set of bands which are due to the 2p levels of oxygen lie approximately 1.2 Rydbergs below the Fermi energy and have a width of about 2.6 ev.

Maximum width for the d bands is about 7.3 ev, with the minimum occurring at  $X_3$  and the maximum occurring halfway along the upper  $Q_1$  band between W and L. The energy bands along the  $\Delta$  axis lie substantially lower than those along other symmetry axes, with the result that



the Fermi level lies above both  $\Gamma_{12}$  and  $\Gamma'_{25}$ , although vanadium has only three 3d electrons per cation in this band picture. One of the 3d electrons and the 4s electron have been used to fill the p bands arising from oxygen.

The energy bands which have been described were computed for  $\lambda = 1$ ; that is, using the full Slater exchange. A band structure was also calculated using  $\lambda = 0.85$ , a value which was suggested by Callaway and Zhang (1969) for nickel. This led to a band structure which had a positive Fermi energy, so we will confine our discussion to the energy bands computed using the full Slater exchange, at least for this configuration.

One of the things of interest about this calculation is that it is the first tight binding calculation on an element this high in the periodic table to include all the core electrons explicitly. Therefore it will be worthwhile to consider the effect of including the core states.

The band structure found from a calculation using the d bands alone is pictured along the  $\Delta$  axis in Figure 2(a). They are the typical transition metal d bands, closely resembling those found for nickel in an earlier calculation (Tyler et al., 1969). Total band width is about 9.7 eV, which is rather large for a transition metal. With no core states, the  $\Delta_1$  band sweeps far down, and the  $\Delta_5$  band lies almost entirely below the  $\Delta_2$  band.

When we included the 2p states from oxygen and the 4s state from vanadium, the band structure shown in Figure 2(b) was found. The  $\Delta_1$  band has hybridized with the 4s band, and the p bands lie below the d bands in about the same position and with the same shape as in the full calculation. It is clear that at least the s states of vanadium must be added to prevent the 4s band from falling back into the core as

it approaches X. The approximation of treating the Bloch functions formed from core states as being orthogonal to those formed from conduction states is not justified here.

In Figure 3 the band structure is given in which all core states except the s states of oxygen were added. The 4s band now begins at -1.5 Rydbergs and passes up through the d bands, causing the s-d hybridization shown. It still contains a downward curve in the vicinity of X, but the  $X_1$  point of the 4s band lies well above the Fermi energy, which is very slightly above  $\Gamma'_{25}$ . The  $\Delta_5$  band has been pushed above the  $\Delta_2$  band for most of the distance between  $\Gamma$  and X.

Comparison of this figure with Figure 1 shows the large effect produced by including the oxygen s states. Their inclusion pushed the 4s band completely above the Fermi energy, so that it plays no role in the ground state conduction picture. The  $\Delta_1$  band of the d bands now lies completely above the  $\Delta_5$  band and the Fermi energy lies above both  $\Gamma_{12}$  and  $\Gamma'_{25}$ .

From the preceding discussion it is clear that in order to obtain an accurate picture of the band structure, it is necessary to include the core states. The calculation was greatly helped in this regard by the fact that Clementi's wave functions use a single set of Slater orbitals to represent all s states from one atom and another single set to represent all p states from one atom, changing only the coefficients in going from 2p to 3p and likewise for s states. This meant that it was possible to obtain all vanadium s-s integrals, for example, by doing a single set of integrals and then combining them by using the various sets of coefficients.

In Figure 4 we have a graphical representation of the density of states near the Fermi energy for the band structure of Figure 1. This exhibits a peak above the Fermi energy arising from the flat  $Z_2$  band of Figure 1, which is a characteristic of the face-centered cubic structure. The density of states curve is based on a sampling of 505 points in  $1/48$  of the Brillouin zone, as shown in Figure 5. The energy grid used was 0.008 Rydbergs.

A second band structure was calculated using the atomic configuration  $3d^3 4s^2 2p^4$  as a starting point. This is the atomic configuration actually described by the wave functions. Results of this calculation along the  $\Delta$  axis are shown in Figure 6(a) for the full Slater exchange.

The Fermi energy has decreased from -0.283 Rydbergs in the  $3d^4 4s^1 2p^4$  configuration to -1.15 Rydbergs in this configuration. The d bands were shifted downward more than the 4s bands, and the 2p bands of oxygen lay about 0.8 Rydbergs below the Fermi energy as compared to 1.2 Rydbergs before. Final configuration for the solid in both calculations was  $3d^3 4s^0 2p^6$ . Ordering of bands was the same in all cases for the two calculations.

Band structures were also run for this configuration using other values of  $\lambda$ . For  $\lambda = 0.85$ , a set of bands very similar in appearance but having Fermi energy -0.77 Rydbergs was obtained, as shown in Figure 6(b). Results for the Kohn-Sham (1965) value of  $2/3$  for  $\lambda$  are shown in Figure 6(c). Fermi energy has risen to -0.304 Rydbergs. In this latter case, the point  $W_2'$  was above the point  $W_3$  (see Figure 1). Otherwise no changes in ordering were found for either  $\lambda = 0.85$  or  $\lambda = 2/3$ .

Comparison of the results for the different values of  $\lambda$  gives a picture of the effect of varying the exchange potential. With less

exchange the energy of the system is higher, of course. Since exchange affects all the d bands by approximately the same amount, changing the value of  $\lambda$  shifts all the d bands about the same, with the result that the overall shape and relative positions of bands tend to remain unchanged. This is evident from the three graphs in Figure 6.

A decrease in the exchange potential means that the crystal field will have a more powerful role in the structure of the bands. This is evidenced by the fact that the band splitting at the  $\Gamma$  point increases with decreasing  $\lambda$ . Band width also increases with decreasing  $\lambda$ . Total band width for the structure of Figure 6(a) is 7.04 ev, for that of Figure 6(b) is 8.4 ev, and for that of Figure 6(c) is 10.0 ev. For the configuration  $3d^4 4s^1 2p^4$ , the band width was 7.3 ev, as mentioned before.

In all respects other than those mentioned above, the bands obtained for the configuration  $3d^3 4s^2 2p^4$  were very similar to those obtained for the configuration  $3d^4 4s^1 2p^4$ . Therefore they will not be discussed further. Instead we will turn our attention to the problem of self-consistency.

The fact that the crystal potential was formed from a superposition of spherically symmetric potentials might affect the accuracy of the crystal field splitting

$$\Delta E = E_{x^2-y^2, x^2-y^2}(000) - E_{xy, xy}(000) \quad .$$

Since there have been no other energy band calculations on VO, it is impossible to make any comparisons here. To best insure that the potential does yield the correct splitting, the band calculation should be done by a self-consistent technique. A procedure for this is now being

developed but is not yet available.

However, it would be useful to obtain at least an estimate of the self-consistent configuration. Both band structures computed at this point have seemed to indicate that the crystal contains doubly ionized atoms. Therefore it was decided to calculate another band structure using the assumed atomic configuration  $3d^3 4s^0 2p^6$ . The same wave functions were used; only the occupation numbers were changed.

The fact that the individual atoms were no longer neutral caused difficulties with the Fourier series convergence. The various central cell integrals were summed to convergence, taking an average of fifty per cent more Fourier coefficients than were used in the first two calculations. Other integrals, however, were summed using the same number of terms as before, with the result that convergence was probably not as good as before. Since the central cell results dominate all other contributions except in the case of the  $4s$  band, this lack of convergence should have little effect. The  $4s$  band lies well above the other bands, so that slight errors in it should not be felt in the vicinity of the Fermi energy.

The band structure obtained for this configuration is shown along the  $\Delta$  and  $Z$  axes in Figure 7. This structure is based on the full Slater exchange. The  $p$  bands have been raised above the  $d$  bands and are completely empty. Electronic configuration for these bands is  $3d^9 4s^0 2p^0$ . This result seems to indicate that the crystal cannot be treated as being doubly ionized.

Before going on to discuss the problem of self-consistency further, it would be appropriate to digress to make a comparison of this last band calculation to one done by Yamashita (1963) for titanium monoxide ( $TiO$ ).

The TiO calculation also used wave functions based on a doubly ionized configuration for the atom. His method was the tight binding approximation as modified by Stern (1959), in which the radial equation is solved numerically in two cells, one centered on titanium and the other on oxygen. His calculation included no exchange, and core functions were not explicitly included. Also, interactions among s, d, and p electrons were not included.

He found the 2p bands lying below the d bands with a band width of 10.7 ev. He calculated values only at the points  $\Gamma$ , X, and L, so his band width result is the maximum spread at those points. These results are in opposition to the results of our calculation. In our case the width of the p bands is difficult to determine since there is a great deal of p-d interaction. However, it is about 3 ev. The difference in results for the two cases is probably due to the differences in method which were mentioned. There is also the fact that in Yamashita's calculations the 2p atomic orbital used was probably inaccurate. He was mainly interested in d bands, and so made little effort to obtain accurate results for the p bands.

We will now return to the problem of the consistency of the calculation. Our goal here is to obtain an estimate of the self-consistent configuration. This is done by the procedure described below.

All three calculations have given an empty 4s band as the final result. Therefore the nine electrons must be spread between the 3d and 2p bands. Since the p levels will not be fully occupied, they must lie very close to the d levels in order to have some portion above the Fermi energy. Therefore we seek a configuration in which the  $\Gamma_{15}$  point lies very close to the  $\Gamma'_{25}$  point. Evidently some small adjustment to the

$3d^3 4s^0 2p^6$  configuration should yield the correct answer.

In order to compute this adjustment, we need to know the effect on relative positions of bands of moving electrons from the p band to the d band. This can be found by examining the s-d, p-d, and s-p energy gaps at  $\Gamma$  for each of the three configurations used so far. One can determine the effect on relative band positions of moving two s electrons to the p bands by considering the difference between energy gaps for the configurations  $3d^3 4s^2 2p^4$  and  $3d^3 4s^0 2p^6$ . Assuming that the effect of moving one electron from 4s to 2p is half this, it was found that changing one electron from s to p raised the 4s band relative to the 3d band by 0.186 Rydbergs, raised the 2p band relative to the 3d band by 0.52 Rydbergs, and raised the 2p band relative to the 4s band by 0.375 Rydbergs. Comparison of the energy gaps for the configurations  $3d^4 4s^1 2p^4$  and  $3d^3 4s^0 2p^6$  allowed us to determine the effect of removing an electron from the 3d band and placing it in the 2p band, if it is assumed that the effect of transferring a 4s and a 3d electron simultaneously to 2p is equal to the sum of the individual effects. From this we found that transferring an electron from 3d to 2p raises the 4s band relative to the 3d band by 0.81 Rydbergs, raises the 2p band relative to the 3d band by 1.02 Rydbergs, and lowers the 4s band relative to the 2p band by 0.285 Rydbergs.

Since the energy gap between  $\Gamma_{15}$  and  $\Gamma'_{25}$  for the configuration  $3d^3 4s^0 2p^6$  was 0.325 Rydbergs, we concluded from the above that a transfer of about 0.3 electrons from the 2p band to the 3d band would bring p and d bands at  $\Gamma$  into close proximity, which is what we desired. The conclusion was therefore that an approximately self-consistent potential for VO is  $3d^{3.3} 4s^0 2p^{5.7}$ .

It should be pointed out here that self-consistency is determined not by the configuration but by the crystalline charge density. In a calculation on TiO by Ern and Switendick (1965) for example, the initial atomic configuration assumed was  $3d^2 4s^1 2p^5$ , while the apparent equivalent configuration for the band electrons was  $3d^2 4s^0 2p^6$ . However, the charge densities, as plotted in the spheres centered on titanium and oxygen, were very nearly the same. Thus comparing configurations may be misleading. An accurate determination of the degree of self-consistency can be made only by comparing charge densities obtained from successive band calculations computed in some sort of iterative procedure.



## CHAPTER V

In this final chapter, we will briefly discuss the possibility of a metal-insulator transition in VO, compare the results of this calculation to those of a calculation on TiO, and suggest certain improvements which could be made in future investigations.

The possibility of a metal-insulator transition in VO due to the onset of antiferromagnetic order was presented by Morin in 1959, when he first reported a transition in VO, and later discussed by Adler (1968) after the work of Warren et al. (1967), although Adler assumed a simultaneous crystalline distortion. One of the requirements for a transition caused by antiferromagnetic splitting of the bands is that there be little or no overlap between  $\Gamma_{12}$  bands and  $\Gamma'_{25}$  bands. Otherwise no gap can be introduced. From Figure 1 it is evident that this criterion is not satisfied in VO. It is more difficult to assess the possibility of a transition caused by crystalline distortion, since band splitting and even the ordering of bands could not be determined without knowing the low temperature structure of VO. However, due to the relatively large band width of the bands arising from  $\Gamma'_{25}$ , a transition of this type appears unlikely. Further discussion of this topic must wait until the question of the existence of the transition has been settled.

An APW calculation of the energy bands of TiO was performed by Ern and Switendick (1965), who assumed the atomic configuration  $3d^2 4s^1 2p^5$ , which corresponds to singly ionized atoms. They found the 2p bands from oxygen to lie 0.35 Rydbergs below the 3d bands at  $\Gamma$ . Based on the interpolation procedure described at the end of Chapter IV, the 2p bands of oxygen would be 0.2 Rydbergs below the 3d bands at  $\Gamma$

for a calculation on VO using singly ionized atoms. Total band width for TiO was found to be about 9 ev, compared to about 7 ev in VO. It is to be expected that the band width in TiO is larger than in VO, since TiO is a better metal and lies further to the left in the transition series.

The conduction in VO is due mainly to the 3d bands, based on the estimate of a self-consistent potential made in Chapter IV. According to that estimate, there would be a small pocket of holes somewhere on the Fermi surface due to the 2p bands of oxygen. This pocket should contain only 0.3 holes, so that the 3d bands will dominate the Fermi surface features. This can be determined definitely only by a self-consistent calculation.

The calculation can be improved in several ways. First, of course, it should be done self-consistently by a systematic process in which each calculation begins with the results of the previous one. This will be done in the near future. A second correction which would be useful is to include more conduction band states, such as a 4p from vanadium and 3s from oxygen. This would no doubt have the effect of repelling some of the higher 3d bands and thus narrowing the band width. Minimum band width possible is about 6.9 ev, since the points  $L_3$  and  $X_3$  are unaffected by s and p bands. Finally, if one wishes to make detailed studies of the Fermi surface, the effects of spin-orbit coupling should be included. There appears to be little value in this at the present time, however. Even if VO is a metal at the temperatures for which Fermi surface measurements are made, it is so highly defective that the mean free path of electrons is probably too short to obtain good results. Since very little is known about VO experimentally, we have not tried to

predict anything other than its gross properties from the band calculation.

We conclude this discussion of results by estimating the overall accuracy of this calculation. In carrying out the various sums and integrations involved in the calculation, we attempted to achieve at least three figure accuracy in each. Allowing for some loss of accuracy in carrying out the computations, it is estimated that the results quoted for the energy bands are accurate within 0.02 Rydbergs, which certainly would not cause any reordering of bands due to numerical errors. As mentioned before, the calculation for the ionic case was not as accurate as the others, although we believe any errors would be small.

Finally, we would like to mention a very important practical consideration in the use of this method to compute energy bands. Computer time on an IBM 360-65 for the first band structure was about 60 hours. Successive iterations require only four to five hours, since the basic integrals can be stored on tape and reused provided the Slater orbitals are not changed. Refinements in technique should permit both the initial time and time for the iterations to be reduced considerably in the future. For future reference we have included three-center programs used in this calculation as Appendix C. The programs are commented with explanations to facilitate usage, and are written in Fortran IV.

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TABLE I

<u>RECIPROCAL LATTICE VECTOR</u>	<u>FOURIER COEFFICIENT OF POTENTIAL</u>	<u>FOURIER COEFFICIENT OF EXCHANGE</u>
000	-1.7134 RYDBERGS	-0.93112 RYDBERGS
111	-0.35346	-0.061622
200	-0.83970	-0.37232
220	-0.60412	-0.10930
131	-0.19506	-0.016291
222	-0.47470	-0.0068919
400	-0.39046	0.016495
331	-0.14030	-0.0045957
240	-0.33097	0.0074661
422	-0.28676	-0.0094525
333	-0.10902	-0.0031780
511	-0.10902	-0.0031780
440	-0.22577	-0.029806
531	-0.088932	-0.0036546
600	-0.20397	-0.029922
442	-0.20397	-0.029922
620	-0.18602	-0.025360
335	-0.075128	-0.0037931
622	-0.17100	-0.018407
444	-0.15827	-0.011053
551	-0.065154	-0.0034620
711	-0.065154	-0.0034620
640	-0.14736	-0.0047153

TABLE I (Cont.)

264	-0.13792	-0.00020310
137	-0.051175	-0.0029812
553	-0.057655	-0.0029812
800	-0.12239	0.0026204
733	-0.051832	-0.0025842
820	-0.11594	0.0014205
644	-0.11594	0.0014205
228	-0.11018	-0.00089478
660	-0.11018	-0.00089478
715	-0.047186	-0.0023371
555	-0.047186	-0.0023371
662	-0.10500	-0.0037961
840	-0.10032	-0.0067985
119	-0.043391	-0.0021985
537	-0.043391	-0.0021985
284	-0.096075	-0.0095067
664	-0.092201	-0.011638
391	-0.040229	-0.0020945
844	-0.085388	-0.013605
933	-0.037549	-0.0019663
755	-0.037549	-0.0019663
771	-0.037549	-0.0019663
860	-0.082376	-0.013412
0010	-0.082376	-0.013412
862	-0.079585	-0.012541

TABLE I (Cont.)

0210	-0.079585	-0.012541
159	-0.035244	-0.0017875
773	-0.035244	-0.0017875
2210	-0.076993	-0.011138
666	-0.076993	-0.011138



TABLE II

	HARTREE-FOCK ATOMIC ENERGY ( <u>RYDBERGS</u> )	POTENTIAL + K.E. ( <u>RYDBERGS</u> )	EXCHANGE ( <u>RYDBERGS</u> )
1s-1s(V)	-403.004		
2s-2s(V)	- 47.7485		
3s-3s(V)	- 6.36626		
2p <sub>x</sub> -2p <sub>x</sub> (V)	- 40.0441		
1s-1s(0)	- 41.3864		
3p <sub>x</sub> -3p <sub>x</sub> (V)		-0.66111	-3.0408
3d <sub>xy</sub> -3d <sub>xy</sub> (V)		1.9921	-2.38687
3d <sub>x<sup>2</sup>-y<sup>2</sup></sub> -3d <sub>x<sup>2</sup>-y<sup>2</sup></sub> (V)		1.9340	-2.4687
4s-4s(V)		-0.73614	-0.85674
2s-2s(0)		-0.45890	-1.940
2p <sub>x</sub> -2p <sub>x</sub> (0)		-0.3096	

TABLE III

<u>INTEGRAL</u>	<u>NEIGHBOR</u>	<u>POTENTIAL + K.E.</u> <u>(RYDBERGS)</u>	<u>EXCHANGE</u> <u>(RYDBERGS)</u>	<u>OVERLAP</u>
xy,xy	110	-0.013533	-0.018263	0.028805
xz,xz	110	0.014814	0.0055322	-0.11041
yz,xz	110	0.022171	0.010374	-0.017641
$x^2-y^2, x^2-y^2$	110	0.054771	0.032921	-0.28683
$3z^2-r^2, 3z^2-r^2$	110	-0.014825	-0.14572	0.14001
xy, $2z^2-r^2$	110	0.0050073	0.010396	-0.12820
xy,xy	002	0.0008587	-0.0004164	0.0007628
xz,xz	002	0.056684	0.0035288	-0.0047653
$x^2-y^2, x^2-y^2$	002	-0.0016500	-0.0008252	0.76283
$3z^2-r^2, 3z^2-r^2$	002	-0.016694	-0.011154	0.0086867
xy,xy	112	-0.0001643	-0.0000274	0.00007008
xz,xz	112	-0.0006314	-0.0002646	0.0004947
xy,xz	112	-0.0003091	-0.0001569	0.0002831
yz,xz	112	-0.0018143	-0.0009518	0.0012314
$x^2-y^2, x^2-y^2$	112	0.0005365	0.0003374	-0.0002932
$3z^2-r^2, 3z^2-r^2$	112	-0.0000445	0.0001120	-0.0000904
xy, $3z^2-r^2$	112	-0.0013379	-0.0008810	0.0008853
yz, $3z^2-r^2$	112	-0.0016076	-0.0009911	0.0010024
xz, $3z^2-r^2$	112	0.0007762	0.0004628	-0.0004435
xy,xy	220	-0.0015149	-0.0006727	0.0007453
xz,xz	220	0.0001228	0.0001079	-0.0001577
yz,xz	220	0.0001935	0.0001388	-0.0001958
$x^2-y^2, x^2-y^2$	220	0.0004439	0.0003498	-0.0003535

TABLE III (Cont.)

$3z^2-r^2, 3z^2-r^2$	220	-0.0006258	-0.0003070	0.0002738
$xy, 3z^2-r^2$	220	0.0009194	0.0004428	-0.0004083

TABLE IV

<u>INTEGRAL</u>	<u>NEIGHBOR</u>	<u>POTENTIAL + K.E.</u> <u>(RYDBERGS)</u>	<u>EXCHANGE</u> <u>(RYDBERGS)</u>	<u>OVERLAP</u>
$2p_x \ 2p_x$	110	0.020731	0.015933	-0.020321
$2p_z \ 2p_z$	110	-0.017555	-0.0092185	0.012889
$2p_y \ 2p_x$	110	0.051565	0.031263	-0.033211
$2p_x \ 2p_x$	002	-0.0025728	-0.0012367	0.0015517
$2p_z \ 2p_z$	002	0.019164	0.011193	-0.0095615
$2p_x \ 2p_x$	112	-0.00016994	0.00008054	-0.0001403
$2p_z \ 2p_z$	112	0.0016397	0.0013462	-0.0015002
$2p_y \ 2p_x$	112	0.00046009	0.0003718	-0.0004533
$2p_y \ 2p_z$	112	0.0010286	0.0007598	-0.0009066
$2p_x \ 2p_x$	220	0.00030285	0.0002958	-0.0003275
$2p_z \ 2p_z$	220	-0.00015390	-0.00006415	0.00008174
$2p_y \ 2p_x$	220	0.00051499	0.0003836	-0.0004093

TABLE V

<u>NEIGHBOR</u>	<u>POTENTIAL + K.E.</u> <u>(RYDBERGS)</u>	<u>EXCHANGE</u> <u>(RYDBERGS)</u>	<u>OVERLAP</u>
110	-0.52091	-0.33517	0.45595
002	-0.33492	-0.19396	0.23305
112	-0.19270	-0.10488	0.12211
220	-0.11210	-0.058653	0.066344
310	-0.064740	-0.033453	0.037266
222	-0.038270	-0.019375	0.021560
321	-0.023260	-0.011616	0.012807
004	-0.014292	-0.0070803	0.0077897
114	-0.0089598	-0.0044028	0.0048412
330	-0.0090094	-0.0044136	0.0048412
420	-0.0057236	-0.0028033	0.0030685
332	-0.0037135	-0.0018078	0.0019804
224	-0.0024475	-0.0011869	0.0012995
431	-0.0016399	-0.0007928	0.0008659
510	-0.0016381	-0.0007935	0.0008659
521	-0.0007623	-0.0003675	0.0004008
440	-0.0005298	-0.0002546	0.0002778
334	-0.0003713	-0.0001784	0.0001947
530	-0.0003717	-0.0001786	0.0001947
442	-0.0002635	-0.0001265	0.0001379
006	-0.0002632	-0.0001265	0.0001379
532	-0.0001888	-0.0000905	0.0000986
116	-0.0001884	-0.0000905	0.0000986

TABLE V (Cont.)

026	-0.0001362	-0.0000653	0.0000712
541	-0.0000993	-0.0000476	0.0000518
226	-0.0000727	-0.0000349	0.0000380
631	-0.0000538	-0.0000258	0.0000280
444	-0.0000400	-0.0000192	0.0000209
710	-0.0000300	-0.0000144	0.0000156
543	-0.0000300	-0.0000143	0.0000156
550	-0.0000300	-0.0000143	0.0000156
640	-0.0000226	-0.0000108	0.0000118

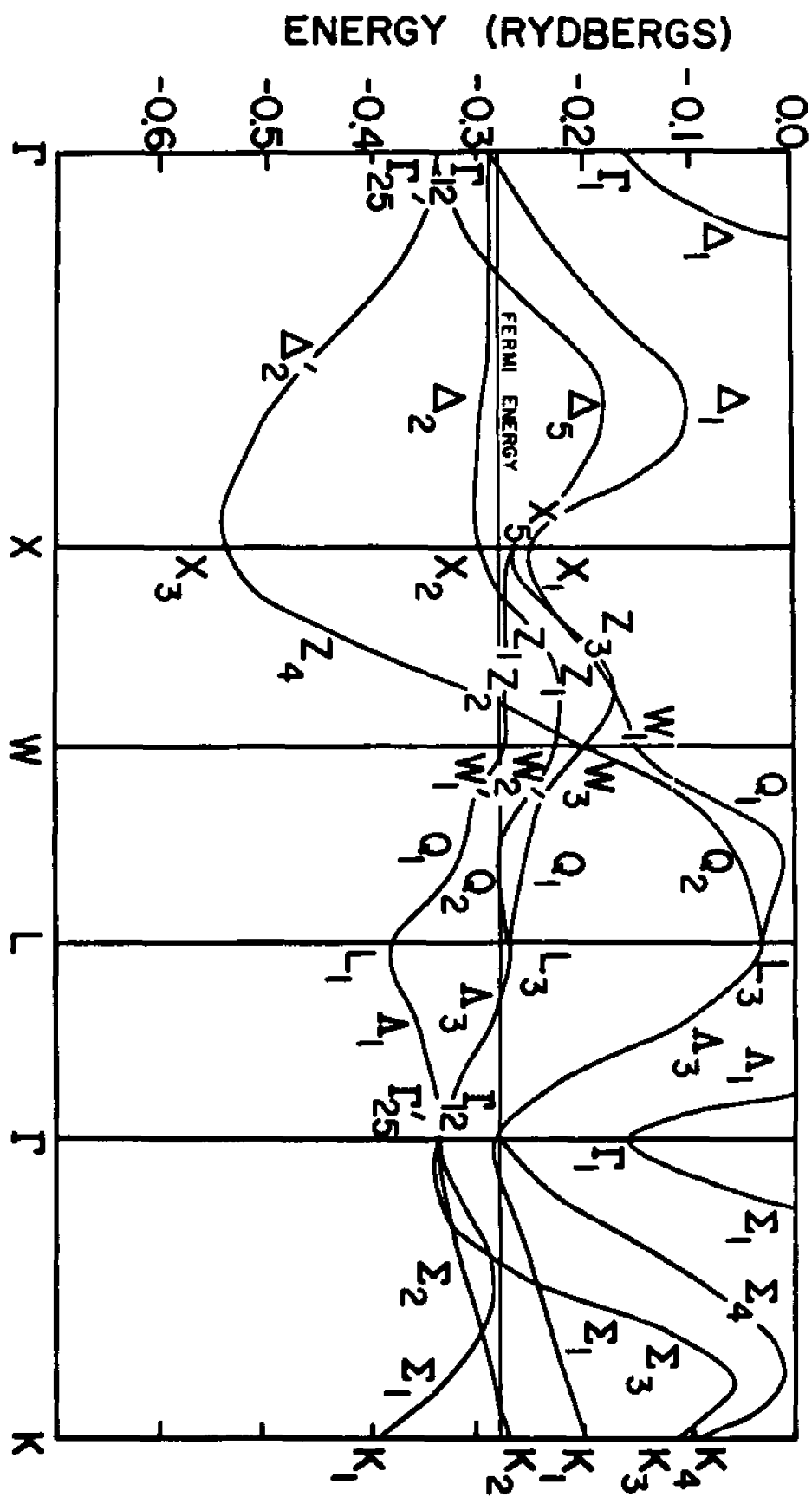


FIGURE 1

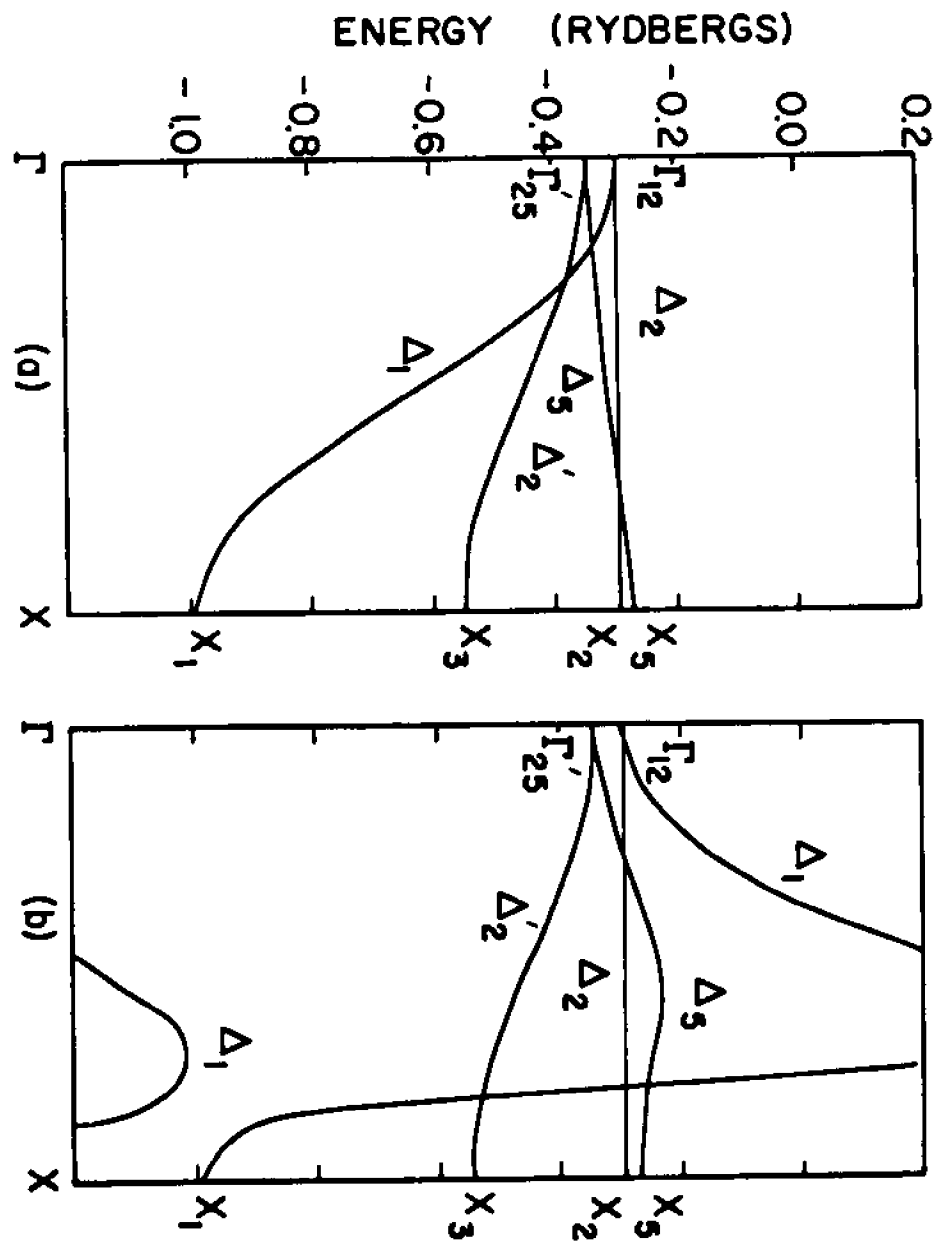


FIGURE 2



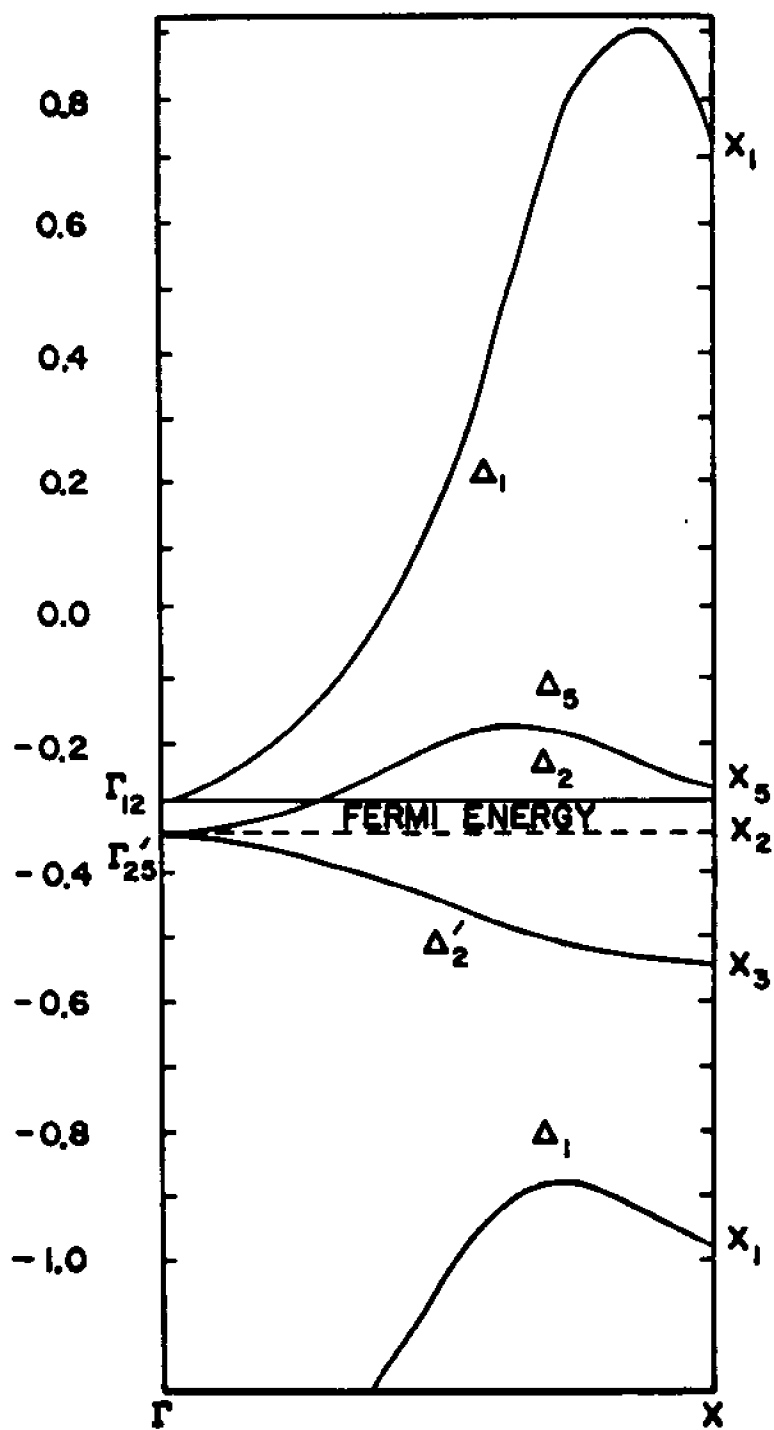


FIGURE 3

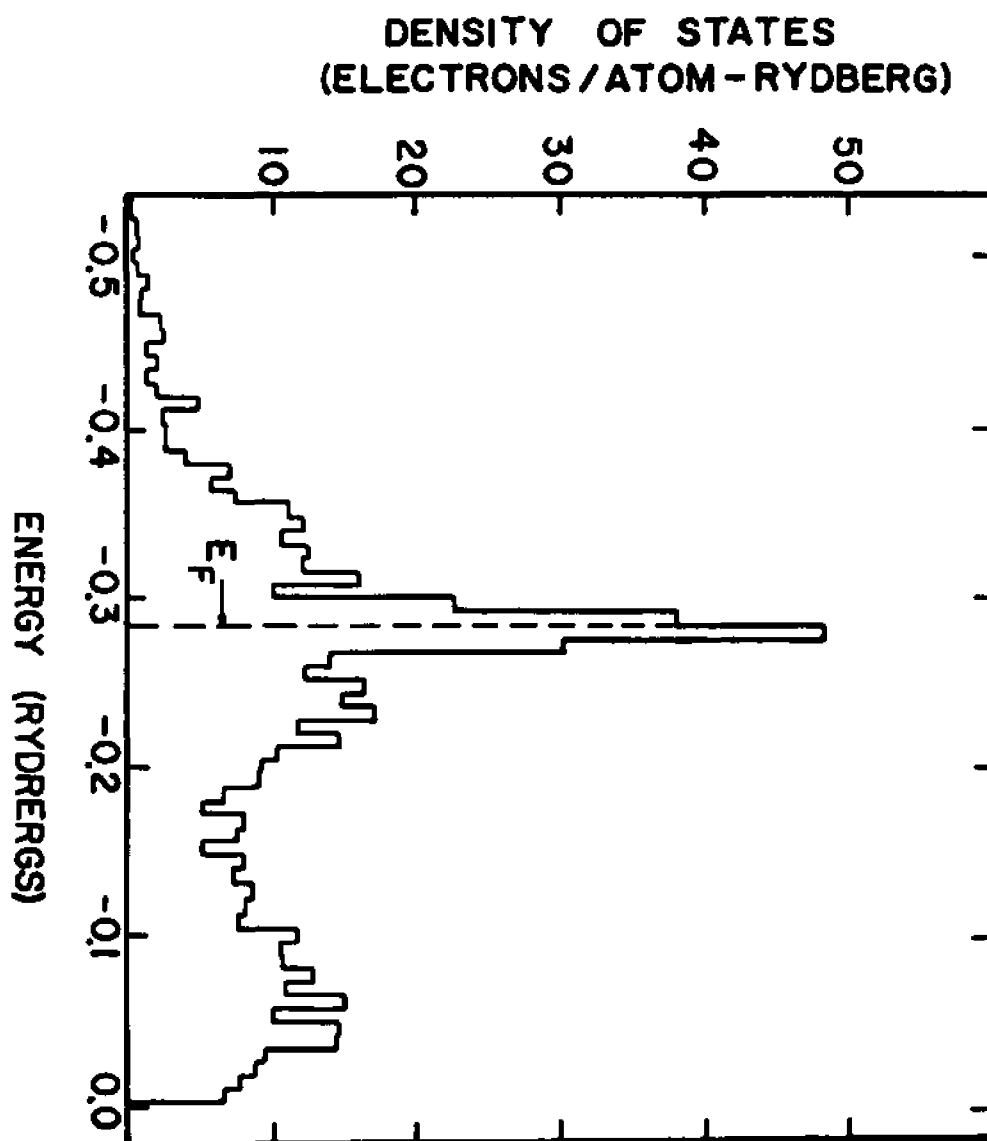


FIGURE 4



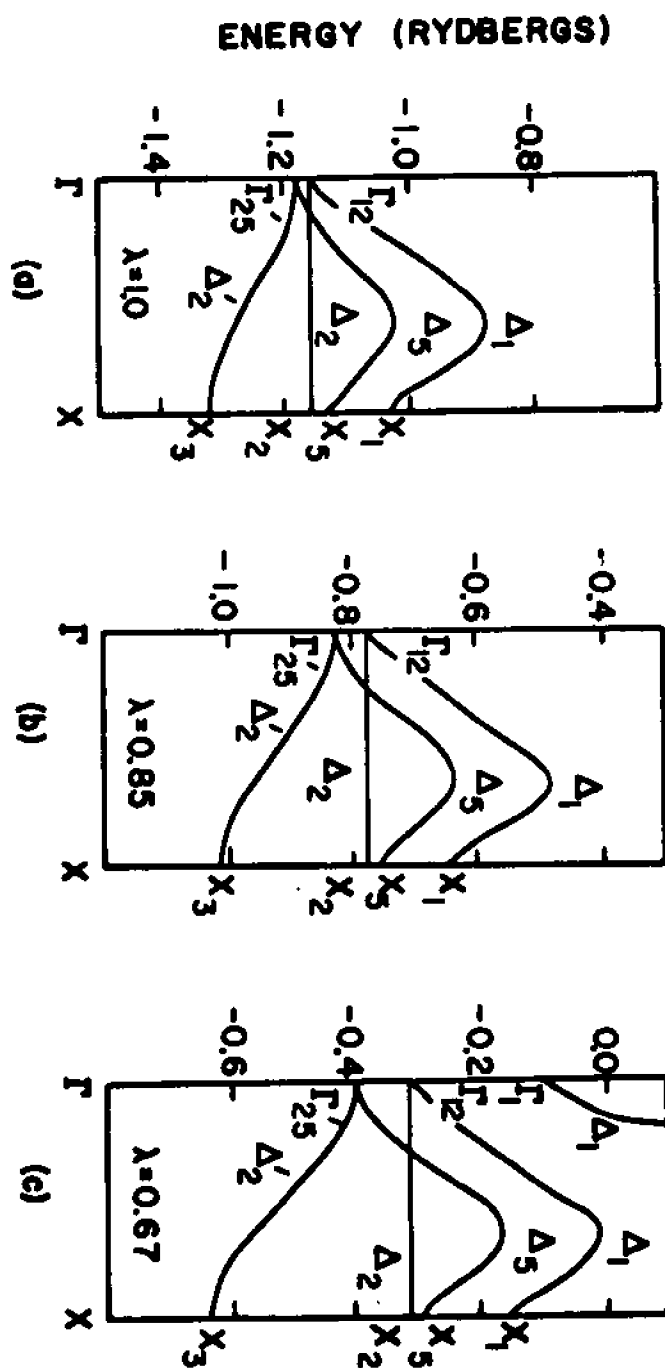


FIGURE 6

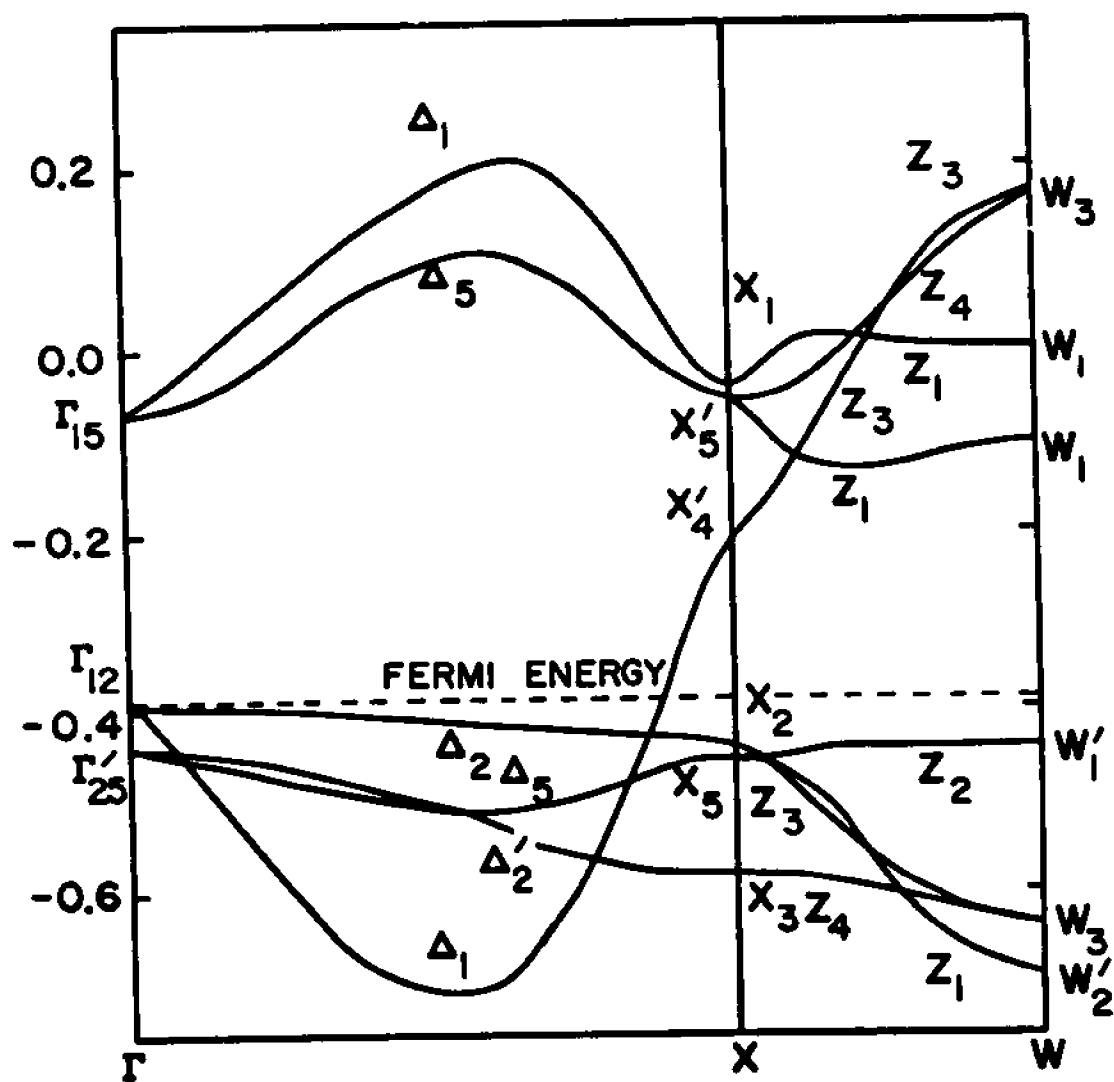


FIGURE 7

## APPENDIX A

The potential integrals involving any Slater orbitals can be obtained from the 1s-1s integral given in the second chapter by taking the proper partial derivatives. These derivatives become very tedious for higher orbitals if carried out step by step. However, a set of recursion relations for the derivatives was discovered empirically, and this greatly simplifies the work. These will be presented in this appendix, together with the integrals used in this calculation. In addition to the definitions given in Chapter II, the following will be useful:

$$A(0) = \frac{1}{(fg)^{3/2}} + \frac{1}{(fg)}$$

$$A(1) = f \left[ \frac{3}{(fg)^{5/2}} + \frac{3}{(fg)^2} + \frac{1}{(fg)^{3/2}} \right]$$

$$A(2) = f^2 \left[ \frac{15}{(fg)^{7/2}} + \frac{15}{(fg)^3} + \frac{6}{(fg)^{5/2}} + \frac{1}{(fg)^2} \right]$$

$$A(3) = f^3 \left[ \frac{105}{(fg)^{9/2}} + \frac{105}{(fg)^4} + \frac{45}{(fg)^{7/2}} + \frac{10}{(fg)^3} + \frac{1}{(fg)^{5/2}} \right]$$

$$A(4) = f^4 \left[ \frac{945}{(fg)^{11/2}} + \frac{945}{(fg)^5} + \frac{420}{(fg)^{9/2}} + \frac{105}{(fg)^4} + \frac{15}{(fg)^{7/2}} + \frac{1}{(fg)^3} \right]$$

$$A(5) = f^5 \left[ \frac{10395}{(fg)^{13/2}} + \frac{10395}{(fg)^6} + \frac{4725}{(fg)^{11/2}} + \frac{1260}{(fg)^5} + \frac{210}{(fg)^{9/2}} + \frac{21}{(fg)^4} + \frac{1}{(fg)^{7/2}} \right]$$

$$A(6) = r^6 \left[ \frac{135135}{(fg)^{15/2}} + \frac{135135}{(fg)^7} + \frac{62370}{(fg)^{13/2}} + \frac{17325}{(fg)^6} + \frac{3150}{(fg)^{11/2}} \right. \\ \left. + \frac{378}{(fg)^5} + \frac{28}{(fg)^{9/2}} + \frac{1}{(fg)^4} \right]$$

$$A(7) = r^7 \left[ \frac{2027025}{(fg)^{17/2}} + \frac{2027025}{(fg)^8} + \frac{945945}{(fg)^{15/2}} + \frac{270270}{(fg)^7} + \frac{51975}{(fg)^{13/2}} \right. \\ \left. + \frac{6930}{(fg)^6} + \frac{630}{(fg)^{11/2}} + \frac{36}{(fg)^5} + \frac{1}{(fg)^{9/2}} \right]$$

$$\cos = \cos \underline{\underline{K}}_v \cdot \underline{\underline{r}}_{CD}$$

$$\sin = \sin \underline{\underline{K}}_v \cdot \underline{\underline{r}}_{CD}$$

$$B_x - A_x = D_x \qquad (\underline{\underline{K}}_v)_x = K_x$$

$$B_y - A_y = D_y \qquad (\underline{\underline{K}}_v)_y = K_y$$

$$B_z - A_z = D_z \qquad (\underline{\underline{K}}_v)_z = K_z$$

$$C = 2\pi\alpha_1\alpha_2r_{AB}^3 \qquad h = u(1-u)$$

In terms of these quantities, the 1s-1s potential integral becomes

$$\langle 1s(A) | \cos \underline{\underline{K}}_v \cdot \underline{\underline{r}}_{CD} | 1s(B) \rangle = C \int_0^1 \cos \underline{\underline{K}}_v \cdot \underline{\underline{r}}_{CD} A(1) e^{-(fg)^{\frac{1}{2}}} du$$

Any potential integral involving Slater-type orbitals not higher than 4s can be written in the following manner:

$$\langle \Phi_1(A) | \cos K_{\nu} \cdot r_{CD} | \Phi_j(B) \rangle = \sum_{N=1}^7 C \int_0^1 F_N A(N) e^{-(fg)^{\frac{1}{2}}} du$$

where  $F_N$  is in general a function of  $D_x, D_y, D_z, K_x, K_y, \alpha_1, \alpha_2, K_z$ , and  $u$ .  $F_N$  will always contain either or both of the factors  $\cos K_{\nu} \cdot r_{CD}$ ,  $\sin K_{\nu} \cdot r_{CD}$ . Before giving the recursion relations for derivatives of this integral, we note the following results.

$$\frac{d}{dA_x} \cos K_{\nu} \cdot r_{CD} = -K_x u \sin K_{\nu} \cdot r_{CD} \quad (1)$$

$$\frac{d}{dB_x} \cos K_{\nu} \cdot r_{CD} = -K_x (1-u) \sin K_{\nu} \cdot r_{CD} \quad (2)$$

There are four derivatives of the integral in which we are interested. The results for these are now presented in equations (3) through (6).

$$\begin{aligned} \frac{d}{dA_x} \langle \Phi_1(A) | \cos K_{\nu} \cdot r_{CD} | \Phi_j(B) \rangle &= \\ \frac{d}{dA_x} \sum_{N=1}^7 C \int_0^1 F_N A(N) e^{-(fg)^{\frac{1}{2}}} du &= \\ \sum_{N=1}^7 C \int_0^1 [D_x h F_N A(N-1) + \frac{dF_N}{dA_x} A(N)] e^{-(fg)^{\frac{1}{2}}} du & \quad (3) \end{aligned}$$

$$\begin{aligned} \frac{d}{dB_x} \langle \Phi_1(A) | \cos K_{\nu} \cdot r_{CD} | \Phi_j(B) \rangle &= \\ \sum_{N=1}^7 C \int_0^1 [-D_x h F_N A(N-1) + \frac{dF_N}{dB_x} A(N)] e^{-(fg)^{\frac{1}{2}}} du & \quad (4) \end{aligned}$$

$$\begin{aligned} \frac{d}{d\alpha_1} \langle \Phi_1(A) | \cos K_{\nu} \cdot r_{CD} | \Phi_j(B) \rangle &= \\ \sum_{N=1}^7 C \int_0^1 [\frac{1}{\alpha_1} \frac{d(\alpha_1 F_N)}{d\alpha_1} A(N) - \frac{\alpha_1}{u} A(N+1)] e^{-(fg)^{\frac{1}{2}}} du & \quad (5) \end{aligned}$$



$$\frac{d}{d\alpha_2} \langle \Phi_1(A) | \cos \mathbf{K}_V \cdot \mathbf{r}_C | \Phi_j(B) \rangle =$$

$$\sum_{N=1}^7 C \int_0^1 \left[ \frac{1}{\alpha_2} \frac{d(\alpha_2 F_N)}{d\alpha_2} A(N) - \frac{\alpha_2}{1-u} A(N+1) \right] e^{-(fg)^{\frac{1}{2}}} du \quad (6)$$

The function  $F_N$  is always simple enough so that its derivatives can be easily found with the aid of equations (1) and (2). Therefore, by repeated application of equations (3), (4), (5), and (6), the potential integral between any two Slater-type orbitals can be written down with a minimum of effort, rather than the hours of tedious labor which would be required to do the partial derivatives for some of the higher orbitals.

In this paper we are not interested in orbitals higher than 4s. The method can readily be extended to treat more involved cases, although the author suspects that even this simple technique would become quite complicated in dealing with functions other than those of s, p, d, or f symmetry.

We will now list the expressions for the three-center potential integrals which were used in this calculation. Those integrals involving two functions of d symmetry have already been listed (Tyler, Norwood, and Fry, 1969) and will not be shown here. Integrals which are identical except for a permutation of coordinates will have only one member given.

#### S-S Integrals

$$\langle 1s(A) | \cos \mathbf{K}_V \cdot \mathbf{r}_C | 1s(B) \rangle = C \int_0^1 \cos A(1) e^{-(fg)^{\frac{1}{2}}} du$$

$$\langle 2s(A) | \cos \mathbf{K}_V \cdot \mathbf{r}_C | 1s(B) \rangle = C \int_0^1 \left[ -\frac{1}{\alpha_1} A(1) + \frac{\alpha_1}{u} A(2) \right] \cos e^{-(fg)^{\frac{1}{2}}} du$$

$$\langle 3s(A) | \cos \tilde{K}_{\nu} \cdot \tilde{r}_C | 1s(B) \rangle = C \int_0^1 \left[ -\frac{3}{u} A(2) + \left(\frac{\alpha_1}{u}\right)^2 A(3) \right] \cos e^{-(fg)^{\frac{1}{2}}} du$$

$$\langle 4s(A) | \cos \tilde{K}_{\nu} \cdot \tilde{r}_C | 1s(B) \rangle = C \int_0^1 \left[ \frac{3}{u \alpha_1} A(2) - \frac{6\alpha_1}{u^2} A(3) + \left(\frac{\alpha_1}{u}\right)^3 A(4) \right] \cos e^{-(fg)^{\frac{1}{2}}} du$$

$$\langle 2s(A) | \cos \tilde{K}_{\nu} \cdot \tilde{r}_C | 2s(B) \rangle = C \int_0^1 \left[ \left(-\frac{\alpha_1}{\alpha_2 u} - \frac{\alpha_2}{\alpha_1(1-u)}\right) A(2) + \frac{\alpha_1 \alpha_2}{h} A(3) \right] \cos e^{-(fg)^{\frac{1}{2}}} du$$

$$\langle 3s(A) | \cos \tilde{K}_{\nu} \cdot \tilde{r}_C | 2s(B) \rangle = C \int_0^1 \left[ \frac{3}{\alpha_2 u} A(2) - \left(\frac{\alpha_1}{\alpha_2 u} + \frac{3\alpha_2}{h}\right) A(3) + \frac{\alpha_1^2 \alpha_2}{uh} A(4) \right] e^{-(fg)^{\frac{1}{2}}} \cos du$$

$$\langle 4s(A) | \cos \tilde{K}_{\nu} \cdot \tilde{r}_C | 2s(B) \rangle = C \int_0^1 \left[ -\frac{3}{\alpha_1 \alpha_2 u} A(2) + \left(\frac{6\alpha_1}{u^2 \alpha_2} + \frac{3\alpha_2}{h \alpha_1}\right) A(3) - \left(\frac{\alpha_1^3}{\alpha_2 u^3} + \frac{6\alpha_1 \alpha_2}{uh}\right) A(4) + \frac{\alpha_1^{3/2}}{u^2 h} A(5) \right] e^{-(fg)^{\frac{1}{2}}} \cos du$$

$$\langle 3s(A) | \cos \tilde{K}_{\nu} \cdot \tilde{r}_C | 3s(B) \rangle = C \int_0^1 \left[ \frac{2}{h} A(3) - \left(\frac{3\alpha_1^2}{uh} - \frac{3\alpha_2^2}{(1-u)h}\right) A(4) + \frac{\alpha_1^2 \alpha_2^2}{h^2} A(5) \right] \cos e^{-(fg)^{\frac{1}{2}}} du$$

$$\langle 4s(A) | \cos \tilde{K}_{\nu} \cdot \tilde{r}_C | 3s(B) \rangle = C \int_0^1 \left[ -\frac{9}{h \alpha_1} A(3) + \left(\frac{18\alpha_1}{uh} + \frac{3\alpha_2^2}{\alpha_1(1-u)h}\right) A(4) - \left(\frac{6\alpha_1 \alpha_2^2}{h^2} + \frac{3\alpha_1^3}{u^2 h}\right) A(5) + \frac{\alpha_1^3 \alpha_2^2}{uh^2} A(6) \right] \cos e^{-(fg)^{\frac{1}{2}}} du$$

$$\begin{aligned}
\langle 4s(A) | \cos \underline{K}_v \cdot \underline{r}_C | 4s(B) \rangle = & C \int_0^1 \left[ \frac{9}{h\alpha_1\alpha_2} A(3) - \frac{18}{h} \left( \frac{\alpha_1}{\alpha_2 u} + \frac{\alpha_2}{\alpha_1(1-u)} \right) A(4) \right. \\
& + \left( \frac{36\alpha_1\alpha_2}{h^2} + \frac{3\alpha_2^3}{\alpha_1(1-u)^2 h} + \frac{3\alpha_1^3}{\alpha_2 u^2 h} \right) A(5) - 6 \left( \frac{\alpha_1^3 \alpha_2}{u h^2} - \frac{\alpha_1 \alpha_2^3}{(1-u) h^2} \right) A(6) \\
& \left. + \frac{\alpha_1^3 \alpha_2^3}{h^3} A(7) \right] \cos e^{-(fg)^{\frac{1}{2}}} du
\end{aligned}$$

### S-P Integrals

$$\langle 1s(A) | \cos \underline{K}_v \cdot \underline{r}_C | 2P_x(B) \rangle = C \int_0^1 [-D_x u \cos A(1) - K_x \sin A(2)] e^{-(fg)^{\frac{1}{2}}} du$$

$$\begin{aligned}
\langle 2s(A) | \cos \underline{K}_v \cdot \underline{r}_C | 2P_x(B) \rangle = & C \int_0^1 \left[ \frac{D_x u}{\alpha_1} \cos A(1) + \left( \frac{K_x}{\alpha_1} \sin - D_x \alpha_1 \cos \right) A(2) \right. \\
& \left. - \frac{K_x \alpha_1}{u} \sin A(3) \right] e^{-(fg)^{\frac{1}{2}}} du
\end{aligned}$$

$$\begin{aligned}
\langle 3s(A) | \cos \underline{K}_v \cdot \underline{r}_C | 2P_x(B) \rangle = & C \int_0^1 \left[ 3 D_x \cos A(2) + \left( \frac{3K_x}{u} \sin - \frac{\alpha_1^2}{u} D_x \cos \right) A(3) \right. \\
& \left. - \frac{K_x \alpha_1^2}{u^2} \sin A(4) \right] e^{-(fg)^{\frac{1}{2}}} du
\end{aligned}$$

$$\begin{aligned}
\langle 4s(A) | \cos \underline{K}_v \cdot \underline{r}_C | 2P_x(B) \rangle = & C \int_0^1 \left[ -\frac{3D_x}{\alpha_1} \cos A(2) + \left( -\frac{3K_x}{u\alpha_1} \sin + \frac{6\alpha_1 D_x}{u} \cos \right) A(3) \right. \\
& \left. + \left( \frac{6K_x \alpha_1}{u^2} \sin - \frac{D_x \alpha_1^3}{u^2} \cos \right) A(4) - \frac{K_x \alpha_1^3}{u^3} \sin A(5) \right] e^{-(fg)^{\frac{1}{2}}} du
\end{aligned}$$

The corresponding 3p integrals can be obtained by a single application of equation (6) to each of the above integrals. They will not be listed here.

### S-D Integrals

Only integrals of the form 1s-3d and 2s-3d will be listed. The 3s-3d and 4s-3d can be obtained by one and two applications of equation (5), respectively.

$$\langle 1s(A) | \cos \vec{K}_v \cdot \vec{r}_C | 3d_{xy}(B) \rangle = C \int_0^1 [D_x D_y u^2 \cos A(1) + (K_x D_y u + K_y D_x u) \sin A(2) - K_x K_y \cos A(3)] e^{-(fg)^{\frac{1}{2}}} du$$

$$\langle 1s(A) | \cos \vec{K}_v \cdot \vec{r}_C | 3d_{x^2-y^2}(B) \rangle = C \int_0^1 [(D_x^2 - D_y^2) u^2 \cos A(1) + 2u \sin (K_x D_x - K_y D_y) A(2) + (K_y^2 - K_x^2) \cos A(3)] e^{-(fg)^{\frac{1}{2}}} du$$

$$\langle 1s(A) | \cos \vec{K}_v \cdot \vec{r}_C | 3d_{3z^2-r^2}(B) \rangle = C \int_0^1 [D_z^2 u^2 \cos A(1) + (\frac{2-u}{1-u} \cos + 2D_z K_z u \sin) A(2) - (\frac{\alpha_2^2}{3(1-u)^2} \cos + K_z^2 \cos) A(3)] e^{-(fg)^{\frac{1}{2}}} du$$

$$\langle 2s(A) | \cos \vec{K}_v \cdot \vec{r}_C | 3d_{xy}(B) \rangle = C \int_0^1 [-\frac{u^2}{\alpha_1} D_x D_y \cos A(1) + (\alpha_1 u D_x D_y \cos - \frac{u \sin}{\alpha_1} (K_x D_y + K_y D_x)) A(2) + (\frac{K_x K_y}{\alpha_1} \cos + \alpha_1 \sin (K_x D_y + K_y D_x)) A(3) - \frac{K_x K_y}{u} \alpha_1 \cos A(4)] e^{-(fg)^{\frac{1}{2}}} du$$

$$\begin{aligned}
\langle 2s(A) | \cos \mathbf{K}_v \cdot \mathbf{r}_C | 3d_{x^2-y^2}(B) \rangle = C \int_0^1 & \left[ \frac{(D_y^2 - D_x^2)}{\alpha_1} u^2 \cos A(1) + ((D_y K_y - \right. \\
& D_x K_x) \frac{2u}{\alpha_1} \sin + \alpha_1 u \cos (D_x^2 - D_y^2)) A(2) + \left( \frac{K_x^2 - K_y^2}{\alpha_1} \cos \right. \\
& \left. + 2\alpha_1 \sin (K_x D_x - K_y D_y)) A(3) + \frac{K_y^2 - K_x^2}{u} \alpha_1 \cos A(4) \right] \\
& e^{-(fg)^{\frac{1}{2}}} du
\end{aligned}$$

$$\begin{aligned}
\langle 2s(a) | \cos \mathbf{K}_v \cdot \mathbf{r}_C | 3d_{3z^2-r^2}(B) \rangle = C \int_0^1 & \left[ - \frac{D_z^2 u^2 \cos}{\alpha_1} A(1) + (\alpha_1 D_z^2 u \cos \right. \\
& - \frac{(2-u) \cos}{(1-u) \alpha_1} - \frac{2D_z K_z u \sin}{\alpha_1} ) A(2) + \left( \frac{\alpha_2^2 \cos}{3\alpha_1 (1-u)^2} + \frac{K_z^2 \cos}{\alpha_1} \right. \\
& \left. + 2\alpha_1 D_z K_z \sin) A(3) - \frac{\alpha_2^2 \alpha_1}{3h(1-u)} + \frac{\alpha_1 K_z^2}{u} \cos A(4) \right] \\
& e^{-(fg)^{\frac{1}{2}}} du
\end{aligned}$$

### P-P Integrals

Only integrals of the form 2p-2p will be listed. Integrals between 2p and 3p functions can be found from these by a single application of either equation (5) or equation (6), and those between two 3p functions can be found by application of both equation (5) and equation (6) to those given.

$$\begin{aligned}
\langle 2p_x(A) | \cos \mathbf{K}_v \cdot \mathbf{r}_C | 2p_x(B) \rangle = C \int_0^1 & \left[ - D_x^2 h \cos A(1) + (D_x K_x (2u-1) \sin + \cos) \right. \\
& \left. A(2) - K_x^2 \cos A(3) \right] e^{-(fg)^{\frac{1}{2}}} du
\end{aligned}$$

$$\begin{aligned}
\langle 2p_y(A) | \cos \mathbf{K}_v \cdot \mathbf{r}_C | 2p_x(B) \rangle = C \int_0^1 & \left[ - D_x D_y h \cos A(1) + (D_x K_y u - D_y K_x (1-u)) \right. \\
& \left. \sin A(2) - K_x K_y \cos A(3) \right] e^{-(fg)^{\frac{1}{2}}} du
\end{aligned}$$

P-D Integrals

Only integrals of the form 2p-3d will be listed. Integrals between 3p and 3d functions can be found by a single application of equation (5) to those given.

$$\begin{aligned} \langle 2p_x(A) | \cos \tilde{K}_v \cdot \tilde{r}_C | 3d_{xy}(B) \rangle = C \int_0^1 [D_x^2 D_y u h \cos A(1) + (-D_y u \cos + \\ ((K_x D_y + K_y D_x) D_x h - D_x D_y K_x u^2) \sin) A(2) + (-K_y \sin \\ + (K_x^2 D_y u + K_x K_y D_x u - D_x K_x K_y (1-u)) \cos) A(3) \\ + K_x^2 K_y \sin A(4)] e^{-(fg)^{\frac{1}{2}} du} \end{aligned}$$

$$\begin{aligned} \langle 2p_z(A) | \cos \tilde{K}_v \cdot \tilde{r}_C | 3d_{xy}(B) \rangle = C \int_0^1 [D_x D_y D_z u h \cos A(1) + (K_x D_y D_z h \\ + K_y D_x D_z h - D_x D_y K_z u^2) \sin A(2) + (K_x K_z D_y u + K_y K_z D_x u \\ - K_x K_y D_z (1-u) \cos A(3) + K_x K_y K_z \sin A(4)] e^{-(fg)^{\frac{1}{2}} du} \end{aligned}$$

$$\begin{aligned} \langle 2p_x(A) | \cos \tilde{K}_v \cdot \tilde{r}_C | 3d_{3z^2-r^2}(B) \rangle = C \int_0^1 [D_z^2 D_x u h \cos A(1) + (D_x (2-u) \cos \\ + (2D_x D_z K_z h - D_z^2 K_x u^2) \sin) A(2) + (-\frac{K_x (2-u)}{1-u} \sin \\ + (2D_z K_z K_x u - \frac{D_x \alpha_2^2}{3(1-u)} - D_x K_z^2 (1-u)) \cos) A(3) \\ + (\frac{K_x \alpha_2^2}{3(1-u)^2} + K_x K_z^2) \sin A(4)] e^{-(fg)^{\frac{1}{2}} du} \end{aligned}$$

$$\begin{aligned}
\langle 2p_z(A) | \cos \tilde{K}_v \cdot \tilde{r}_C | 3d_{3z^2-r^2}(B) \rangle = C \int_0^1 [D_z^3 u h \cos A(1) + (D_z(2-3u) \cos \\
+ D_z^2 K_z u(2-3u) \sin) A(2) + (\frac{-4+3u}{1-u} K_z \sin + (K_z^2(3u-1) \\
- \frac{\alpha_z^2}{3(1-u)} D_z \cos) A(3) + (\frac{K_z \alpha_z^2}{3(1-u)^2} + K_z^3) \sin A(4)] \\
e^{-(fg)^{\frac{1}{2}}} du
\end{aligned}$$

$$\begin{aligned}
\langle 2p_x(A) | \cos \tilde{K}_v \cdot \tilde{r}_C | 3d_{x^2-y^2}(B) \rangle = C \int_0^1 [(D_x^2 - D_y^2) D_x u h \cos A(1) \\
+ (2h(K_x D_x^2 - K_y D_x D_y) \sin + u^2 (D_y^2 - D_x^2) K_x \sin \\
- 2u D_x \cos) A(2) + (-2K_x \sin + (1-u)(K_y^2 - K_x^2) D_x \cos \\
- 2u K_x \cos (K_y D_y - K_x D_x)) A(3) + K_x (K_x^2 - K_y^2) \sin A(4)] \\
e^{-(fg)^{\frac{1}{2}}} du
\end{aligned}$$

$$\begin{aligned}
\langle 2p_z(A) | \cos \tilde{K}_v \cdot \tilde{r}_C | 3d_{x^2-y^2}(B) \rangle = C \int_0^1 [(D_x^2 - D_y^2) D_z u h \cos A(1) + 2h D_z \\
(K_x D_x - K_y D_y) + u^2 K_z (D_y^2 - D_x^2) \sin A(2) + ((1-u) D_z \\
(K_y^2 - K_x^2) + 2u K_z (K_x D_x - K_y D_y)) A(3) \cos + K_z (K_x^2 - K_y^2) \\
\sin A(4)] e^{-(fg)^{\frac{1}{2}}} du
\end{aligned}$$

## APPENDIX B

In this appendix our purpose is to describe the operations of the cubic group and to give the transformation properties of the d functions under them. Although the transformation matrix for the d functions is five by five, it can be broken down into a two by two block and a three by three block which never mix. These two blocks are treated separately.

In the tabulation below, the notation for the operations is that of Slater (1965). Only 24 of the 48 operations of the cubic group are given. The others can be found by applying the inversion operator to those shown. The d functions are invariant under the inversion operator. Each operation is identified by showing its effect on an arbitrary function of x,y,z. Then its effect on the two sets of d functions is given. Matrices are denoted by brackets rather than parentheses.

<u>Operation</u>	<u>Effect on d functions</u>
$R_1 f(x,y,z) = f(x,y,z)$	$R_1 \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix} = \begin{vmatrix} 1 & 0 \\ 0 & 1 \end{vmatrix} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix}$
	$R_1 \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$
$R_2 f(x,y,z) = f(x,-y,-z)$	$R_2 \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix} = \begin{vmatrix} 1 & 0 \\ 0 & 1 \end{vmatrix} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix}$



$$R_2 \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_3 \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \begin{vmatrix} 1 & 0 \\ 0 & 1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_3 \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_4 \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \begin{vmatrix} 1 & 0 \\ 0 & 1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_4 \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_5 \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & \sqrt{3} \\ -\sqrt{3} & -1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_5 \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_6 \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & \sqrt{3} \\ -\sqrt{3} & -1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_6 f(x, y, z) = f(-y, z, -x)$$

$$R_6 \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 0 & -1 \\ 1 & 0 & 0 \\ 0 & -1 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_7 f(x, y, z) = f(-y, -z, x)$$

$$R_7 \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & \sqrt{3} \\ -\sqrt{3} & -1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_7 \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 0 & 1 \\ -1 & 0 & 0 \\ 0 & -1 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_8 f(x, y, z) = f(y, -z, -x)$$

$$R_8 \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & \sqrt{3} \\ -\sqrt{3} & -1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_8 \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 0 & -1 \\ -1 & 0 & 0 \\ 0 & 1 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_9 f(x, y, z) = f(z, x, y)$$

$$R_9 \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & -\sqrt{3} \\ \sqrt{3} & -1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_9 \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{10} f(x, y, z) = f(-z, -x, y)$$

$$R_{10} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & -\sqrt{3} \\ \sqrt{3} & 1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_{10} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 1 & 0 \\ 0 & 0 & -1 \\ -1 & 0 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{11} f(x, y, z) = f(z, -x, -y)$$

$$R_{11} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & -\sqrt{3} \\ \sqrt{3} & 1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_{11} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & -1 & 0 \\ 0 & 0 & -1 \\ 1 & 0 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{12} f(x, y, z) = f(-z, x, -y)$$

$$R_{12} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & -\sqrt{3} \\ \sqrt{3} & -1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_{12} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & -1 & 0 \\ 0 & 0 & 1 \\ -1 & 0 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{13} f(x, y, z) = f(-x, z, -y)$$

$$R_{13} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & -\sqrt{3} \\ -\sqrt{3} & 1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_{13} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & -1 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{14} f(x, y, z) = f(-x, -z, y)$$

$$R_{14} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & -\sqrt{3} \\ -\sqrt{3} & 1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_{14} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & -1 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{15} f(x, y, z) = f(-z, -y, x) \quad R_{15} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & \sqrt{3} \\ \sqrt{3} & 1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_{15} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 0 & 1 \\ 0 & -1 & 0 \\ -1 & 0 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{16} f(x, y, z) = f(z, -y, -x) \quad R_{16} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & \sqrt{3} \\ \sqrt{3} & 1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_{16} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 0 & -1 \\ 0 & -1 & 0 \\ 1 & 0 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{17} f(x, y, z) = f(y, -x, -z) \quad R_{17} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \begin{vmatrix} 1 & 0 \\ 0 & -1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_{17} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} -1 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{18} f(x, y, z) = f(-y, x, -z) \quad R_{18} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix} = \begin{vmatrix} 1 & 0 \\ 0 & -1 \end{vmatrix} \begin{vmatrix} 3z^2 - r^2 \\ x^2 - y^2 \end{vmatrix}$$

$$R_{18} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} -1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{19}^{f(x,y,z) = f(x,z,y)}$$

$$R_{19} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & -\sqrt{3} \\ -\sqrt{3} & 1 \end{vmatrix} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix}$$

$$R_{19} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{20}^{f(x,y,z) = f(x,-z,-y)}$$

$$R_{20} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & -\sqrt{3} \\ -\sqrt{3} & 1 \end{vmatrix} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix}$$

$$R_{20} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & -1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 1 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{21}^{f(x,y,z) = f(z,y,x)}$$

$$R_{21} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & \sqrt{3} \\ \sqrt{3} & 1 \end{vmatrix} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix}$$

$$R_{21} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{22}^{f(x,y,z) = f(-z,y,-x)}$$

$$R_{22} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix} = \frac{1}{2} \begin{vmatrix} -1 & \sqrt{3} \\ \sqrt{3} & 1 \end{vmatrix} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix}$$

$$R_{22} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 0 & 0 & -1 \\ 0 & 1 & 0 \\ -1 & 0 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{23}f(x,y,z) = f(y,x,z) \quad R_{23} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix} = \begin{vmatrix} 1 & 0 \\ 0 & -1 \end{vmatrix} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix}$$

$$R_{23} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

$$R_{24}f(x,y,z) = f(-y,-x,z) \quad R_{24} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix} = \begin{vmatrix} 1 & 0 \\ 0 & -1 \end{vmatrix} \begin{vmatrix} 3z^2-r^2 \\ x^2-y^2 \end{vmatrix}$$

$$R_{24} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix} = \begin{vmatrix} 1 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & -1 & 0 \end{vmatrix} \begin{vmatrix} xy \\ xz \\ yz \end{vmatrix}$$

## APPENDIX C

In this Appendix we have included listings of computer programs to evaluate kinetic energy, potential energy (three-center) and overlap integrals of the type ss or dd on vanadium. The programs are similar, the principle difference being in Subroutine D3INT. Subroutines which are identical in both programs are not shown and the Gaussian coefficients,  $V(I)$  and  $W(I)$ , have been deleted from the second listing to conserve space. Comment cards in the listing describe the programs, and differences between ss and dd programs illustrate how other programs such as sp or sd may be constructed. These programs are not included, but can be built around the basic ss and dd programs using formulas given in Appendix A.

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V(7)=	C.11369585011C665920911	035
V(8)=-	0.11369585011C665920911	036
V(9)=	0.145973714654896541989	037
V(10)=-	145973714654896941989	038
V(11)=	0.178096882367618602759	039
V(12)=-	178096882367618602759	040
V(13)=	0.210031310460567203603	041
V(14)=-	210031310460567203603	042
V(15)=	0.241743156163840012328	043
V(16)=-	241743156163840012328	044
V(17)=	C.273198812591049141487	045
V(18)=-	273198812591049141487	046
V(19)=	0.304364944354496353024	047
V(20)=-	304364944354496353024	048
V(21)=	0.335208522892625422616	049
V(22)=-	335208522892625422616	050
V(23)=	0.365696861472313635031	051
V(24)=-	365696861472313635031	052
V(25)=	0.395797649828908603285	053
V(26)=-	395797649828908603285	054
V(27)=	0.425478988407300545365	055
V(28)=-	425478988407300545365	056
V(29)=	0.454709422167743008636	057
V(30)=-	454709422167743008636	058
V(31)=	0.483457973920596359768	059
V(32)=-	483457973920596359768	060
V(33)=	0.511694177154667673586	061
V(34)=-	511694177154667673586	062
V(35)=	0.539388108324357436227	063
V(36)=-	539388108324357436227	064
V(37)=	0.566510418561397168404	065
V(38)=-	566510418561397168404	066
V(39)=	0.593032364777572080684	067
V(40)=-	593032364777572080684	068
V(41)=	0.618925840125468570386	069
V(42)=-	618925840125468570386	070

V(43)=0.644163403784967106798	071
V(44)=-.644163403784967106798	072
V(45)=0.668718310043916153953	073
V(46)=-.668718310043916153953	074
V(47)=0.692564536642171561344	075
V(48)=-.692564536642171561344	076
V(49)=0.715676812348967626225	077
V(50)=-.715676812348967626225	078
V(51)=0.738030643744400132851	079
V(52)=-.738030643744400132851	080
V(53)=0.759602341176647498703	081
V(54)=-.759602341176647498703	082
V(55)=0.780369043867433217604	083
V(56)=-.780369043867433217604	084
V(57)=0.800308744139140817229	085
V(58)=-.800308744139140817229	086
V(59)=0.819400310737931675539	087
V(60)=-.819400310737931675539	088
V(61)=0.837623511228187121494	089
V(62)=-.837623511228187121494	090
V(63)= 0.854959033434601455463	091
V(64)=-0.854959033434601455463	092
V(65)= 0.871388505909296502874	093
V(66)=-0.871388505909296502874	094
V(67)= 0.886894517402420416057	095
V(68)=-0.886894517402420416057	096
V(69)= 0.901460635315852341319	097
V(70)=-0.901460635315852341319	098
V(71)= 0.915071423120898074206	099
V(72)=-0.915071423120898074206	100
V(73)= 0.927712456722308690965	101
V(74)=-0.927712456722308690965	102
V(75)= 0.939370339752755216932	103
V(76)=-0.939370339752755216932	104
V(77)= 0.950032717784437635756	105
V(78)=-0.950032717784437635756	106

V(79)= 0.9596882914487425393  
V(80)=-0.9596832914487425393  
V(81)= 0.968326328463264212174  
V(82)=-0.968326828463264212174  
V(83)= 0.975939174585136466453  
V(84)=-0.975939174585136466453  
V(85)= 0.982517263563014677447  
V(86)=-0.982517263563014677447  
V(87)= 0.988054126325623799481  
V(88)=-0.988054126325623799481  
V(89)= 0.99254390323762624572  
V(90)=-0.99254390323762624572  
V(91)= 0.995981842987209290650  
V(92)=-0.995981842987209290650  
V(93)= 0.998364375863181677724  
V(94)=-0.998364375863181677724  
V(95)= 0.99968550383230766828  
V(96)=-0.99968550383230766828  
W(1)=0.032550614452363166242  
W(2)=0.032550614452363166242  
W(3)=0.032516118713868835987  
W(4)=0.032516118713868835987  
W(5)=0.032447163714064269364  
W(6)=0.032447163714064269364  
W(7)=0.032343822568575928429  
W(8)=0.032343822568575928429  
W( 9)=0.032206204794030250669  
W(10)=0.032206204794030250669  
W(11)=0.032034456231992663218  
W(12)=0.032034456231992663218  
W(13)=0.031828758894411006535  
W(14)=0.031828758894411006535  
W(15)=0.03158933070727168558  
W(16)=0.03158933070727168558  
W(17)=0.031316425596861355813  
W(18)=0.031316425596861355813

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W(19)=0.031010332586313837423  
 W(20)=0.031010332586313837423  
 W(21)=0.030671376123669149014  
 W(22)=0.030671376123669149014  
 W(23)=0.030299915420827593794  
 W(24)=0.030299915420827593794  
 W(25)=0.029896344136328385984  
 W(26)=0.029896344136328385984  
 W(27)=0.029461089958167905970  
 W(28)=0.029461089958167905970  
 W(29)=0.028994614150555236543  
 W(30)=0.028994614150555236543  
 W(31)=0.028497411065085385646  
 W(32)=0.028497411065085385646  
 W(33)=0.02797007616848334440  
 W(34)=0.02797007616848334440  
 W(35)=0.027412962726029242823  
 W(36)=0.027412962726029242823  
 W(37)=0.026826866725591762198  
 W(38)=0.026826866725591762198  
 W(39)=0.026212340735672413913  
 W(40)=0.026212340735672413913  
 W(41)=0.025570036005349361499  
 W(42)=0.025570036005349361499  
 W(43)=0.02490063322483610288  
 W(44)=0.02490063322483610288  
 W(45)=0.024204841792364691282  
 W(46)=0.024204841792364691282  
 W(47)=0.023483395085926219842  
 W(48)=0.023483395085926219842  
 W(49)=0.022737065658329374001  
 W(50)=0.022737065658329374001  
 W(51)=0.02196664438744349195  
 W(52)=0.02196664438744349195  
 W(53)=0.021172939892191298988  
 W(54)=0.021172939892191298988

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W(55)=0.020356797154333324595  
 W(56)=0.020355737154333324595  
 W(57)=0.019519091140145022410  
 W(58)=0.019519091140145022410  
 W(59)=0.018660679627411467385  
 W(60)=0.018660679627411467385  
 W(61)=0.017782502316045260838  
 W(62)=0.017782502316045260838  
 W(63)=0.016885479864245172450  
 W(64)=0.016885479864245172450  
 W(65)=0.015970562902562291381  
 W(66)=0.015970562902562291381  
 W(67)=0.015038721026994938006  
 W(68)=0.015038721026994938006  
 W(69)=0.014090941772314860916  
 W(70)=0.014090941772314860916  
 W(71)=0.013128229566961572637  
 W(72)=0.013128229566961572637  
 W(73)=0.012151604671088319635  
 W(74)=0.012151604671088319635  
 W(75)=0.011162102099838498591  
 W(76)=0.011162102099838498591  
 W(77)=0.010160770535008415758  
 W(78)=0.010160770535008415758  
 W(79)=0.009148671230783386633  
 W(80)=0.009148671230783386633  
 W(81)=0.009126876925698759217  
 W(82)=0.008126876925698759217  
 W(83)=0.007096470791153865269  
 W(84)=0.007096470791153865269  
 W(85)=0.006058545504235961683  
 W(86)=0.006058545504235961683  
 W(87)=0.005014202742927517693  
 W(88)=0.005014202742927517693  
 W(89)=0.003964554338444686674  
 W(90)=0.003964554338444686674

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	W(91)=0.002910731817934946408	215
	W(92)=0.002910731817934946408	216
	W(93)=0.001853960738946921732	217
	W(94)=0.001853960738946921732	218
	W(95)=0.000796792065552012429	219
	W(96)=0.000796792065552012429	220
C	READ OR WRITE THE FILE IDENTIFICATION ON THE TAPE	221
C	WRITE(3,16)	222
C	16 FORMAT(1X,'D-D INTEGRALS V-V')	223
	READ(3,16)(GER(N),N=1,6)	224
	16 FORMAT(6A3)	225
	READ(5,6)(GER(I),I=1,30)	226
	6 FORMAT(26A3/4A3)	227
C	READ NUMBER OF TERMS TO BE INCLUDED IN FOURIER SERIES	228
	READ(5,101)JJJJ	229
	101 FORMAT(I5)	230
C	PARAMETER KKR	231
C	1 COMPUTES EXCHANGE ONLY	232
C	2 COMPUTES POTENTIAL ONLY	233
C	3 COMPUTES POTENTIAL AND EXCHANGE	234
	READ(5,101)KKR	235
C	READ A/2	236
	READ(5,641)QA	237
	641 FORMAT(F9.6)	238
	GOTO (17,11,11),KKR	239
C	READ FOURIER COEFFICIENTS OF POTENTIAL	240
C	SUM ONLY OVER VECTORS HAVING NON-NEGATIVE COMPONENTS.PARAMETER IOU	241
C	DENOTES THE NUMBER OF ZERO COMPONENTS IN EACH VECTOR	242
	11 DO 400 J=1,JJJJ	243
	400 READ(5,5)VK(J),B(J),C(J),D(J),IOU(J)	244
	5 FORMAT(1XD16.8,3F6.2,I5)	245
	IF(KKR.EQ.2)GOTO 402	246
C	READ FOURIER COEFFICIENTS OF EXCHANGE	247
	17 DO401 J=1,JJJJ	248
	401 READ(5,5)RK(J) ,B(J),C(J),D(J),IOU(J)	249
	402 CONTINUE	250

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C      READ D WAVEFUNCTIONS FOR VANADIUM IN FOLLOWING ORDER
C      EXPON   COEFF
C      3.0248 5.16807
C      9.971 38.69851
C      4.792 31.22514
C      2.0756 2.39639
C      1.172 0.09345
C      DO10 J=1,5
C      10 READ(5,3)ALD(J),CD(J)
C      3 FOFMAT(2F8.5)
C      READ DIRECT LATTICE COORDINATES
C      60 READ(5,4)AX,AY,AZ,BX,BY,BZ,CX,CY,CZ
C      4 FOFMAT(9F8.5)
C      DO398J=1,15
C      SS(J)=0.00
C      SOL1(J)=0.0
C      398 SOLD(J)=0.00
C      CON1=BX-AX
C      CON2=BY-AY
C      CON3=BZ-AZ
C      CON4=CCN1**2
C      CON5=CON2**2
C      CON6=CON3**2
C      CON7=CCN1**3
C      CON8=CON2**3
C      CON9=CON3**3
C      CON10=CON1*CON2
C      CON11=CON2*CON3
C      CON12=CON1*CON3
C      CON13=CCN4+CON5
C      CON14=CON4+CON6
C      CON15=CON5+CON6
C      CON16=CON4*CON5
C      CON17=CON4*CON6
C      CON18=CON5*CON6
C      CON19=CON4*CON2

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CON20=CON4*CON3
CON21=CON5*CON1
CON22=CON5*CON3
CON23=CON6*CON1
CON24=CON6*CON2
CON25=CON10*CON3
CON26=CON19*CON3
CON27=CON21*CON3
CON28=CON23*CON2
CON34=CON4-CON5
IF(DABS(CON34).GT.0.00001)GO TO 26
CON34=0.00
26 CON30=CON34**2
CON29=CON34*CON10
CON80=CON34*CON11
CON31=CON8*CON3
CON32=CON9*CON2
CON33=CON9*CON1
C SUM OVER DIFFERENT PAIRS OF TERMS IN THE WAVEFUNCTIONS
DO40L=1,5
DO 40 K=1,L
IF(L.EQ.K)GO TO 30
CDC=2.0*CD(L)*CD(K)
GO TO 50
30 CDC=CD(L)*CD(K)
50 ALPHA1=ALD(L)
ALPHA2=ALD(K)
CON35=ALPHA1*ALPHA1/3.0
CON36=ALPHA2*ALPHA2/3.0
CON37=CON9*CON3
CON38=ALPHA1*ALPHA1
CON39=ALPHA2*ALPHA2
CON40=CON38*CON39
CON41=CON40*ALPHA1
CON42=CON38*ALPHA1
CON43=CON39*ALPHA1

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CON44=CCN40*ALPHA2	323
CON45=CON42*ALPHA2	324
CON46=CON39*ALPHA2	325
CON47=CCN46*ALPHA1	326
CON48=ALPHA1*ALPHA2	327
CON49=CON46/ALPHA1	328
CON50=CON42/ALPHA2	329
CON51=ALPHA1/ALPHA2	330
CON52=ALPHA2/ALPHA1	331
CON53=1.0/ALPHA1	332
CON54=1.0/ALPHA2	333
CON55=1.0/CUN48	334
CON56=CON39*CON53	335
CON57=CON51*ALPHA1	336
CON58=CUN48*ALPHA1	337
CON59=ALPHA1*CON4	338
CON60=ALPHA1*CON2	339
CON61=CON2*CON60	340
CON62=ALPHA1*CUN1	341
CUN63=CON2/ALPHA1	342
CON64=CUN3/ALPHA1	343
CON65=CON3*ALPHA1	344
CON66=CON65*CON3	345
CON67=CON53*CON25	346
CON68=CCN25*ALPHA1	347
CON69=CON34*CUN1	348
CON70=CCN34*CON2	349
CON71=CON53*CON23	350
CON72=CON53*CON24	351
CON73=CON36*CON10	352
CON74=CON36*CON11	353
CON75=CON36*CON12	354
CON76={CON35+CON36}*CON6	355
CON77=CUN3*CON36	356
CON78=CON3*CON35	357
CON79=CON35*CUN36	358

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359 CONR1=CON12*CON34
360 CONR2=CON19*CON53
361 CONR3=CON53*CON10
362 CONR4=CON53*CON4
363 CONR5=ALPHA1*CON10
364 CONR6=ALPHA1*CON4
365 CONR7=CON21*CON53
366 CONR8=CON5*CON53
367 CONR9=CON48*CON6
368 RAB=DSQRT(CON1**2+CON2**2+CON3**2)
369 AA=2.*3.14159*ALPHA1*ALPHA2*RAB**3
370 DO100N=1,15
371 SUM1(N)=0.0
372 100 SUM(N)=0.00
373 SUM THE FOURIER SERIES
374 DO500 J=1,JJJJ
375 DO80KOL=1,1440
376 DO FUKOL=0.00
377 N=100(J)
378 AI=3.14159265/QA*B(J)
379 AJ=3.14159265/QA*C(J)
380 AK=3.14159265/QA*D(J)
381 SK =DSQRT(AI*AI+AJ*AJ+AK*AK)
382 RS=SK*SK
383
384 SUBROUTINES D3INT,AGET,AND INTEG COMPUTE THE THREE-CENTER INTEGRAL
385 IF THE INTEGRALS HAVE BEEN STORED ON TAPE,READ THE TAPE INSTEAD OF
386 CALLING D3INT
387
388 CALL D3INT
389 WRITE(3)(S(N),N=1,15)
390 READ(3)(S(N),N=1,15)
391 DO146N=1,15
392 IF(J.NE.1)GOTO145
393 SS(N)=SS(N)+S(N)*CDC
394 145 GOTO (147,148,148),KKR

```

	STOP	395
148	SUM1(N)=SUM1(N)+RK(J)*S(N)	396
	IF(KKR.EQ.2)GOTO 146	397
147	SUM(N)=SUM(N)+VK(J)*S(N)	398
146	CONTINUE	399
500	CONTINUE	400
	DO408J=1,15	401
	SOL1(J)=SUM1(J)*CDC+SOL1(J)	402
408	SOLD(J)=SUM(J)*CDC+SOLD(J)	403
40	CONTINUE	404
	WRITE(6,2)AX,AY,AZ	405
	WRITE(6,2)BX,BY,BZ	406
	WRITE(6,2)CX,CY,CZ	407
2	FORMAT(1X,3F10.6)	408
C		409
C	COMPUTE KINETIC ENERGY INTEGRALS	410
C	SEE SUBROUTINE SIZE FOR DESCRIPTION OF THE PARAMETERS	411
C		412
	NA=3	413
	NB=3	414
	LA=2	415
	LB=2	416
	CA=28.	417
	CB=28.	418
	VOL(1)=1.00/6.00	419
	VOL(2)=2.00/3.00	420
	VOL(3)=1.00	421
	VOL(4)=VOL(2)	422
	VOL(5)=VOL(1)	423
	CALL FACGEN(BF)	424
	CALL AGENT(INDA,AFUN)	425
	OD=RA8	426
	DDP=0.0	427
	DDD=0.0	428
	DDS=0.0	429
	DO 200 J=1,5	430

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DO 200 K=1,5
XA=ALD(J)
XB=ALD(K)
COJ=5.6250/((ALD(J)**3.5)*(ALD(K)**3.5))
DO 300 L=1,5
MA=L-3
MB=MA
ABK=VOL(L)
CALL SZE(E,ZA,ZB,NA,LA,MA,XA,CA,NB,LB,MB,XB,CB,ABK,INDA,AFUN)
300 QL(L)=E
C COMPUTE PARAMETERS FOR THE TWO-CENTER INTEGRALS
DDS=DDS+QL(3)*CD(J)*CD(K)*2.0*COJ
DDP=DDP+(QL(2)+QL(4))*CD(J)*CD(K)*COJ
DDD=DDD+(QL(1)+QL(5))*CD(J)*CD(K)*COJ
200 CALL NIT(EK,CON1,CON2,CON3,RAB,DDC,DDP,DDS)
N1=AX/QA+0.1
N2=AY/QA+0.1
N3=AZ/QA+0.1
PUNCH 108,N1,N2,N3
108 FORMAT(3I5)
DO407J=1,15
IF(KKR.NE.3)GOTO 28
PUNCH 107,SOLD(J),EK(J),SOL1(J),SS(J)
107 FORMAT(4E16.8)
28 KK=2 *J-1
WRITE(6,8)GER(KK),GER(KK+1)
8 FORMAT(1X,2A3)
WRITE(6,1)SOLD(J),EK(J),SOL1(J),SS(J)
1 FORMAT(1X,'POTENTIAL=',E16.8,' KINETIC ENERGY=',E16.8,' EXCHANGE='
1,E16.8,' OVERLAP=',E16.8)
WRITE(6,974)
974 FORMAT(1X,1H )
407 CONTINUE
GO TO 60
END

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000 SUBROUTINE AGET (EAX)
001 IMPLICIT REAL*8(A-H,C-Z)
002 DIMENSION S (20)
003 DIMENSION FU(144),A(9),W(96),V(96)
004 COMMON/BIGD/CON73,CON74,CON75,CON76,CCN82,CON83,CON84,CON85,CON86,
005 1CON87,CON88,CON89,W,V,S,CON77,CUN78,CON80,CON79,
006 3CON81, CON1,CON2,CON3,CCN4,CON5,CON6,CON7,CON8,CON9,CON10,CON11,
007 1CON12,CUN13,CON14,CCN15,CON16,CON17,CON18,CON19,CON20,CON21,CON22,
008 2CON23,CUN24,CON25,CON26,CON27,CUN28,CON29,CON30,CON31,CON32,CON33,
009 3CON34,CON35,CON36,CCN37,CON38,CON39,CON40,CON41,CON42,CON43,CON44,
010 4CON45,CON46,CON47,CON48,CON49,CON50,CCN51,CON52,CON53,CON54,CON55,
011 5CON56,CON57,CUN58,CCN59,CON60,CON61,CON62,CCN63,CCN64,CON65,CON66,
012 6CON67,CON68,CON69,CCN70,CON71,CON72,U,A,RAB,ALPHA1,ALPHA2,RS,FU,
013 7AI,AJ,AK,AX,AY,AZ,BX,BY,BZ,CX,CY,CZ,AA,N
014 F=U*(1.-U)*RAB**2
015 G=RS+ALPHA2**2/(1.-U)+ALPHA1**2/U
016 FG=F*G
017 COB=CSQRT(FG)
018 D1=D08**3
019 D2=D08*D1
020 D3=D2*D08
021 D4=D1*D1
022 D5=D1*D2
023 D6=D2*D2
024 C7=D1*D4
025 D8=D3*D3
026 D9=D1*D6
027 C10=D4*D4
028 D11=D4*D5
029 A(1)=F*(3./D3+3./D2+1./D1)
030 A(2)=F*(15./D5+15./D4+6./D3+1./D2)
031 A(3)=F**3*(105./D7+105./D6+45./D5+10./D4+1./D3)
032 A(4)=F**4*(945./D9+945./D8+420./D7+105./D6+15./D5+1./D4)
033 A(5)=F**5*(10395./D11+10395./D10+4725./D9+1260./D8+210./D7+21./D6
034 1+1./D5)

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A(6)=F**2*(105./C7+105./D6+45./D5+10./E4+1./D3)*RAB**2
A(7)=F**2*(945./D9+945./D8+420./D7+105./D6+15./C5+1./D4)*RAB**4
A(8)=F**3*(10395./D11+10395./D10+4725./D9+1260./D8+210./D7+21./D6
1+1./D5)*RAB**4
IF(DOR.GT.40.)GOTO 222
EAX=DEXP(-DOR)
GOTO 223
222 EAX=C.C
223 RETURN
END
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	SUBROUTINE INTEG	000
	IMPLICIT REAL*8(A-H,O-Z)	001
	COMMON S (20)	002
	DIMENSION FU(144C),A(5),K(96),V(96)	003
	COMMON/RIGD/CON73,CON74,CON75,CON76,CON82,CON83,CON84,CON85,CON86,	004
	1CON87,CUN98,CON89,K,V,S,CON77,CON78,CON80,CON79,	005
	3CON81, CON1,CON2,CON3,CCN4,CON5,CCN6,CCN7,CON8,CON9,CON10,CON11,	006
	1CON12,CON13,CON14,CON15,CON16,CON17,CON18,CON19,CUN20,CON21,CON22,	007
	2CON23,CON24,CON25,CCN26,CON27,CUN28,CON29,CON30,CON31,CON32,CON33,	008
	3CON34,CCN35,CUN36,CCN37,CCN38,CON39,CCN40,CON41,CCN42,CCN43,CON44,	009
	4CON45,CON46,CON47,CCN48,CON49,CUN50,CON51,CON52,CON53,CUN54,CON55,	010
	5CON56,CCN57,CON58,CCN59,CON60,CON61,CON62,CON63,CON64,CUN65,CUN66,	011
	6CON67,CUN68,CON69,CCN70,CON71,CON72,U,A,PAB,ALPHA1,ALPHA2,RS,FU,	012
	7AI,AJ,AK,AX,AY,AZ,BX,BY,BZ,CX,CY,CZ,AA,N	013
	DO 49 I=1,15	014
	DO 50 J=1,96	015
	M=J+(I-1)*96	016
50	S(I)=S(I)+W(J)*FU(M)	017
49	CONTINUE	018
	DO 51 I=1,15	019
51	S(I)=S(I)*0.500	020
	RETURN	021
	END	022

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SUBROUTINE D3INT
IMPLICIT REAL*8(A-H,C-Z)
DIMENSION S (20)
DIMENSION FU(1440),A(9),K(96),V(96)
COMMON/BIQU/CUN73,CUN74,CUN75,CUN76,CUN77,CUN78,CUN79,
1CUN97,CCN86,CCN89,W,V,S,CUN77,CUN78,CUN80,CUN79,
3CCN91, CUN1,CUN2,CCN3,CCN4,CUN5,CCN6,CCN7,CUN8,CCN9,CCN10,CCN11,
1CUN12,CUN13,CUN14,CUN15,CUN16,CUN17,CUN18,CUN19,CUN20,CUN21,CUN22,
2CCN23,CUN24,CUN25,CCN26,CUN27,CUN28,CUN29,CUN30,CUN31,CUN32,CUN33,
3CCN34,CUN35,CUN36,CCN37,CCN38,CUN39,CCN40,CUN41,CCN42,CCN43,CUN44,
4CCN45,CUN46,CUN47,CUN48,CUN49,CUN50,CUN51,CUN52,CUN53,CUN54,CUN55,
5CCN56,CUN57,CUN58,CCN59,CUN60,CUN61,CUN62,CUN63,CUN64,CUN65,CUN66,
6CCN67,CCN69,CUN69,CUN70,CUN71,CUN72,U,A,RAB,ALPHA1,ALPHA2,RS,FU,
7AI,AJ,AK,AX,AY,AZ,BX,BY,BZ,CX,CY,CZ,AA,N
DO150J=1,15
150 S(J)=0.0
ACON1=AI*AI
ACON2=AJ*AJ
ACON3=AK*AK
ACON4=ACON1*AI
ACON5=ACON2*AJ
ACON6=ACON3*AK
ACON7=ACON6*AK
ACON8=AI*AJ
ACON9=AI*AK
ACON10=AK*AJ
ACON11=ACON1*AJ
ACON12=ACON1*AK
ACON13=ACON2*AI
ACON14=ACON2*AK
ACON15=ACON3*AI
ACON16=ACON3*AJ
ACON17=ACON1*ACON2
ACON18=ACON1*ACON3
ACON19=ACON2*ACON3

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	ACON20=A1*AJ*AK	035
	ACON21=ACON1+ACON2	036
	ACON22=ACON1+ACON3	037
	ACON23=ACON2+ACON3	038
	ACON24=ACON13*AK	039
	ACON25=ACON16*AI	040
	ACON26=ACON11*AK	041
	ACON27=ACON1-ACON2	042
	IF(DABS(ACON27).GT.0.00001)GO TO 38	043
	ACON27=0.00	044
38	ACON28=ACON27**2	045
	ACON29=ACON6*AI	046
	ACON30=ACON6*AJ	047
	ACON31=ACON8*ACON27	048
	ACON32=ACON27*ACON10	049
	ACON34=ACON27*ACON9	050
	ACON33=-ACON34	051
	GOTO(10,11,12,15),N	052
15	CON=1.00	053
	GOTO 206	054
10	CON=8.00	055
	GOTO 206	056
11	CON=4.00	057
	GOTO 206	058
12	CON=2.00	059
C	SUM OVER THE 96 POINTS IN THE GAUSSIAN INTEGRATION FORMULA	060
206	DO 750 J=1,96	061
	U=(V(J)+1.00)*0.500	062
	DX=AX*L+BX*(1.-U)	063
	DY=AY*U+BY*(1.-U)	064
	DZ=AZ*L+BZ*(1.-U)	065
	RCDX=DX-CX	066
	RCDY=DY-CY	067
	RCDZ=DZ-CZ	068
	S1=DSIN(AI*RCDX)	069
	S2=DSIN(AJ*RCDY)	070

S3=DSIN(AK*RCOZ)	071
C1=DCJS(AI*RCDX)	072
C2=ECDS(AJ*RCDY)	073
C3=DCDS(AK*RCOZ)	074
SA=S1*C2*C3*CON	075
SB=C1*S2*C3*CON	076
SC=C1*C2*S3*CON	077
SD=-S1*S2*S3*CON	078
CE=C1*C2*C3*CON	079
CC=-S1*S2*C3*CON	080
CB=-S1*C2*S3*CON	081
CA=-C1*S2*S3*CON	082
CAB=1.0-U	083
HH=U*(1.0-U)	084
HH=HH*HH	085
CER=1.-3.*U	086
CIL=2.-3.*U	087
CUR=-4.+3.*U	088
COP=2.0*U-1.0	089
CAB2=U*U	090
CAS3=CAB*CAB	091
CAB4=U*HH	092
CAB5=CAB*HH	093
CAB6=HH*U	094
CAB7=CAB2*HH	095
CAB8=CAB3*HH	096
CAB9=CAB2*U	097
CAB10=CAB3*CAB	098
CAB11=2.-U	099
CAB12=CAB11/CAB	100
CAB13=CAB11/HH	101
CAB14=CAB11/CAB4	102
CAB15=CAB14/U	103
CAB16=U/CAB	104
CAB17=COR*U	105
CAB18=COR/U	106

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107 CAB19=CAB11/U
108 CAB20=CUR/CAB
109 CAB21=CUR/HH
110 CAB22=CER/U
111 CUR=COR*COR
112 CALL AGT(EAX)
113 IF(EAX.F,0.0)GOTO750
114
115 FOLLOWING ARE THE EXPRESSIONS FOR THE 15 DIFFERENT D-D INTEGRALS
116
117 COR78=CON37*CO
118 CUR79=-2.0*CON9*AK*SC
119 COR80=CON6*CO
120 COR81=6.0*CON3*AK*SC
121 COR82= CON6*ACON3*CO
122 COR83=(3.0-CON76)*CO
123 COR84=2.0*CO
124 COR85=2.0*ACON6*CON3*SC
125 COR86=2.0*AK*CON77*SC
126 COR87=-2.0*AK*CON78*SC
127 COR88=ACON3*CO
128 COR89=-2.0*CON35*CO
129 COR90=-2.0*CON36*CO
130 COR91=(CON36+CON35)*CO
131 COR92=ACON7*CO
132 COR93=ACON3*CON36*CO
133 COR94=ACON3*CON35*CO
134 COR95=CON79*CO
135 FU(J)=(COR78*HH+A(1)+(CUR*COR79*HH+(2.0-6.*HH)*COR80)*A(2)+(HH*(
136 1COR81*COR+COR82*(6.0*HH-1.00)+COR83)+COR84)*A(6)+( -HH*COR85*COR+
137 2HH*(COR86*U+COR87*CAB)+( -1.0-6.0*HH)*HH*COR88+COR89*CAB+COR90*U
138 3 *COR91*HH)*A(7)+(COR92*HHH+COR93*CAB2+COR94*CAB3+CUR95)*A(8))*EAX
139 4 *AA#0.89524655
140 M=J+96
141 Q1=CON16*CO
142 B2=-CON13*CO

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H3=A1*CON21*SA+AJ*CON19*SB      143
B4=ACONR*CC*CON10                144
B5=(ACON1*CON5+CON4+ACON21*CO    145
B6=AJ*CON2*SB+A1*CON1*SA          146
B7=-ACON21*CO                      147
B8=ACON11*CON2*SB+ACON13*CON1*SA  148
B9=ACON17*CO                      149
FU(M)=(Q1*HHH*A(1)+(B2-B3*COR)*HH*A(2)+(-B4*COB+CO+B5*HH+B6*COR)*A    150
1(3)+(B7-B3*COR)*A(4)+B9*A(5))*EAX*AA*1.19366207 151
M=J+132                             152
W1=CON17*CO                        153
W2=-CON14*CO                       154
W3=AK*CON20*SC+SA*A1*CON23        155
W4=ACON9*CON12*CB                  156
W5=(ACON3*CON4+CON6*ACON11*CO     157
W6=A1*CON1*SA+AK*CON3*SC           158
W7=-ACON22*CO                      159
W8=ACON15*CON1*SA+ACON12*CON3*SC  160
W9=ACON18*CO                       161
FU(M)=(W1+HHH*A(1)+(W2-W3*COR)*HH*A(2)-(W4*COB-W5*HH-CO-W6*COR)*A(    162
13)+(W7-W8*COR)*A(4)+W9*A(5))*EAX*AA*1.19366207 163
M=J+288                             164
COR1=CON2*AJ*SB+CON3*AK*SC        165
COR2=(ACON10*CON11*CA)             166
COR3=(ACON2*CON6+ACON3*CON5)*CO   167
COR4=CON18*CO                      168
COR5=-CON15*CO                     169
COR6=CON24*AJ*SB+CON22*AK*SC      170
COR7=-ACON23*CO                    171
COR8=ACON14*CON3*SC+ACON16*CON2*SB 172
COR9=ACON19*CO                     173
FU(M)=((COP1*COR-COB*COR2+COR2*HH+CO)*A(3)+COR4*HHH*A(1)+(COR5-    174
1COR6*COR)*HH*A(2)+(COP7-COR8*COR)*A(4)+COR9*A(5))*EAX*AA*1.1936620    175
27                                  176
M=J+384                             177
COR38=CON30*CO                     178

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COR39=-4.0*CON13*CO
COR40=2.0*(AI*CON21*SA+AJ*CON19*SB-CON7*AI*SA-CCN8*AJ*SB)
COR41={CON4*ACON1+ACON2*CON5}*CO
COR42={ACON2*CON4+ACON1*CON5}*CO
COR43=-8.0*CON10*ACON8*CC
COR44=4.0*CO
COR45=2.*{-CON1*ACON4*SA+ACON13*CON1*SA+(ACON11*CON2-CON2*ACON5)*
1SB)
COR46=-4.0*ACUN21*CO
COR47=ACUN28*CO
COR48=4.0*(AI*CON1*SA+AJ*CCN2*SB)
FU(M)=(COR38*HHH*A(1)+(COR39+COR*COR40)*HH*A(2)+(COR*COR48+(6.0*HH
1-1.C)*COR41+(1.0-2.0*HH)*COR42+HH*COR43+COR44)*A(3)+(COR*COR45
2+COR46)*A(4)+COR47*A(5))*EAX*AA*0.29841552
M=J+480
COR24=CON28*CO
COR25=-CON10*CO
COR26=CON24*AI*SA
COR27=-CON23*AJ*SB
COR28=CON25*AK*SC
COR29=ACON9*CON11*CB+ACON3*CON10*CC+ACCN8*CON6*CC+ACON10*CON12*CA
COR30=-ACON9*CON11*CB
COR31=-ACON10*CON12*CA
COR32=-AI*CON2*SA
COR33=CON1*AJ*SB
COR34=ACON15*CON2*SA+ACON20*CON3*SO
COR35=-ACON20*CON3*SD-ACON16*CON1*SB
COR36=-ACON8*CC
COR37=ACON25*CC
FU(M)=(COR24*HHH*A(1)+(COR25+COR26*CB+COR27*U-COR28*COR)*HH*A(2)
1+(HH*COR29+COR30*CB+COR31*CA+COR32*CB+COR33*U)*A(3)+(CA8*COR
234+U*COR35+COR36)*A(4)+COR37*A(5))*EAX*AA*1.19366207
M=J+576
T1={CON17-CON18}*CO
T2=2.*CON34*CO
T3=-CON34*CO

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215 T4=AK*CON3*CON34*SC*2.0*HH
216 T5=2.*(AJ*CON24*SB-AI*CON23*SA)
217 T6=-2.*(AJ*CON2*SB-AI*CON1*SA)
218 T7=CON36*CON5*CO
219 T8=-ACON3*CON34*CO
220 T9=-CON6*ACON27*CO
221 T10=-CON36*CON4*CO
222 T11=4.*(ACON9*CON12*CH-ACON10*CON11*CA)
223 T12=2.*(ACON15*CON1*SA-CON2*SB*ACON16)
224 T13=2.*(ACON14*CON3-ACON12*CON3)*SC
225 T14=2.*CON36*(CON1*AI*SA-CON2*AJ*SB)
226 T15= ACON27*CO
227 T16=ACON3*ACON27*CO
228 T17=CON36*ACON27*CO
229 FU(M)=(T1*HHH*A(1)+(T2*CAB+T3*HH+T4*CAB+T5*U*HH)*A(2)+(T6*(U-2.))+T
230 17+T8*CAB3+T9*CAB2+T10+T11*HH)*A(3)+(T12*CAB*HHH+T13*CAB6+T14*CAB4+
231 2T15*(U-2.)*CAB4)*A(7)+(T16*HHH+T17*CAB2)*A(8))*EAX*AA*.51687084
232 M=J+672
233 U1=CON26*CO
234 U2=-CON11*CO
235 U3=CON19*AK*SC+CON25*AI*SA
236 U4=-CON20*AJ*SB-CON25*AI*SA
237 U5=ACON9*CON10*CB+ACON1*CON11*CO+ACON10*CON4*CA+ACON8*CON12*CC
238 U6=-ACON9*CON10*CB
239 U7=-ACON8*CON12*CC
240 U8=-AK*CON2*SC
241 U9=AJ*CON3*SB
242 U10=ACON12*CON2*SC+ACON20*CON1*SD
243 U11=-ACON20*CON1*SD-ACON11*CON3*SB
244 U12=-ACON10*CA
245 U13=ACON26*CA
246 FU(M)=(U1*HHH*A(1)+(U2+U3*CAB+U4*U)*A(2)*HH+(U5*HH+U6*CAB3+U7*CAB2
247 1+U8*CAB+U9*U)*A(3)+(U10*CAB+U11*U+U12)*A(4)+U13*A(5))*EAX*AA
248 2*1.19366207
249 M=J+768
250 E1=CON29*CO

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251 E2=AI*CON19*SA-AJ*CON21*S8
252 E3=AJ*CON7*SR-AI*CON8*SA
253 E4=ACON27*CON10*CO
254 E5=ACON8*CON34*CC*U
255 E6=2.*(AJ*CON1*SB-AI*CON2*SA)
256 E7=ACON4*CON2*SA-ACON5*CON1*S8
257 E8=ACON11*CON1*SB-ACON13*CON2*SA
258 E9=ACON31*CC
259 FU(M)=(E1*HHH*A(1)+(E2*CIL-E3*U)*HH*A(2)+(-E4*CAR*CER+E5*CIL+E6)*A
260 1(3)+(E7*CAR+E8*CER)*A(4)+E9*A(5))*EAX*AA*0.59683104
261 M=J+864
262 H1=CON32*CO
263 H2=AK*CON24*SC
264 H3=-AJ*CON9*S8
265 H4=CON11*CO
266 H5=AK*CON2*SC
267 H6=-AJ*CON3*S8
268 H7=ACON3*CON11*CO
269 H8=ACON10*CON6*CA*U
270 H9=-CON74*CO
271 H10=ACON10*CA
272 H11=ACON6*CON2*SC
273 H12=ACON16*CON3*S8
274 H13={AK*CON2*SC+AJ*CON3*S8)*CON36
275 H14=ACON30*CA
276 H15=CON36*ACON10*CA
277 FU(M)={H1*HHH*A(1)+(H2*HH*CIL+H3*U*HH+H4*(-3.*H+2.*CAB))*A(2)+[H5
278 1*CUB+H6*CIL-H7*CAB*CER+H8*CIL+H9)*A(3)+[H10*CUB*CAB4+[H11*CAB+H12
279 2*CER)*HHH+H13*CAB4)*A(7)+[H14*HHH+H15*CAB2)*A(8))*EAX*AA*1.0337416
280 M=J+960
281 COK63=CON33*CO
282 COR64=CON23*AK*SC
283 COR65=-AI*CON9*SA
284 COR66=CON12*CO
285 COR67=CON1*AK*SC
286 COR68=CON3*AI*SA

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287 COR69=(ACON3*CON12)*CO
288 COR70=CON6*ACON9*CB*U
289 COR71=-CON75*CO
290 COR72=ACON9*CB
291 COP73=ACON6*CON1*SC
292 CUP74= ACON15*CON3*SA
293 COR75=CON36*(AK*CON1*SC+AI*CON3*SA)
294 COR76=ACON29*CB
295 COR77=CON36*ACON9*CB
296 FU(M)=(COR63*HHH*A(1))+((CIL*COR64+U*COR65)*HH+(-3.0*HH+2.0*CB)*
297 1CUP66)*A(2)+(CUB*COR67-CIL*CUR68-CAB*CER*COR69+CIL*COR70+COR71)*
298 2A(3)+((COR72*CUB+COR75)*CB4+(COR73*CB8+COR74*CER)*HH)*A(7)+(
299 3COR76*HH+COR77*CB82)*A(8))*EAX*AA*1.03374168
300 M=J+1056
301 G1=CON28*CO
302 G2=CON10*CO
303 G3=2.*AK*CON25*SC
304 G4=-AJ*CON23*SB-AI*CON24*SA
305 G5=CON1*AJ*SB+AI*CON2*SA
306 G6=-ACON3*CON10*CO
307 G7=2.*(ACON10*CON12*CA+ACON9*CON11*CB)
308 G8=-ACON8*CON6*CC
309 G9=-CJN73*CO
310 G10=ACON8*CC
311 G11=ACON16*CON1*SB+ACON15*CON2*SA
312 G12=-2.*CON3*ACON20*SD
313 G13=CON36*(AI*CON2*SA+AJ*CON1*SB)
314 G14=ACON25*CC
315 G15=CON36*ACON8*CC
316 FU(M)=(G1*HHH*A(1))+((G2*(CIL+CB82)+(G3*CB3+G4*U)*HH)*A(2)+(G5*(U-2.
317 1)+ G6*CB83+G7*HH+G8*CB82+G9)*A(3)+ ((U-2.)*G10*CB4+(G11*CB8+G12*U
318 2)*HHH+G13*CB4)*A(7)+(G14*HHH+G15* CB82)*A(8))*EAX*AA*1.03374168
319 M=J+1152
320 COR10=CON27*CO
321 COR11=-CON12*CO
322 COR12=CON21*AK*SC

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323 COR13=CON22*AI*SA
324 COR14=CON25*AJ*SB
325 COR15=CON10*ACCN10*CA+CON12*ACON2*CO+CCN5*ACON9*CB+CON11*ACON8*CC
326 COR16=-(CON10*ACON10)*CA
327 COR17=-(CON11*ACON8)*CC
328 COR18=-CON1*AK*SC
329 COR19=CON3*AI*SA
330 COR20=ACON14*CCN1*SC+ACON20*CON2*SD
331 COR21=-ACON20*CON2*SD-ACON13*CON3*SA
332 COR22=-ACON9*CB
333 COR23=ACON24*CB
334 FU(M)=(COR10*HH*A(1)+(COR11+COR12*CAB-COR13*U-COR14*COR )*HH
335 1*A(2)+(HH*COR15+COR16*CAB3+COR17*CAB2+COR18*CAB+COR19*U)*A(3)+(
336 2COR20*CAB+COR21*U+COR22)*A(4)+COR23*A(5))*EAX*AA*1.19366207
337 M=J+1248
338 F1=CON80*CO
339 F2=2.*(AI*CON25*SA-AJ*CON22*SB)
340 F3=-CON34*(AJ*CON3*SB+AK*CCN2*SC)
341 F4=2.*CON11*CO
342 F5=2.*AJ*CON3*SB
343 F6=-2.*AK*CON2*SC
344 F7=-CON11*ACON27*CO
345 F8=2.*(ACON8*CCN12*CC+ACON9*CON10*CB-ACON2*CON11*CO)
346 F9=-ACON10*CON4*CA
347 F10=-ACON10*CON5*CA*U
348 F11=ACON11*CON3*SB+ACON12*CON2*SC-ACON14*CON2*SC-ACON5*CON3*SB
349 F12=-2.*(ACON20*CON1*SD-ACON14*CON2*SC)
350 F13=2.*ACON10*CA
351 F14=ACON32*CA
352 FU(M)=(F1*HH*A(1)+(F2*CAB+F3*U+F4)*HH*A(2)+(F5*CAB+F6*U+F7*CAB3+F
353 18*HH+F9*CAB2+F10*CIL)*A(3)+(F11*CAB+F12*U+F13)*A(4)+F14*A(5))*EAX
354 2*AA*0.59683104
355 M=J+1344
356 COR49=CON81*CO
357 COR50=2.*(CON20*AI*SA-CON25*AJ*SB)
358 COR51=-CON34*(CON3*AI*SA+CON1*AK*SC)

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COR52=2.0*CON12*CO
COR53=-2.0*AI*CON3*SA
COR54=2.0*AK*CON1*SC
COR55=-CON12*ACON27*CO
COR56=-2.0*(CON11*ACON8*CC+CON10*ACON10*CA-ACON1*CON12*CO)
COR57=ACUN9*CON5*CB
COR58=ACON9*CON4*CB*U
COR59=-ACON13*CON3*SA-ACON14*CON1*SC+ACON12*CON1*SC+ACON4*CON3*SA
COR60=2.0*(ACON20*CON2*SD-ACON12*CCN1*SC)
COR61=-2.0*ACON9*CB
COR62=ACON34*CB
FU(M)=(COR49*HHH*A(1)+(COR50*CB8+U*COR51-COR52)*HH*A(2)+(CA8*COR53
1+U*COR54+CA83*COR55+HH*COR56+CB2*COR57+CIL*COR58)*A(3)+(CA8*COR59
2+U*COR60+COR61)*A(4)+COR62*A(5))*EAX*AA*0.59683104
750 CONTINUE
CALL INTEG
RETURN
END

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C
C      THIS SUBROUTINE COMPUTES THE KINETIC ENERGY INTEGRALS USING THE
C      RELATIONS DEVELOPED BY SLATER AND KOSTER
SUBROUTINE NIT(EK,CCN1,CON2,CON3,RAB,DDD,DDP,DDS)
  IMPLICIT REAL * 8 ( A - H ,   O - Z )
  DIMENSION EK(15)
  A=CON1/RAB
  B=CON2/RAB
  C=CON3/RAB
  AA=A*A
  BB=B*B
  CC=C*C
  CJN1=3.0*BB*CC
  CCN1=CCN1*DDS
  CON2=BB+CC-4.*BB*CC
  CON2=CON2*DDP
  CGN3=AA+BB*CC
  CON3=CON3*DDO
  EK(4)=CON1+CON2+CON3
  CON1=3.*CC*AA
  CON1=CCN1*DDS
  CON2=CC+AA-4.*AA*CC
  CON2=CON2*DDP
  CON3=BB+AA*CC
  CON3=CON3*DDO
  EK(3)=CON1+CON2+CON3
  CON1=3.*B*CC*A
  CON1=CON1*DDS
  CON2=8*A*(1.-4.*CC)
  CON2=CON2*DDP
  CGN3=A*B*(CC-1.)
  CON3=CON3*DDO
  EK(6)=CON1+CON2+CON3
  CON1=3.*AA*BB
  CON1=CON1*DDS
  CCN2=AA+BB-4.*AA*BB

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CUN2=CCN2*DDP	035
CUN3=CC+AA*BB	036
CUN3=CCN3*DDD	037
EK(2)=CON1+CON2+CUN3	038
CUN1=3.*A*BB*C	039
CUN1=CON1*DDS	040
CUN2=A*C*(1.-4.*BB)	041
CUN2=CCUN2*DDP	042
CUN3=A*C*(BB-1.)	043
CUN3=CCN3*DDD	044
EK(13)=CON1+CON2+CON3	045
CUN1=3.*AA*B*C	046
CUN1=CCUN1*DDS	047
CUN2=8*C*(1.-4.*AA)	048
CUN2=CCUN2*DDP	049
CUN3=B*C*(AA-1.)	050
CUN3=CCUN3*DDD	051
EK(8)=CUN1+CON2+CON3	052
CUN1=1.5*A*B*(AA-BB)	053
CUN1=CON1*DDS	054
CUN2=2.*A*B*(BB-AA)	055
CUN2=CCUN2*DDP	056
CUN3=.5*A*B*(AA-BB)	057
CUN3=CCUN3*DDD	058
EK(9)=CON1+CON2+CON3	059
CUN1=1.5*3*C*(AA-BB)	060
CUN1=CON1*DDS	061
CUN2=-B*C*(1.+2.*(AA-BB))	062
CUN2=CCUN2*DDP	063
CUN3=B*C*(1.+0.5*(AA-BB))	064
CUN3=CCUN3*DDD	065
EK(14)=CON1+CON2+CON3	066
CUN1=1.5*A*C*(AA-BB)	067
CUN1=CON1*DDS	068
CUN2=A*C*(1.-2.*(AA-BB))	069
CUN2=CCUN2*DDP	070

CON3=-A*C*(1,-.5*(AA-BB))	071
CON3=CON3*DD	072
EK(15)=CON1+CON2+CON3	073
E=DSORT(3,DD)	074
CON1=E*A*B*(CC-.5*(AA+BB))	075
CON1=CON1*DDS	076
CON2=-2.*E*A*3*CC	077
CON2=CON2*DDP	078
CON3=.5*E*A*3*(1.+CC)	079
CON3=CON3*DD	080
EK(12)=CON1+CON2+CON3	081
CON1=E*B*C*(CC-.5*(AA+BB))	082
CON1=CON1*DDS	083
CON2=E*R*C*(AA+BB-CC)	084
CON2=CON2*DDP	085
CON3=-.5*E*9*C*(AA+BB)	086
CON3=CON3*DD	087
EK(10)=CON1+CON2+CON3	088
CON1=E*A*C*(CC-.5*(AA+BB))	089
CON1=CON1*DDS	090
CON2=E*A*C*(AA+BB-CC)	091
CON2=CON2*DDP	092
CON3=-.5*E*A*C*(AA+BB)	093
CON3=CON3*DD	094
EK(11)=CON1+CON2+CON3	095
CON1=.75*(AA-BB)**2	096
CON1=CON1*DDS	097
CON2=AA+BB-(AA-BB)**2	098
CON2=CON2*DDP	099
CON3=CC+.25*(AA-BB)**2	100
CON3=CON3*DD	101
EK(5)=CON1+CON2+CON3	102
CON1=.5*E*(AA-BB)*(CC-.5*(AA+BB))	103
CON1=CON1*DDS	104
CON2=E*C*C*(BB-AA)	105
CON2=CON2*DDP	106

CON3=.25*E*(1.+CC)*(AA-BB)	107
CON3=CON3*DDD	108
EK(7)=CON1+CON2+CON3	109
CON1=(CC-.5*(AA+BB))**2	110
CON1=CON1*DDS	111
CON2=3.*CC*(AA+BB)	112
CON2=CON2*DDP	113
CON3=.75*(AA+BB)**2	114
CON3=CON3*DDD	115
EK(1)=CON1+CON2+CON3	116
RETURN	117
END	118

	SUBROUTINE SZE(E,ZA,ZB,NA,LA,MA,XA,CA,NB,LB,MB,XB,CB,ABK,INDA,AFUN	000
	1)	001
	IMPLICIT REAL * 8 ( A - H , O - Z )	002
	COMMON/CFUN/P(40),Q(40),R(40),D(9),DX,SCRAP(150),FAA(10),FBB(10),O	003
	*D/FACL/BF(100)	004
	DIMENSION INDA(5,5,5),AFUN(9,9,55)	005
C		006
C	FUNCTION SZE IS SET UP TO CALCULATE THE ONE-ELECTRON INTEGRALS	007
C	(NA,LA,MA,XA/OP/NB,LB,MB,XB) , WHERE OP IS THE OPERATOR FOR	008
C	RESPECTIVELY - OVERLAP,KINETIC ENERGY,AND EXCHANGE NUCLEAR	009
C	ATTRACTION (RA**=1 AND RB**=1)	010
C	THE PARAMETERS ARE--	011
C	E KINETIC ENERGY INTEGRAL	012
C	ZA NUCLEAR ATTRACTION (A/RA**=1/B)*CA	013
C	ZB NUCLEAR ATTRACTION (A/RB**=1/B)*CB	014
C	NA,LA,MA,XA,CA ARE RESPECTIVELY THE DEFINING QUANTUM NOS.	015
C	ORBITAL EXPONENT AND CORE CHARGE FOR A	016
C	NB,LB,MB,XB,CB DITTO FOR CENTRE B	017
C	ABK IS THE FACTOR KAB FROM THE PAPERS BY N.M.KLIMENKO AND M.E.	018
C	DYATKINA IN J.STRUCT.CHEM. (ENGLISH TRANSLATION) VOL.6.,573 AND	019
C	714(1965).	020
C	THE ALGORITHM IS ESSENTIALLY THAT OUTLINED BY THEM,BUT MODIFIED TO	021
C	TAKE ADVANTAGE OF SIMILARITIES IN THE REQUIRED C-FUNCTIONS.	022
C	NOTE THAT THE ONE CENTRE NUCLEAR ATTRACTION INTEGRAL IS NON-ZERO	023
C	IN ZA,ZERO IN ZB.	024
C		025
	E = 0.0	026
	ZA = 0.0	027
	ZB = 0.0	028
	INIT=0	029
	IF(MA-MB)1000,20,1000	030
20	IF(OD)60,40,60	031
40	IF(LA-LB)1000,60,1000	032
60	M=IABS(MA)	033
	V=XA/XB	034

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V1=V*(NA+0.5)
V2=V1/V
V3=V2/V
FACS=A8K*V1
FACE=-C.5*X1*XA*A8K
FACZA=XA*A6K*V2
FACZB=-FACZA*CB
FACZA=-FACZA*CA
FACE1=V1
FACE2=-V2*(NA+NA)
NL1=NA-LA-1
IF(NL1)100,100,80
80 FACE3=(NA+LA)*NL1*V3
   LF=3
   GO TO 120
100 FACE3=0.
   LF=2
120 LAM=LA-M
   L9M=LB-M
   LAM1=LAM-1
   LB1=LB1-1
   FAA(1) = 1.
   F88(1) = 1.
   LIMA = 0
   LIMB = 0
   IF(LAM1)160,16C,140
140 FAA(2) = -1.D+00 * LAM * LAM1 / ( 4.D+00 * LA - 2.D+00 )
   LIMA=1
160 IF(LB1)200,200,180
180 F88(2) = -1.D+00 * LB1 * LB1 / ( 4.D+00 * LB - 2.D+00 )
   LIMB=1
200 NAB=NA+NB
   LAB=LA+LB
   IF(OD)300,300,500
C   PD ZERO
300 XX=XB/(XA+XB)

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	DX=0.	071
	IF 1=LAM+LBM+1	072
	JF 1=LAR/2+1	073
	KF 1=(LAM+LBM)/2+1	074
	LF 1=LAB+2	075
	DO 320 KA=INIT,LMA	076
	FA=FAA(KA+1)	077
	KA2=KA+KA	078
	IF 2=IF 1-KA2	079
	JF 2=JF 1-KA	080
	KF 2=KF 1-KA	081
	LF 2=LF 1-KA2	082
	DO 320 KB=INIT,LMB	083
	F3=FB8(KB+1)*FA	084
	KB2=KB+KB	085
	IF 3=IF 2-KB2	086
	JF 3=JF 2-KB	087
	KF 3=KF 2-KB	088
	LF 3=LF 2-KB2	089
320	DX=DX+FB*3F(IF 3)*BF(JF 3)/(BF(KF 3)*BF(LF 3))	090
	NAB2=NAB+2	091
	DX=DX*2.*(M+M+1)*BF(M+1)*XX**NAB2	092
360	DO 380 I1=1,LF	093
	CX=DX/XX	094
380	D(I1)=DX*HF(NAB2-I1)	095
400	ZA=FACZA*D(2)	096
	E=FACE1*D(1)+FACE2*D(2)	097
	IF(FACE3)420,440,420	098
420	E=E+FACE3*D(3)	099
440	E=E*FACE	100
	GO TO 1000	101
C	PO NON-ZERO	102
C500	D(1)=D(2)=D(3)=DB=0.	103
500	D(1) = 0.0	104
	D(2) = 0.0	105
	D(3) = 0.0	106

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DB      = 0.0
RH=X8*OD
RA=XA*OD
RU=RA+RB
IF(R0.GT.160.)GO TO 1000
RT=RA-RB
EPRB = DEXP ( RB )
ENRB=1./EPRB
NMAX=NAB+LAB
NMIN=0
IREF=1-NMIN
CALL PCFN(R0,RT,NMIN,NMAX,P,Q,R,BF,IREF)
N1=NA-LA+IREF
LAM1=LAM+1
LBM1=LBM+1
M1=M+1
NLB=NB-LB
DO 900 KA=INIT,LIMA
KA2=KA+KA
FA=FAA(KA+1)
LAM2=LAM1-KA2
LR1=LAB-KA2
N2=N1+KA2
DO 500 KB=INIT,LIMB
KB2=KB+KB
FB=FA*FB(KB+1)
LBM2=LBM1-KB2
NLB1=NLB+KB2
LR=LR1-KB2
IND=INDA(LAM2,LBM2,M1)
DO 900 IR=INIT,LR
IR1=IR+1
IR2=IR+IR
N3=N2+IR2
LS=LR-IR
DO 880 IS=INIT,LS

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143 IS1=IS+1
144 IS2=IS+IS
145 AA=AFUN(IR1,IS1,IND)
146 IF(AA)520,880,520
147 FC=FR*AA
148 LK=NLB1+IS2
149 FC2=FC*BF(LK)
150 FC1=FC*BF(LK+1)
151 FD=FC1
152 N4=N3+1
153 DO 700 IK=INIT,LK
154 FD=FD/RB
155 LL=LK-IK
156 DO 700 IL=INIT,LL
157 LLL=LL-IL
158 FE=FD/(BF(IL+1)*BF(LL+1))
159 N5=N4+IL
160 FP = 1.
161 FR = 1.
162 IF(LLL/2*2-LLL)600,620,600
163 FP=-1.
164 IF(IL/2*2-IL)640,660,640
165 FR=-1.
166 DO 700 IF=1,LF
167 N=N5-IF
168 D(IF)=D(IF)+FE*(FP*EPRB*P(N)+ENR3*(FR*R(N)-Q(N)))
169 N4=N3
170 FD=FC2
171 LK=LK-1
172 DO 800 IK=INIT,LK
173 FD=FD/RB
174 LL=LK-IK
175 DO 900 IL=INIT,LL
176 LLL=LL-IL
177 FE=FD/(BF(IL+1)*BF(LL+1))
178 N=N4+IL

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	FP = 1.	179
	FR = 1.	180
	IF(LLL/2*2-LLL)720,740,720	181
720	FP=-1.	182
740	IF(IL/2*2-IL)760,800,760	183
760	FR=-1.	184
800	DB=DB+FE*(FP*EPRB*P(N)+ENRB*(FR*R(N)-Q(N)))	185
880	CONTINUE	186
900	CONTINUE	187
	FD=RB**NAB	188
	DB=DB*FD	189
	FD=FD*RB	190
	D(1)=D(1)*FD	191
	DO 920 KA=2,LF	192
	FD=FD/RB	193
920	D(KA)=D(KA)*FD	194
	ZB=FACZb*DB	195
	GO TO 400	196
1000	RETURN	197
	END	198

	SUBROUTINE FACGEN(BF)	000
	IMPLICIT REAL * 8 ( A - H , O - Z )	001
C	** THIS SUBROUTINE HAS BEEN CHANGED FROM 100 TO 57 JHRS **	002
C	FACTORIAL GENERATOR	003
	DIMENSION BF(100)	004
	BF(1)=1.	005
C	DO 10 I=1,99	006
	DO 10 I=1,56	007
10	BF(I+1)=I*BF(I)	008
C		009
	DO 20 I = 57 , 94 , 5	010
	BF(I+1) = + 10.D-78	011
	BF(I+2) = - 10.D+63	012
	BF(I+3) = 0.D+00	013
	BF(I+4) = - 10.D-78	014
	BF(I+5) = + 10.D+63	015
20	CONTINUE	016
	RETURN	017
	END	018

SUBROUTINE AGENT(INDA,AFUN)	000
IMPLICIT REAL * 8 ( A - H , O - Z )	001
COMMON/FACL/BF(100)	002
DIMENSION INDA(5,5,5),AFUN(9,9,55)	003
INIT = 0	004
N = 0	005
DO 100 M=INIT,4	006
MM=M+1	007
LIM=4-M	008
DO 100 K=INIT,LIM	009
KK=K+1	010
DO 100 L=INIT,LIM	011
LL=L+1	012
N=N+1	013
INDA(KK,LL,MM)=N	014
LIMR=K+L+M+M	015
DO 100 IR=INIT,LIMR	016
JR=IR+1	017
LIMS=LIMR-IR	018
DO 100 IS=INIT,LIMS	019
JS=IS+1	020
AA=A(IR,IS,K,L,M)	021
IF(DABS(AA)-1.E-8)90,90,100	022
90 AA=0.	023
100 AFUN(JR,JS,N)=AA	024
RETURN	025
END	026

	SUBROUTINE PQRN(RD,RT,NMIN,NMAX,P,Q,R,BF,IREF)	000
	IMPLICIT REAL * 8 ( A - H , O - Z )	001
	DIMENSION P(40),Q(40),R(40),BF(100)	002
	NMINA=1ABS(NMIN)	003
	IREF=NMINA+1	004
	RUN=-RD	005
	ER = DEXP ( RUN )	006
	NMINA1=NMINA-1	007
C	PN,QN BY FORWARD RECURRENCE ALL N.GT.0,ALL RD.GT.0	008
	P(IREF) = ER / RD	009
	PP = ER / RD	010
	Q(IREF) = 1.0D+00 / RD	011
	RPD = 1.0D+00 / RD	012
	IF(NMAX)50,50,20	013
20	DO 40 I=1,NMAX	014
	RR=RPD*I	015
	IRI=IREF+I	016
	P(IRI)=PP+RR*P(IRI-1)	017
40	Q(IRI)=RR*Q(IRI-1)	018
50	IF(NMIN)60,480,480	019
C	PN,QN N.LT.0	020
60	IF(RD-4.)80,300,300	021
C	PN BY FORWARD RECURRENCE EX PO N.LT.0,RD.LT.4	022
80	CALN=DLOG(RD)+0.577215664901533	023
	IF(RD-1.)100,140,140	024
100	NT=10.-DLOG10(RD)	025
	IF(NT-5)120,160,160	026
120	NT=5	027
	GO TO 160	028
140	NT=6.*RD+11.	029
C	EI BY SERIES EXPANSION	030
160	TB = RON	031
	EI = RON	032
	DO 180 I=2,NT	033
	TB=TB*RON/I	034

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180 EI=EI+TR/I
180 EI=EI+CALN
180 P(IREF-1)=-EI
180 Q(IREF-1)=-CALN
180 IF(NMINA1)480,480,200
180 PN FORWARD RECURRENCE, SET Q APPRAY TO ZERO
180 DO 220 N=2,NMINA
180 IRN=IREF-N
180 C(IRN)=0.
180 P(IRN)=(ER-RO*P(IRN+1))/(N-1.)
180 ENTRY FOR QN CALCULATION, WITH Q(-1) SET, OTHERS ZERO
180 Q2=1.
180 F1=RO
180 Q(IREF-2)=F1*(CALN-Q2)+Q2
180 IF(NMINA-2)480,480,260
180 DO 280 I=3,NMINA
180 IRI=IREF-I
180 Q3 = 1.0+00 / ( I - 2. )
180 XI = 1.0+00 / ( I - 1. )
180 F1=F1*RON*XI
180 Q2=Q2+XI
180 Q(IRI)=Q(IRI)+F1*(CALN-Q2)+XI
180 DO 280 J=I,NMINA
180 IRJ=IREF-J
180 Q3=Q3*RON/(J-I+1.)
180 Q(IRJ)=Q(IRJ)+Q3
180 GO TO 480
180 PN BY BACK RECURRENCE EX CONTINUED FRACTION N.LT.0,RO.GE.4
180 IF(RO-100.)340,340,320
180 NT=-5
180 GO TO 400
180 IF(RO-6.)380,360,360
180 NT=-100./RO-6.
180 GO TO 400
180 NT=-24
180 NL=-1

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	CF=0.	071
	DO 420 I=NT,NL	072
	CC=1.+CF	073
420	CF=-I*CC/(RO*CC+NMINA-I)	074
	P(I)=ER*(1.+CF)/(RO*(1.+CF)+NMINA)	075
C	SET Q(-1)	076
	CALN=DLOG(RO)+0.577215664901533	077
	Q(IREF-1)=-CALN	078
	IF(NMINA)480,48C,440	079
C	PN BACK RECURRENCE,SET Q ARRAY TO ZERO	080
440	DO 460 I=2,NMINA	081
	IRN=I-IREF	082
	P(I)=(ER+IRN*P(I-1))/RO	083
460	Q(-IRN)=0.	084
	GO TO 240	085
C	START ON RN CALCULATION FIRSTLY,IS TAU ZERO	086
480	IF(RT)580,500,580	087
C	RN=1/(N+1) FOR N.GT.0	088
500	R(IREF)=1.	089
	IF(NMAX)530,530,510	090
510	DO 520 I=1,NMAX	091
520	R( IREF + I ) = 1.0+CO / ( I + 1. )	092
530	IF(NMIN)540,3000,3000	093
540	DO 560 I=1,NMINA	094
560	R(I)=0.	095
	GO TO 3000	096
C	SORT OUT VALUES OF RT FOR RN(N.GE.0)	097
580	ART = DABS ( RT )	098
	RTN=-RT	099
	ERT = DEXP ( RTN )	100
	IF(RT-80.)680,60C,600	101
C	RN=QN FOR N.GE.0,RT.GE.80	102
600	RRT = 1.0+CO / RT	103
	R(IREF) = RRT	104
	IF(NMAX)660,66C,620	105
620	DO 640 I=1,NMAX	106

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107 IPI=IREF+I
108 R(IRI)=ERT*R(IRI-1)*I
109 IF(NMIN)1200,3000,3000
110 680 IF(RT-9.)780,700,700
111 C RN BY DYATKINA N.GE.0,9.GE.RT.LT.80 AND -40.GT.RT.LE.-8
112 700 P1=1./RT
113 R2=ERT/RT
114 R(IREF)=R1-R2
115 IF(NMAX)760,76C,720
116 720 DO 740 I=1,NMAX
117 IRI=IREF+I
118 R(IRI)=R1*(I*R(IRI-1)-ERT)
119 GO TO 1160
120 IF(RT)56C,960,800
121 C ESTABLISH LIMIT FOR RN AS SERIES N.GE.0,0.GE.RT.LT.9
122 800 IF(RT-1.)820,860,960
123 820 NT=2.*DLOG10(ART)+13.
124 IF(NT-7)840,88C,880
125 840 NT=7
126 GO TO 880
127 NT=2.8*ART+19.-NMAX*(ART+6.)/18.
128 880 F1 = ERT / ( NMAX + 1. )
129 A1 = F1
130 K=NMAX+2
131 KK=NMAX+NT
132 DO 900 J=K,KK
133 F1=F1*RT/J
134 A1=A1+F1
135 R(IREF+NMAX)=A1
136 IF(NMAX)1160,116C,920
137 C RN BACK PECURENCE
138 920 NNM=-NMAX
139 LIM=-1
140 DO 940 I=NNM,LIM
141 IRI=IREF-I
142 940 R(IRI-1)=(RT*R(IRI)+ERT)/(-1)

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	GO TO 1160	143
960	IF(RT+8.)1080,1080,980	144
C	ESTABLISH LIMIT FOR RN AS SERIES N.GE.0,-8.GT.RT.LT.0	145
980	IF(RT+0.01)1020,1020,1000	146
1000	NT=6	147
	GO TO 880	148
1020	IF(RT+1.)1060,1060,1040	149
1040	NT=18.+6.*DLOG10(ART)	150
	GO TO 880	151
1060	NT=ART*(5.036-0.131*ART)+14.+0.1*(NMAX*RT-40.)	152
	GO TO 880	153
1080	IF(RT+40.)1100,1100,700	154
C	RN=PN N.GE.0,RT.LE.-40	155
1100	PP = - ERT / RT	156
	R(IREF) = PP	157
	IF(NMAX)1160,1160,1120	158
1120	RR=1./RT	159
	DO 1140 I=1,NMAX	160
	IRI=IREF+I	161
1140	R(IRI)=PP+RR*I*R(IRI-1)	162
C	FINIS OF RN,N.GE.0 START NOW ON N.LT.0	163
1160	IF(NMIN)1180,3000,3000	164
1180	IF(RT-40.)1300,1200,1200	165
C	RN=QN N.LT.0,RT.GE.40	166
1200	CALN=DLOG(ART)+0.577215664901533	167
	R(IREF-1)=-CALN	168
	IF(NMINA1)3000,3000,1220	169
1220	DO 1240 I=2,NMINA	170
	IRN=IREF-I	171
1240	R(IRN)=0.	172
	R2=1.	173
	F1=RT	174
	R(IREF-2)=F1*(CALN-R2)+R2	175
	IF(NMINA-2)3000,3000,1260	176
1260	DO 1280 I=3,NMINA	177
	IRI=IREF-I	178

R3 = 1.0+00 / ( I - 2. )	179
XI = 1.0+00 / ( I - 1. )	180
F1=F1*RTN*XI	181
R2=R2+XI	182
R( IRI )=P( IRI )+F1*(CALN-R2)+XI	183
DO 1280 J=1,NMINA	184
IRJ=IREF-J	185
R3=P3*RTN/(J-1+1. )	186
1280 R( IRJ )=R( IRJ )+R3	187
GO TO 3000	188
1300 IF( RT-2. )1500,1320,1320	189
C ESTABLISH LIMITS FOR E( CONTINUED FRACTION ) IN RN BY DYATKINA	190
C N.LT.0,2.GE.RT.LT.60	191
1320 IF( RT-6. )1360,1360,1340	192
1340 NT=-100./ART-6.	193
GO TO 1380	194
1360 NT=-24	195
1380 NL=-1	196
CF=0.	197
DO 1400 I=NT,NL	198
CC=1.+CF	199
1400 CF=-I*CC/(RT*CC+1-I)	200
CF=-ERT*(1.+CF)/(RT*(1.+CF)+1.)	201
CALNEI=0.577215664901533+DLOG(ART)-CF	202
R( IREF-1 )=-CALNEI	203
IF( NMINA1 )3000,3000,1420	204
1420 R( IREF-2 )=RT*(CALNEI-1.)+1.-ERT	205
IF( NMINA-2 )3000,3000,1440	206
1440 T1=1.	207
F1=RT	208
DO 1480 I=3,NMINA	209
T4 = 1.0+00 / ( I - 1. )	210
A = T4	211
XI = T4	212
F1=F1*RTN*XI	213
T1=T1+XI	214

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215 AA=1.
216 T3=0.
217 I2=I-2
218 DO 1460 J=1,12
219   RR=I-J-1.
220   A=A*RTN/RR
221   AA=AA*RTN/J
222   T4=T4+A
223   1460 T3=T3+AA/98
224   1480 R(IREF-I)=F1*(CALNEI-I1)+T3+XI-ERT*T4
225   GO TO 3000
226   1500 IF(RT+2.)1700,1700,1520
227   RN BY SERIES,ALL N,LT.0, -2.GT.RT.LT.2,RT.NE.0
228   1520 IF(ART-0.1)1540,1580,1580
229   1540 NT=10-DLOG10(ART)
230   IF(NT-5)1560,1600,1600
231   156C NT=5
232   GO TO 1600
233   158C NT=6.*ART+11.
234   1600 TA=1.
235   DC 1620 N=1,NMINA
236   TA=TA/N
237   1620 R(IREF-N)=RTN*TA
238   TA=RTN
239   DO 1640 I=2,NT
240     I1=I-1
241     TA=TA*RTN/I1
242     TB=TA/I
243     DO 1640 N=1,NMINA
244       IRN=IREF-N
245       TB=TB/(I1+N)
246     1640 R(IRN)=R(IRN)+TB
247     IF(NMINA)3000,3000,1660
248   C FINISH OFF RN BY SERIES
249   1660 TA=1.
250   DO 168C N=2,NMINA

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IFN=IREF-N	251
TA=TA*RTN	252
1680 R(IRN)=R(IRN)*TA	253
GO TO 3000	254
1700 IF(RT+60.)1980,1980,1760	255
C RN BY MIXED FORMULA FIRSTLY CALCULATE RN FOR MOST NEGATIVE N	256
C NEXT,CALCULATE QN IN R ARRAY,ALL N,N.LT.0	257
C THEN,PN=QN-RN FOR MOST NEGATIVE N,AND BACK RECURRENCE OF PN	258
C (STABLE) GIVES RN BY QN-PN	259
1760 NDEC=NMIN	260
IF(RT+4.)1800,1800,1780	261
1780 NDEC=NMIN/2	262
1800 NT=ART*(2.981-0.0215*ART)+19.+NDEC	263
F1 = 1.	264
A1 = 1.	265
DO 1820 I=2,NT	266
F1=F1*RTN/(NMINA1+I)	267
1820 A1=A1+F1/I	268
RR=A1*RTN**NMINA/8F(IREF)	269
IF(NMINA1)1940,1940,1860	270
1940 R(1)=RR	271
GO TO 3000	272
C PUT QN INTO R ARRAY	273
1860 DO 1880 I=1,NMINA	274
1880 R(I)=0.	275
CALN=DLOG(ART)+0.577215664901533	276
R(IREF-1)=-CALN	277
R2=1.	278
F1=RT	279
R(IREF-2)=F1*(CALN-R2)+R2	280
IF(NMINA-2)1940,1940,1900	281
1900 DO 1920 I=3,NMINA	282
IRI=IREF-I	283
R3 = 1.D+00 / ( I - 2. )	284
XI = 1.D+00 / ( I - 1. )	285
F1=F1*FTN*XI	286

R2=R2+XI	287
R(IRI)=R(IRI)+F1*(CALN-R2)+XI	288
DO 1920 J=1,NMINA	289
IRJ=IREF-J	290
R3=R3*RTN/(J-1+1.)	291
1920 R(IRJ)=R(IRJ)+R3	292
C BACK RECURRENCE FOR PN AS PP	293
1940 PP=R(1)-RR	294
R(1)=RR	295
DO 960 I=2,NMINA	296
IRN=(-1)EF	297
PP=(ERT+IRN*PP)/RT	298
1960 R(I)=R(I)-PP	299
GO TO 3000	300
C RN=PN N.LT.0, RT.LE.-60 PN BY CONTINUED FRACTION AND BACK REC.	301
1980 NT=-7	302
NL=-1	303
CF=0.	304
DO 2000 I=NT,NL	305
CC=1.+CF	306
2000 CF=-1*CC/(RT*CC+NMINA-I)	307
R(1)=-ERT*(1.+CF)/(RT*(1.+CF)+NMINA)	308
IF(NMINA1)3000,3000,2020	309
2020 DO 2040 I=2,NMINA	310
IRN=I-IREF	311
2040 R(I)=(-ERT+IRN*R(I-1))/RT	312
3000 RETURN	313
END	314

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000 DOUBLE PRECISION FUNCTION A ( IR , IS , IT , JT , MT )
001 IMPLICIT REAL * 8 ( A - H , O - Z )
002 COMMON/FACL/BF(100)
003 A=0.
004 INIT=0
005 DO 420 I2=INIT,IT
006 IR1=IR-I2
007 IF(IR1)440,20,20
008 IL3=IT-I2
009 F1=BF( I2+1)
010 DO 400 I3=INIT,IL3
011 IS1=IS-I3
012 IF(IS1)420,40,40
013 I1=IL3-I3
014 F2=F1*BF( I3+1)*BF( I1+1)
015 DO 380 J2=INIT,JT
016 IR2=IR1-J2
017 IF(IR2)400,60,60
018 JL3=JT-J2
019 F3=F2*BF( J2+1)
020 I3J2=I3+J2
021 DO 360 J3=INIT,JL3
022 IS2=IS1-J3
023 IF(IS2)380,80,80
024 J1=JL3-J3
025 F4=F3*BF( J3+1)*BF( J1+1)
026 DO 340 M2=INIT,MT
027 IR3=IR2-M2
028 IF(IR3)360,100,100
029 LM3=MT-M2
030 F5=F4*BF( M2+1)
031 DO 320 M3=INIT,LM3
032 IS3=IS2-M3
033 IF(IS3)340,120,120
034 LM4=LM3-M3

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F6=F5*BF(M3+1)
M23=M2+M3
DO 300 M4=INIT,LM4
IR4=IR3-M4
IF(IR4)320,14C,14C
IS4=IS3-M4
IF(IS4)320,160,160
LM5=LM4-M4
F7=F6*BF(M4+1)
M234=M23+M4
DO 280 M5=INIT,LM5
IR5=IR4-M5-M5
IF(IR5)300,180,280
LM6=LM5-M5
F8=F7*BF(M5+1)
DO 260 M6=INIT,LM6
IS5=IS4-M6-M6
IF(IS5)290,200,260
M1=LM6-M6
F9=F8*BF(M6+1)*BF(M1+1)
AA=2.*M234/F9
NEG=M234+13J2
IF(NEG/2#2-NEG)220,240,220
220 AA=-AA
240 A=A+AA
260 CONTINUE
280 CONTINUE
300 CONTINUE
320 CONTINUE
340 CONTINUE
360 CONTINUE
380 CONTINUE
400 CONTINUE
420 CONTINUE
440 A=A*BF(IT+1)*BF(JT+1)*BF(MT+1)/2.**((IT+JT+MT+MT)
IF(MT/2#2-MT)460,480,460

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460 A=-A  
480 RETURN  
END

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C		000
C	\$ \$ \$ \$ \$ \$ \$ \$ \$ \$ MAIN PROGRAM VANSS * * * * *	001
C		002
	IMPLICIT REAL*8(A-F,H,O-Z)	003
	DIMENSION RK(500),B(500),C(500),D(500),IOU(500),ALD(11),CD(44),VK(	004
	1500)	005
	DIMENSION SS(10),INDA(125),AFUN(4455)	006
	DIMENSION SUM(10),SOL1(10),SUM1(10),COJ(10),EK(10)	007
	DIMENSION SOLD(10),CDC(16),IEX(11)	008
	DIMENSION S (20)	009
	DIMENSION GER(20),SINT(2)	010
	DIMENSION FU(100) ,A(9),W(96),V(96)	011
	COMMON/CFUN/XXX(300),GD/FACL/BF(100)	012
	1/RIGS/ CON35,CON36,CCN38,CCN39,CON40,CON41,CON42,CON43,CON44,CON45	013
	1,CON46,CON47,CCN48,CON49,CON50,CON51,CCN52,CON53,CON54,CON55,CON56	014
	2,CON57,CON58,S,W,V,U,A,RAB,ALPHA1,ALPHA2,RS,FU,AI,AJ,AK,AX,AY,AZ,	015
	3BX,BY,BZ,CX,CY,CZ,AA,N,IE,IX	016
C		017
C	THIS PROGRAM COMPUTES POTENTIAL,KINETIC ENERGY,EXCHANGE,AND OVERLA	018
C	INTEGRALS BETWEEN FUNCTIONS OF S SYMMETRY FROM 1S THROUGH 4S	019
C		020
C		021
C	AT THIS POINT INSERT V(I) AND W(J) FROM MAIN PROGRAM VANDO,	022
C	LINES 026 THROUGH 217	023
C		024
C		025
C	READ OR WRITE THE FILE IDENTIFICATION ON THE TAPE	026
C	WRITE(3,16)	027
C	16 FORMAT(1X,'SS INTEGRALS VV')	028
C		029
	READ(3,16)(GER(N),N=1,4)	030
16	FORMAT(4A4)	031
	KJ=1	032
	READ(5,6)(GER(I),I=1,20)	033
6	FOPMAT(20A3)	034

C	READ NUMBER OF TERMS TO BE INCLUDED IN FOURIER SERIES	035
	READ(5,101)JJJJ	036
101	FORMAT(I5)	037
C	PARAMETER KKR	038
C	1 COMPUTE EXCHANGE ONLY	039
C	2 COMPUTE POTENTIAL ONLY	040
C	3 COMPUTE POTENTIAL AND EXCHANGE	041
	READ(5,101)KKR	042
C	READ A/2	043
	READ(5,641)QA	044
541	FORMAT(F9.6)	045
	GOTO (17,11,11),KKR	046
C	READ FOURIER COEFFICIENTS OF POTENTIAL	047
C	THE FOURIER SERIES SUM INCLUDES ONLY ONE VECTOR FROM EACH SET OF P	048
C	IOU IS THE NUMBER OF VECTORS OF THE TYPE (B,C,D)	049
	11 DO 400 J=1,JJJJ	050
400	READ(5,5)VK(J),B(J),C(J),D(J),IOU(J)	051
	5 FORMAT(1XE16.8,3F6.2,I5)	052
	IF(KKR.EQ.2)GOTO 402	053
C	READ FOURIER COEFFICIENTS OF EXCHANGE	054
	17 DO401 J=1,JJJJ	055
401	READ(5,5)RK(J),B(J),C(J),D(J),IOU(J)	056
402	CONTINUE	057
C	READ S WAVEFUNCTIONS FOR VANADIUM IN THE FOLLOWING ORDER	058
C	EXPONENT 1-S COEFF 2-S COEFF 3-S COEFF 4-S COEFF	059
C	22.77 0.20690816D 03 -0.64393167D 02 0.23074581D 02 -0.50849298D 01	060
C	36.05 0.91571364D 01 0.82262692D-01 0.18184385D 00 -0.19916231D 00	061
C	19.54 0.67068863D 02 -0.30786187D 03 0.12745228D 03 -0.31692523D 02	062
C	9.374 0.12892283D 01 0.32329185D 03 -0.13076502D 03 0.30935265D 02	063
C	7.905 -0.21432959D 01 0.73645055D 02 -0.14871312D 03 0.31669247D 02	064
C	5.130 0.71161337D 00 -0.30063086D 01 0.94460229D 02 -0.14181990D 02	065
C	3.902 -0.28160109D 00 0.13778162D 01 0.77294301D 01 -0.94140875D 01	066
C	3.75 0.92348347D-01 -0.41729371D 00 0.17225125D 02 0.13075490D 00	067
C	1.658 -0.15334215D-03 0.72289872D-03 0.10985870D-01 0.58007146D 00	068
C	0.995 0.11004437D-04 -0.50620409D-04 -0.51610808D-03 0.59293006D-01	069
C	0.676 -0.77644540D-06 0.36881156D-05 0.34745931D-04 0.53225332D-03	070

DO10 J=1,11	071
I=4*J-3	072
K=4*J-2	073
L=4*J-1	074
M=4*J	075
10 READ(5,3)ALD(J),CD(I),CD(K),CD(L),CD(M),IEX(J)	076
3 FORMAT(F8.5,4D16.8,15)	077
C READ DIRECT LATTICE COORDINATES	078
60 READ(5,4)AX,AY,AZ,BX,BY,BZ,CX,CY,CZ	079
4 FORMAT(9F8.5)	080
DO398J=1,10	081
SOL1(J)=0.0	082
SS(J)=0.00	083
398 SOLD(J)=0.00	084
C SUM OVER ALL DIFFERENT PAIRS OF EXPONENTS	085
DO40L=1,11	086
DO 40 K=1,L	087
CDC(1)=CD(4*K-3)*CD(4*L-3)	088
CDC(2)=CD(4*K-3)*CD(4*L-2)	089
CDC(3)=CD(4*K-3)*CD(4*L-1)	090
CDC(4)=CD(4*K-3)*CD(4*L)	091
CDC(5)=CD(4*K-2)*CD(4*L-2)	092
CDC(6)=CD(4*K-2)*CD(4*L-1)	093
CDC(7)=CD(4*K-2)*CD(4*L)	094
CDC(8)=CD(4*K-1)*CD(4*L-1)	095
CDC(9)=CD(4*K-1)*CD(4*L)	096
CDC(10)=CD(4*K)*CD(4*L)	097
CDC(11)=CD(4*K-2)*CD(4*L-3)	098
CDC(12)=CD(4*K-1)*CD(4*L-3)	099
CDC(13)=CD(4*K) *CD(4*L-3)	100
CDC(14)=CD(4*K-1)*CD(4*L-2)	101
CDC(15)=CD(4*K) *CD(4*L-2)	102
CDC(16)=CD(4*K) *CD(4*L-1)	103
IF(L.EQ.K)GOTO 50	104
CDC(1) =2.0*CDC(1)	105
CDC(2) =CDC(2)+CDC(11)	106

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CDC(3) =CDC(3)+CDC(12)
CDC(4) =CDC(4)+CDC(13)
CDC(5) =2.0*CDC(5)
CDC(6) =CDC(6)+CDC(14)
CDC(7) =CDC(7)+CDC(15)
CDC(8) =2.0*CDC(8)
CDC(9) =CDC(9)+CDC(16)
CDC(10)=2.0*CDC(10)
50 ALPHA1=ALD(L)
ALPHA2=ALD(K)
IE=IEX(K)
IX=IEX(L)
CON35=ALPHA1*ALPHA1/3.0
CON36=ALPHA2*ALPHA2/3.0
CON38=ALPHA1*ALPHA1
CON39=ALPHA2*ALPHA2
CON40=CON38*CON39
CON41=CON40*ALPHA1
CON42=CON38*ALPHA1
CON43=CON39*ALPHA1
CON44=CON40*ALPHA2
CON45=CON42*ALPHA2
CON46=CON39*ALPHA2
CON47=CON46*ALPHA1
CON48=ALPHA1*ALPHA2
CON49=CON46/ALPHA1
CON50=CON42/ALPHA2
CON51=ALPHA1/ALPHA2
CON52=ALPHA2/ALPHA1
CON53=1.0/ALPHA1
CON54=1.0/ALPHA2
CON55=1.0/CON48
CON56=CON39*CON53
CON57=CON51*ALPHA1
CON58=CON48*ALPHA1
PAB=DSQRT((HX-AX)**2+(BY-AY)**2+(BZ-AZ)**2)

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	AA=C.5D0	*ALPHA1*ALPHA2*RAB**3	143
	DO100N=1,10		144
	SUM1(N)=0.0		145
100	SUM(N)=0.00		146
C	SUM THE FOURIER SERIES		147
	DO500 J=1,JJJJ		148
	DO60KOL=1,100		149
90	FU(KOL)=0.00		150
	N=100(J)		151
	AI=3.14159265/QA*B(J)		152
	AJ=3.14159265/QA*C(J)		153
	AK=3.14159265/QA*D(J)		154
	SK= DSQRT(AI*AI+AJ*AJ+AK*AK)		155
	RS=SK*SK		156
C			157
C	SUBROUTINES D3INT,AGET,AND INTEG COMPUTE THE THREE-CENTER INTEGRAL		158
C	IF THE INTEGRALS HAVE BEEN STORED ON TAPE,READ THE TAPE INSTEAD OF		159
C	CALLING D3INT		160
C			161
C	CALL D3INT		162
C	SINT(KJ)=S(1)		163
C	KJ=KJ+1		164
C	IF(KJ.NE.3)GOTO 631		165
C	WRITE(3)SINT(1),SINT(2)		166
C	KJ=1		167
C	IF(KJ.EQ.2)GOTO 819		168
C	READ INTEGRALS IN PAIRS FROM TAPE		169
	READ(3) SINT(1),SINT(2)		170
	KJ=1		171
819	S(1)=SINT(KJ)		172
	KJ=KJ+1		173
631	DO 777N=1,9		174
777	S(N+1)=S(1)		175
	DO146N=1,10		176
	IF(J.NE.1)GOTO145		177
	SS(N)=SS(N)+S(N)*CDC(N)		178

145	GOTO (147,148,148),KKR	179
	STOP	180
149	SUM1(N)=SUM1(N)+RK(J)*S(N)	181
	IF(KKR.EQ.2)GOTO 146	182
147	SUM(N)=SUM(N)+VK(J)*S(N)	183
146	CONTINUE	184
500	CONTINUE	185
	DO4C8J=1,10	186
	SOL1(J)=SOL1(J)+SUM1(J)*CDC(J)	187
408	SOLD(J)=SUM(J)*CDC(J)+SELD(J)	188
40	CONTINUE	189
	WRITE(6,2)AX,AY,AZ	190
	WRITE(6,2)BX,BY,BZ	191
	WRITE(6,2)CX,CY,CZ	192
2	FORMAT(1X,3F10.6)	193
C		194
C	COMPUTE KINETIC ENERGY INTEGRALS	195
C	SEE SUBROUTINE SZE FOR DESCRIPTION OF THE PARAMETERS	196
C		197
	LA=0	198
	LB=0	199
	MB=0	200
	MA=0	201
	CA=23.	202
	CB=23.	203
	CALL FACGEN(RF)	204
	CALL AGENT(INDA,AFUN)	205
	OD=RA8	206
	DO700 J=1,10	207
700	EK(J)=0.D0	208
	DO200 J=1,11	209
	DO200 K=1,J	210
	XA=ALD(J)	211
	XB=ALD(K)	212
	IX=IEX(J)	213
	IE=IEX(K)	214



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215 NA=IX+1
216 N8=IF+1
217 CCJ(1)=0.2500/(XA*XB)**1.5
218 CCJ(2)=0.43301270/(XA**2.5*XB**1.5)
219 CCJ(3)=1.18585412/(XA**3.5*XB**1.5)
220 CCJ(4)=4.43705983/(XA**4.5*XB**1.5)
221 CCJ(5)=7.68521307/(XA**4.5*XB**2.5)
222 CCJ(6)=21.04682276/(XA**4.5*XB**3.5)
223 CCJ(7)=78.7500/(XA*XB)**4.5
224 CCJ(8)=2.05395953/(XA**3.5*XB**2.5)
225 CCJ(9)=5.62500/(XA*XB)**3.5
226 CCJ(10)=0.7500/(XA*XB)**2.5
227 IF(IX.NE.0.OR.IE.NE.0) GOTO 2000
228 ABK=2.00
229 I=1
230 GOTO 749
231 2000 IF(IX.NE.1.OR.IE.NE.0) GOTO 2001
232 ABK=2.00/DSQRT(3.00)
233 I=2
234 GOTO 749
235 2001 IF(IX.NE.2.OR.IE.NE.0) GOTO 2002
236 ABK=4.00/(3.00*DSQRT(10.00))
237 I=3
238 GOTO 749
239 2002 IF(IX.NE.3.OR.IE.NE.0) GOTO 2003
240 ABK=2.00/(3.00*DSQRT(35.00))
241 I=4
242 GOTO 749
243 2003 IF(IX.NE.1.OR.IE.NE.1) GOTO 2005
244 ABK=2.00/3.00
245 I=10
246 GOTO 749
247 2005 IF(IX.NE.2.OR.IE.NE.1) GOTO 2006
248 ABK=4.00/(3.00*DSQRT(30.00))
249 I=8
250 GOTO 749

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2006 IF(IX.NE.3.OR.IE.NE.1) GOTO 2007
ABK=2.D0/(3.D0*DSQRT(105.D0))
I=5
GOTO 749
2007 IF(IX.NE.2.OR.IE.NE.2)GOTO 2010
ABK=4.D0/45.D0
I=9
GOTO 749
2010 IF(IX.NE.3.OR.IE.NE.2)GOTO 2011
ABK=4.D0/(45.D0*DSQRT(14.D0))
I=6
GOTO 749
2011 ABK=2.D0/315.D0
I=7
749 CALL SZE(E,ZA,ZB,NA,LA,MA,XA,CA,NB,LB,MB,XB,CB,ABK,INDA,AFUN)
E=E*CDJ(I)*2.0
CDC(1)=CD(4*K-3)*CD(4*J-3)
CDC(2)=CD(4*K-3)*CD(4*J-2)
CDC(3)=CD(4*K-3)*CD(4*J-1)
CDC(4)=CD(4*K-3)*CD(4*J)
CDC(5)=CD(4*K-2)*CD(4*J-2)
CDC(6)=CD(4*K-2)*CD(4*J-1)
CDC(7)=CD(4*K-2)*CD(4*J)
CDC(8)=CD(4*K-1)*CD(4*J-1)
CDC(9)=CD(4*K-1)*CD(4*J)
CDC(10)=CD(4*K)*CD(4*J)
CDC(11)=CD(4*K-2)*CD(4*J-3)
CDC(12)=CD(4*K-1)*CD(4*J-3)
CDC(13)=CD(4*K)*CD(4*J-3)
CDC(14)=CD(4*K-1)*CD(4*J-2)
CDC(15)=CD(4*K)*CD(4*J-2)
CDC(16)=CD(4*K)*CD(4*J-1)
IF(K.EQ.J)GOTO 710
CDC(1)=2.0*CDC(1)
CDC(2)=CDC(2)+CDC(11)
CDC(3)=CDC(3)+CDC(12)

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CDC(4) =CDC(4)+CDC(13)	287
CDC(5) =2.0*CDC(5)	288
CDC(6) =CDC(6)+CDC(14)	289
CDC(7) =CDC(7)+CDC(15)	290
CDC(8) =2.0*CDC(8)	291
CDC(9) =CDC(9)+CDC(16)	292
CDC(10)=2.0*CDC(10)	293
710 DO701 L=1,10	294
701 EK(L)=EK(L)+E*CDC(L)	295
200 CONTINUE	296
N1=AX/QA+0.1	297
N2=AY/QA+0.1	298
N3=AZ/QA+0.1	299
PUNCH 108,N1,N2,N3	300
108 FORMAT(3I5)	301
DO407J=1,10	302
IF(KKR.NE.3)GOTO 28	303
PUNCH 107,SGLD(J),EK(J),SOL1(J),SS(J)	304
107 FORMAT(4E16.8)	305
28 KK=2 *J-1	306
WRITE(6,8)GER(KK),GER(KK+1)	307
8 FORMAT(1X,2A3)	308
WRITE(6,1)SGLD(J),EK(J),SOL1(J),SS(J)	309
1 FORMAT(1X,'POTENTIAL=',E16.8,' KINETIC ENERGY=',E16.8,' EXCHANGE='	310
1,E16.8,' OVERLAP=',E16.8)	311
WRITE(6,974)	312
974 FORMAT(1X,1H )	313
407 CONTINUE	314
GO TO 60	315
END	316

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000  SUBROUTINE AGET (EAX)
001  IMPLICIT REAL*8(A-H,O-Z)
002  DIMENSION S (20)
003  DIMENSION FU(100) , A(9),W(96),V(96)
004  COMMON/BIGS/CON35,CCN36,CCN38,CCN39,CON40,CON41,CON42,CON43,
005  1CON44,CON45
006  1,CON46,CON47,CCN48,CCN49,CON50,CON51,CON52,CON53,CON54,CON55,CON56
007  2,CON57,CON58,S,W,V,U,A,PAB,ALPHA1,ALPHA2,RS,FU,AI,AJ,AK,AX,AY,AZ,
008  3BX,BY,BZ,CX,CY,CZ,AA,N,IE,IX
009  F=U*(1.-U)*PAB**2
010  G=RS+ALPHA2**2/(1.-U)+ALPHA1**2/U
011  FG=F*G
012  DO8=DSOPT(FG)
013  O1=DO8**3
014  O2=DO8*O1
015  O3=O2*DO8
016  O4=O1*O1
017  O5=O1*O2
018  O6=O2*O2
019  O7=O1*O4
020  O8=O3*O3
021  O9=O1*O6
022  O10=O4*O4
023  O11=O4*O5
024  O12=O5*O5
025  O13=O5*O6
026  O14=O6*O6
027  O15=O6*O7
028  A(1)=F*(3./O3+3./O2+1./O1)
029  A(2)=F*F*(15./O5+15./O4+6./O3+1./O2)
030  A(3)=F**3*(105./O7+105./O6+45./O5+10./O4+1./O3)
031  A(4)=F**4*(945./O9+945./O8+420./O7+105./O6+15./O5+1./O4)
032  A(5)=F**5*(10395./O11+10395./O10+4725./O9+1260./O8+210./O7+21./O6
033  1+1./O5)
034  A(6)=F**6*(135135./O13+135135./O12+62370./O11+17325./O10+3150./O9

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1+378./D8+28./D7+1./D61
A(7)=F**7*(2027025./D15+2027025./D14+945945./D13+270270./D12+
151975./D11+6930./D10+630./D9+36./D8+1./D7)
IF(D08.GT.90.)GOTO 222
EAX=DEXP(-D08)
GOTO 223
222 EAX=0.0
223 RETURN
END
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SUBROUTINE INTEG
  IMPLICIT REAL*8(A-H,C-Z)
  DIMENSION S (20)
  DIMENSION FU(100) , A(9),%(96),V(96)
  COMMON/HIGS/CON35,CON36,CON38,CON39,CCN40,CON41,CON42,CON43,
1CON44,CON45
  1,CON46,CON47,CCN48,CCN49,CCN50,CON51,CON52,CON53,CON54,CON55,CON56
  2,CON57,CON58,S,W,V,U,A,RAB,ALPHA1,ALPHA2,RS,FU,AI,AJ,AK,AX,AY,AZ,
  3BX,BY,BZ,CX,CY,CZ,AA,N,IE,IX
  DO 50 J=1,96
    S(1)=S(1)+W(J)*FU(J)
    S(1)=S(1)*0.500
  RETURN
  END
50

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	SUBROUTINE D3INT	000
	IMPLICIT REAL*8(A-H,C-Z)	001
	DIMENSION S (20)	002
	DIMENSION FU(100),A(9),W(96),V(96)	003
	COMMON/BIGS/CON35,CON36,CON38,CON39,CON40,CON41,CON42,CON43,	004
	1CON44,CON45	005
	1,CON46,CON47,CON48,CON49,CCN50,CCN51,CCN52,CCN53,CCN54,CCN55,CCN56	006
	2,CCN57,CCN58,S,X,V,U,A,RAB,ALPHA1,ALPHA2,RS,FU,AI,AJ,AK,AX,AY,AZ,	007
	3BX,BY,BZ,CX,CY,CZ,AA,N,IE,IX	008
	S(1)=0.DO	009
	TUM1=CON42*CON46	010
	TUM2=CON46*ALPHA1	011
	CON=N	012
206	DO 750 J=1,96	013
	U=(V(J)+1.DO)*C.5D0	014
	DX=AX*U+BX*(1.-U)	015
	DY=AY*L+BY*(1.-U)	016
	DZ=AZ*U+BZ*(1.-U)	017
	RCDX=DX-CX	018
	RCDY=DY-CY	019
	RCDZ=DZ-CZ	020
	C1=DCOS(AI*RCDX)	021
	C2=DCOS(AJ*RCDY)	022
	C3=DCOS(AK*RCDZ)	023
	CO=C1*C2*C3*CON	024
	CAB=1.0-U	025
	HH=U*(1.0-U)	026
	HHH=HH*HH	027
	HHHH=HHH*HH	028
	CAB2=U*U	029
	CAB3=CAB*CAB	030
	CAB4=U*HH	031
	CAB5=CAB*HH	032
	CAB6=HHH*U	033
	CAB7=CAB2*HH	034

	CAB8=CAB3*HH	035
	CAB9=CAB2*U	036
	CAB10=CAB3*CAB	037
	CALL AGET(EAX)	038
	IF(EAX.EQ.0.0)GOTO750	039
C		040
C	FOLLOWING ARE THE 10 DIFFERENT S-S INTEGRAL FORMULAS FOR 1S THROUG	041
C		042
	IF(IX.NE.0.OR.IE.NE.0) GOTO 2000	043
	FU(J)=A(1)*CO*EAX*AA	044
	GOTO 750	045
2000	IF(IX.NE.1.OR.IE.NE.0) GOTO 2001	046
	FU(J)=(-CON53*A(1)+ALPHA1/U*A(2))*CO*EAX*AA	047
	GOTO 750	048
2001	IF(IX.NE.2.OR.IE.NE.0) GOTO 2002	049
	FU(J)=(-3./U*A(2)+CON38/CAB2*A(3))*CO*EAX*AA	050
	GOTO 750	051
2002	IF(IX.NE.3.OR.IE.NE.0) GOTO 2003	052
	FU(J)=(3.*CON53/U*A(2)-6.*ALPHA1/CAB2*A(3)+CON42*A(4)/CAB9)	053
	1*EAX*CO*AA	054
	GOTO 750	055
2003	IF(IX.NE.1.OR.IE.NE.1) GOTO 2004	056
	FU(J)=(A(1)-(CON39/CAB8+CON38/U)*A(2)+CON40/HH*A(3))*EAX*CO*CON55	057
	1*AA	058
	GOTO 750	059
2004	IF(IX.NE.2.OR.IE.NE.1) GOTO 2005	060
	FU(J)=(3./U*A(2)-(CCN38/CAB2+3.*CON39/HH)*A(3)+CON40/CAB4*A(4))*	061
	1CON54*CO*EAX*AA	062
	GOTO 750	063
2005	IF(IX.NE.3.OR.IE.NE.1) GOTO 2006	064
	FU(J)=(-3.*CON55*A(2)/U +(6.*CON51/CAB2+3.*CON52/HH)*A(3)-(CON50/	065
	1CAB9+6.*CON48/CAB4)*A(4)+CON45/CAB7*A(5))*CO*EAX*AA	066
	GOTO 750	067
2006	IF(IX.NE.2.OR.IE.NE.2)GOTO 2007	068
	FU(J)=(9./HH*A(3)+CCN40/HHH*A(5)-3.*(CCN38/CAB4+CON39/CAB5)*A(4))	069
	1*EAX*CO*AA	070



GOTO 750	071
2007 IF(IX.NE.3.UR.IE.NE.2)GOTO 2008	072
FU(J)=(CCN41*A(6)/CAB6-16.*CON43/HH+3.*CON42/CAB7)*A(5)+(18.*	073
1ALPHA1/CAB4+3.*CON39+CON53/CAB5)*A(4)-9.*CON53/HH*A(3))*CO*EAX*AA	074
GOTO 750	075
2008 FU(J)=(TCM1 /HHH*A(7)-6.*(CCN45/CAB6+TOM2 /(HHH*CAB))*A(6)+(36.	076
1*CON48/HH+3.*CON49/CAB8+3.*CON50/CAB7)*A(5)-18.*(CON51/CAB4+CON52	077
2/CAB5)*A(4)+9.*A(3)/(HH*CCN48))*CO*EAX*AA	078
750 CONTINUE	079
CALL INTEG	080
RETURN	081
END	082
C	083
C	084
C	085
C	086
C	087
C	088
C	089

AT THIS POINT INSERT THE KINETIC ENERGY GROUP OF SUBROUTINES AS LISTED EARLIER. THESE ARE SIZE, FACGEN, AGENT, PORN, AND DOUBLE PRECISION FUNCTION.

## VITA

Thomas E. Norwood was born on January 9, 1943 in Montgomery, Alabama. He graduated from Robert E. Lee High School, Montgomery, Alabama, in 1961. He received the degree of Bachelor of Arts in Physics from Vanderbilt University, Nashville, Tennessee in 1965. He is now a candidate for the degree of Doctor of Philosophy in the Department of Physics and Astronomy.

## EXAMINATION AND THESIS REPORT

Candidate: Thomas Ellis Norwood, III

Major Field: Physics

Title of Thesis: Energy Bands of Vanadium Oxide

Approved:

John L. Fry  
Major Professor and Chairman

Max Goodrich  
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Date of Examination:

7 October 1969